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and Astronomers Accelerator School

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Ultrafast Laser Technology and Synchronization

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outline

- **What is a laser?**
- What is an ultrafast laser?
- Mathematic description of an ultrafast laser pulse;
- Ultrafast lasers for S³FEL;
- Femtosecond timing and synchronization;

LASER

Light Amplification by Stimulated Emission of Radiation

All kinds of sources of light



sunlight



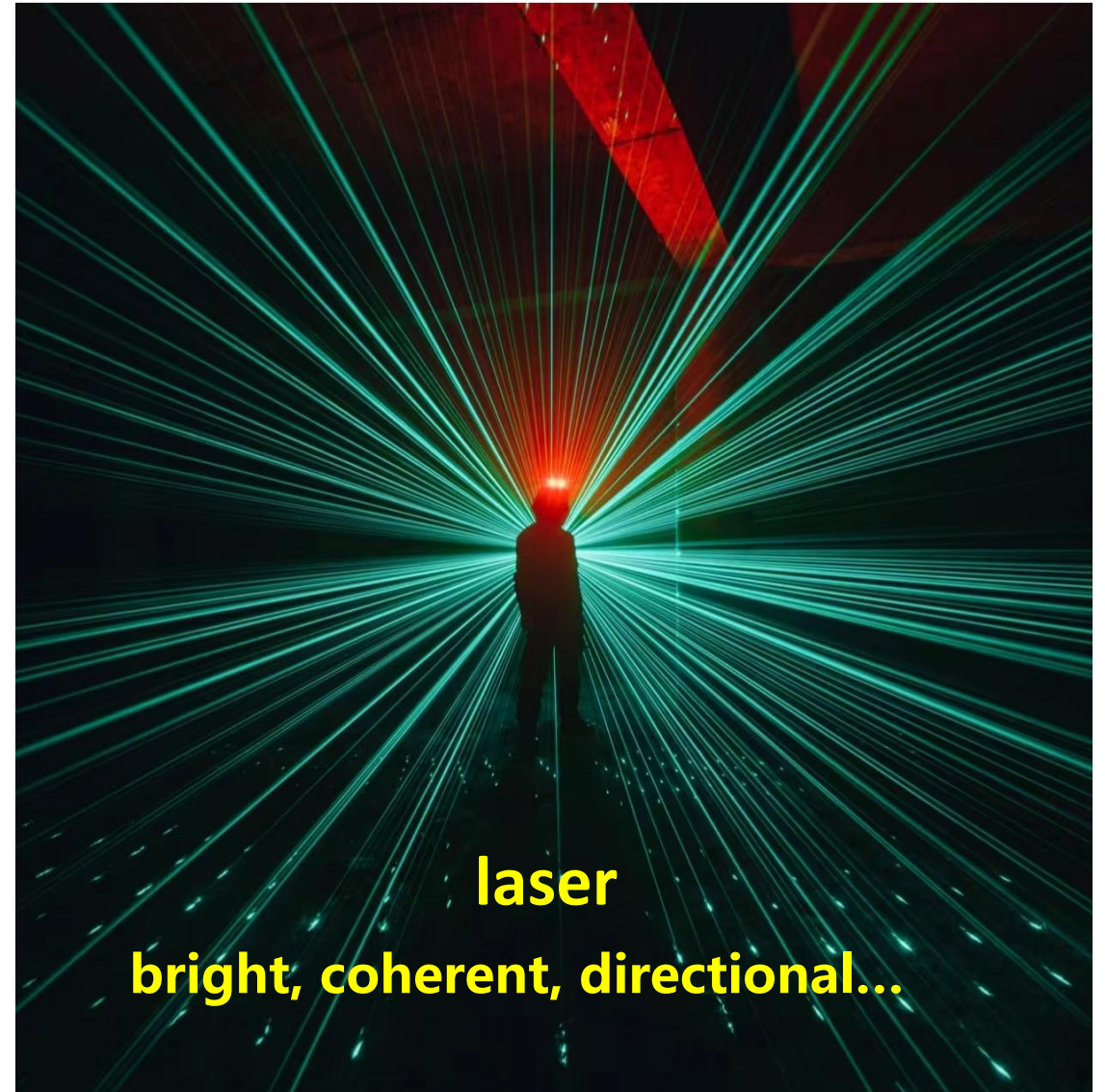
moonlight



lamp



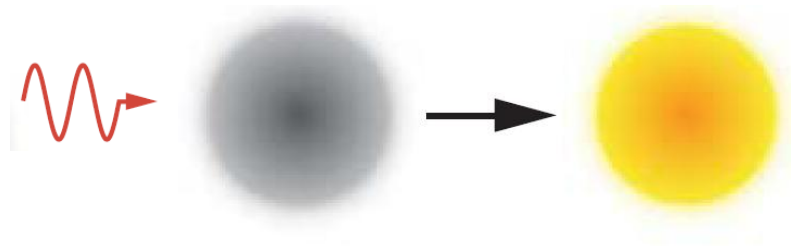
fire



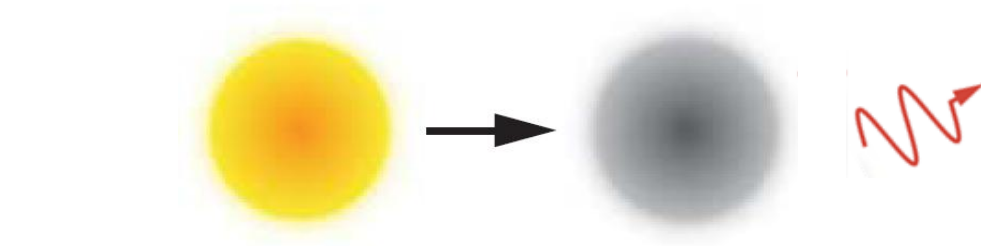
laser

bright, coherent, directional...

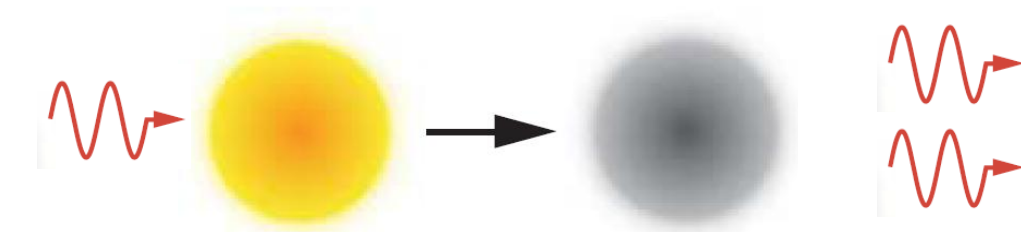
The basic principle of laser: interaction of light with matter



Absorption



Spontaneous emission

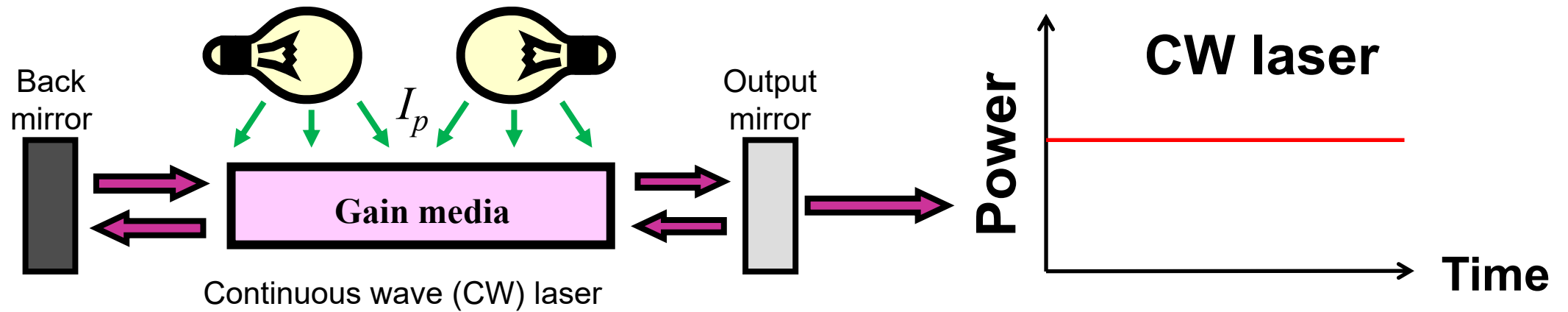


Stimulated emission

1917: *on the quantum theory of radiation* – Einstein's paper

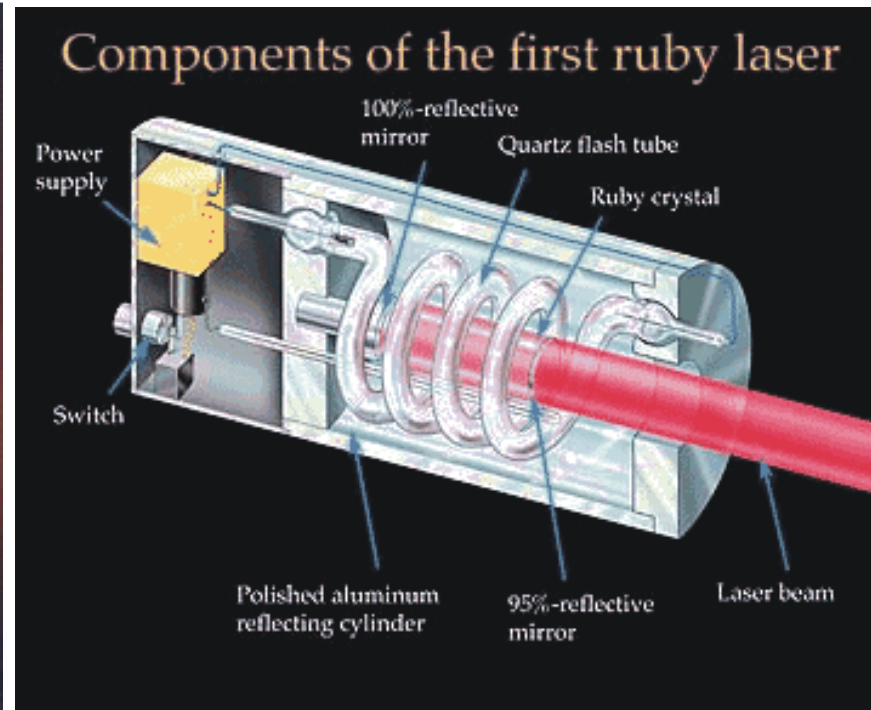
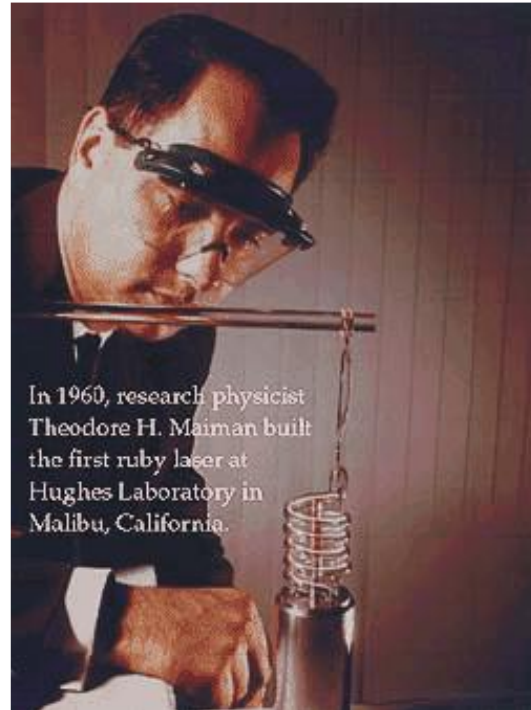
THE LASER

Three key elements to implement a laser



- **PUMP:** provide energy;
- **GAIN MEDIA:** storage of the energy provided by pump through population inversion;
- **RESONANT CAVITY:** Light bounces back and forth within a cavity, gaining energy with each pass through the gain medium and transmitting part of the energy out of the cavity in the form of laser.

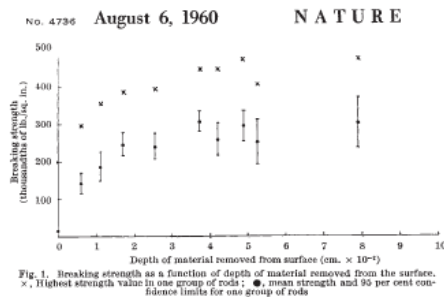
Invention of laser



- May 16, 1960: first coherent optical beam generated from the LASER constructed by Theodore Maiman working at Hughes Laboratory.
- Submitted the results to PRL on June 22nd, and got rejected on 24th.
- July 7th, Public demonstration to the press at Hotel Delmonico in New York.

Invention of laser

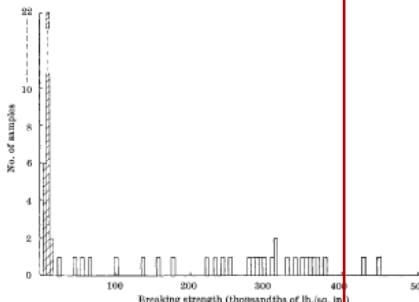
The most
valuable
paper ever
published
in Nature



particular etching solution removes material from a glass surface it is possible to study the strength of etched specimens as a function of the depth of material removed from the surface. Such a study may give some information about the size and nature of the surface imperfections.

Commercially available soda-glass rods, of 6.8 mm. diameter, have been etched and broken in four point bending over a constant bending moment span of 1 in. The rod diameters and loads at fracture were measured and the breaking stresses calculated using the simple bending formula. Groups of rods (containing 16-32 rods) were given different periods of etching, and the depth of material removed from the surface of the rods was calculated for each group.

The variation of the mean breaking strength of these groups of rods, with depth of material removed from the surface, is shown in Fig. 1. Also shown on Fig. 1 are the 95 per cent confidence limits on the mean strength and the highest strength value recorded in each group of rods. Fig. 2 is a histogram comparing the distribution of breaking stresses for a group of rods which have been etched for 40 min. with that for unetched rods. The maximum strengths



obtained with these bulk glass specimens are in the region 450,000-500,000 lb./sq. in. and closely approach the value obtained by Thomas³ for fine glass fibres.

The glass rod used in these experiments had the following approximate composition by weight (percentages): SiO_2 , 69; Na_2O , 16; CaO , 4; Al_2O_3 , 3; MgO , 3. The etching solution contained about 15 per cent hydrofluoric acid, 15 per cent sulphuric acid by weight and the remainder water.

Experiments are being continued to determine the effect on these results of varying the concentration, temperature and nature of the etchant; and of changing the size and composition of the glass.

B. A. PROCTOR

Rolls-Royce, Ltd.,
Aerophysics Laboratory,
Littleover,
Derby.
June 22.

¹ Greene, C. H., *J. Amer. Cer. Soc.*, **39**, 66 (1956).

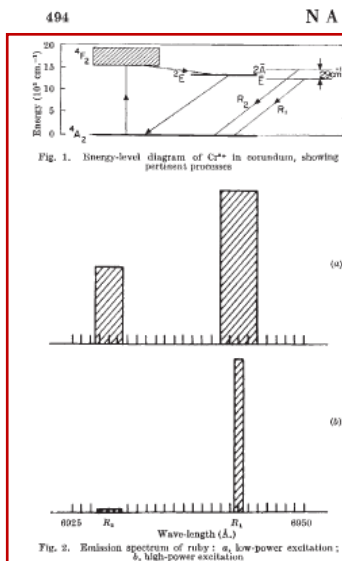
² Thomas, W. F., *Nature*, **181**, 1006 (1958); *Phys. and Chem. Glasses*, **1**, 4 (1960).

Stimulated Optical Radiation in Ruby

Schawlow and Townes¹ have proposed a technique for the generation of very monochromatic radiation in the infra-red optical region of the spectrum using an alkali vapour as the active medium. Javan² and Sanders³ have discussed proposals involving electron-excited gaseous systems. In this laboratory an excited gaseous system has been successfully applied to a fluorescent solid resulting in the attainment of negative temperatures and stimulated optical emission at a wave-length of 6943 Å.; the active material used was ruby (chromium in corundum).

A simplified energy-level diagram for triply ionized chromium in this crystal is shown in Fig. 1. When this material is irradiated with energy at a wave-length of about 5500 Å., chromium ions are excited to the 4F_3 state and then quickly lose some of their excitation energy through non-radiative transitions to the 2E non-radiative state. This state then slowly decays by spontaneously emitting a sharp doublet the components of which at 300° K. are at 6943 Å. and 6929 Å. (Fig. 2a). Under very intense excitation the population of this metastable state (2E) can become greater than that of the ground-state; this is the condition for negative temperatures and consequently amplification via stimulated emission.

To demonstrate the above effect a ruby crystal of 1 cm. dimensions coated on two parallel faces with silver was irradiated by a high-power flash lamp;



the emission spectrum obtained under these conditions is shown in Fig. 2b. These results can be explained on the basis that negative temperatures were produced and regenerative amplification ensued. I expect, in principle, a considerably greater ($\sim 10^3$) reduction in line width when mode selection techniques are used. I gratefully acknowledge helpful discussions with C. Birnbaum, R. W. Hellwarth, L. C. Levitt, and R. A. Satten and am indebted to I. J. D'Huysens and C. K. Asawa for technical assistance in obtaining the measurements.

T. H. MAIMAN

Hughes Research Laboratories,
A Division of Hughes Aircraft Co.,
Malibu, California.

¹ Schawlow, A. L., and Townes, C. H., *Phys. Rev.*, **112**, 1940 (1958).
² Javan, A., *Phys. Rev. Letters*, **3**, 87 (1959).
³ Sanders, J. H., *Phys. Rev. Letters*, **3**, 56 (1959).
⁴ Maiman, T. H., *Phys. Rev. Letters*, **4**, 564 (1960).

METALLURGY

A Simple Method of Investigating the Creep of Metals under Simple Shear

DR. K. IL JOLLIFFE and I¹ have investigated the creep of metals under simple shear by the use of a disk of the metal in question, in which is cut a concentric circular annulus, the metal external to the annulus being securely gripped, while that internal to the annulus is subjected to a constant torque. We

have shown that the behaviour of the metal in these circumstances is in many ways simpler and more informative than that exhibited by wires or rods under tension.

From the point of view of the industrial study of creep, the method has the disadvantage that the specimens are somewhat troublesome to prepare and measure. To get over this difficulty I have devised a method in which the specimen has the form shown in Fig. 1. In the plate $ABCD$ (Fig. 1a) are cut rectangular grooves $MNOP$, $QRST$, as shown in cross-section in Fig. 1b. The metal $ABNM$, $TSCD$ is securely held, and a force F applied to the metal $PORQ$ in a direction parallel to the grooves. Under these conditions the shear stress distribution in the rectangular plates $MNOP$, $QRST$ is not strictly uniform, as it is in the disk method, which is clear from the fact that the shear stress over the free ends must be zero. The distribution has been worked out by me² and by C. E. Inglis³.

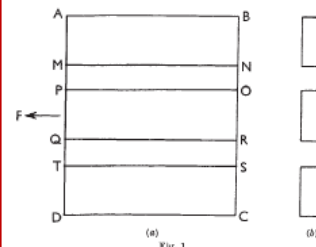


Fig. 1

Mr. D. B. Gilding has been working under my direction on the best form of the plate and has established that if the ratio of MN to NO is in the region of 7, the results on creep obtained with the disposition described correspond closely to those obtained by the method of Andrade and Jolliffe⁴. It seems possible that the new method may be of use in further investigations of creep and may also have applications to the problem of fatigue.

E. N. DA C. ANDRADE

Department of Metallurgy,
Imperial College of Science and Technology,
London, S.W.7.

¹ Andrade, E. N. da C., and Jolliffe, K. H., *Proc. Roy. Soc., A*, **213**, 3 (1952); **254**, 291 (1960).
² Andrade, E. N. da C., *Proc. Roy. Soc., A*, **85**, 448 (1911).
³ Inglis, C. E., *Proc. Roy. Soc., A*, **105**, 539 (1923).

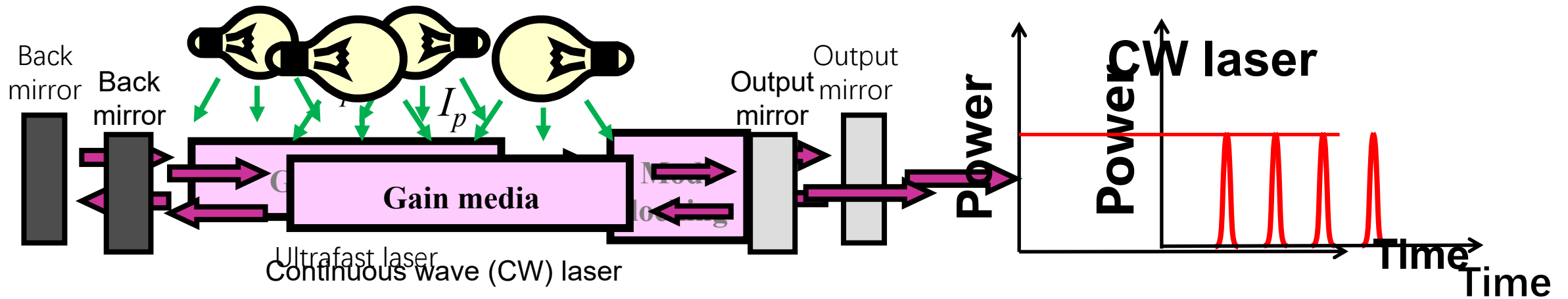
An Improvement in the Ductility of Beryllium at High Temperatures

THE outstanding problem in beryllium metallurgy is the lack of ductility exhibited by the metal, both at room temperature and at elevated temperatures. At room temperature, brittleness can be attributed to the cause of cleavage of basal planes of the hexagonal lattice, and to the high yield-strength of the prismatic planes, while at temperatures above 400° C., intergranular failure predominates. A ductility maximum occurs at intermediate temperatures depending on the strain-rate used, tensile specimens failing with

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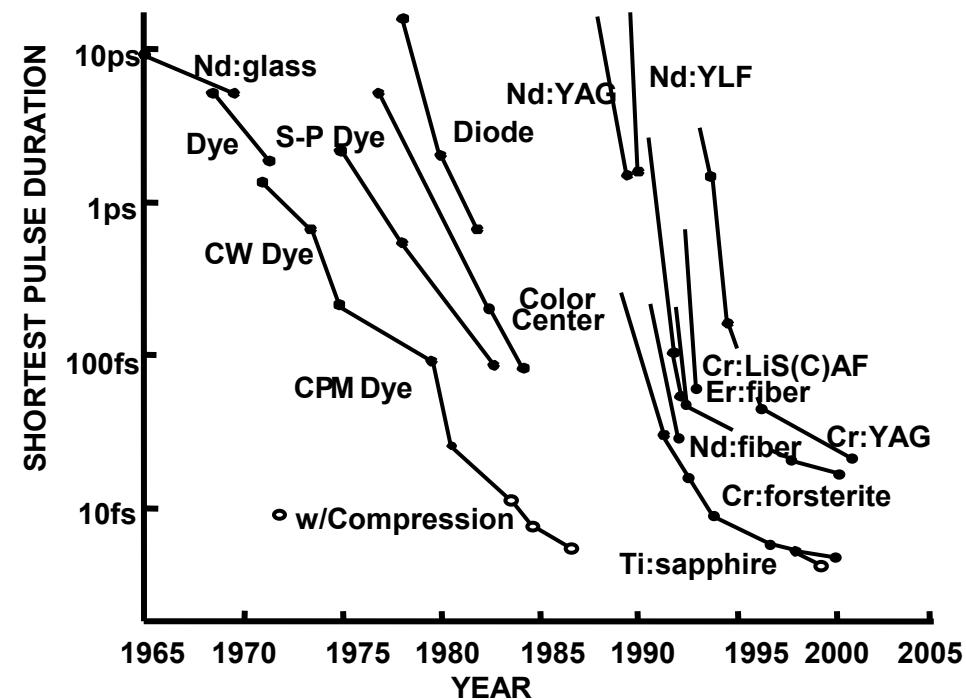
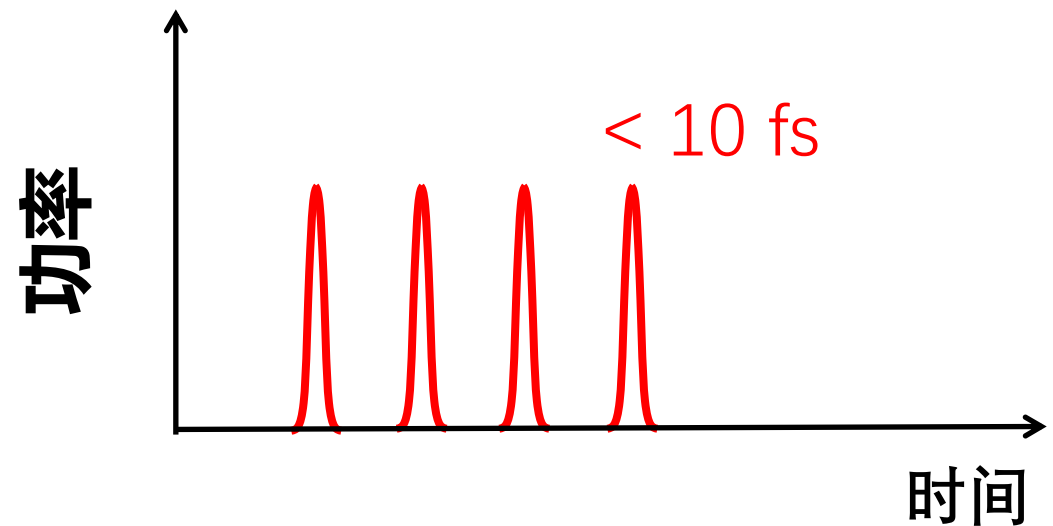
Generation of ultrafast laser pulses: mode locking



Four key elements to implement an ultrafast laser

- **PUMP:** provide energy;
- **GAIN MEDIA:** storage of the energy provided by pump through population inversion;
- **RESONANT CAVITY:** Light bounces back and forth within a cavity, gaining energy with each pass through the gain medium and transmitting part of the energy out of the cavity in the form of laser;
- **MODE LOCKING:** generate phase relation between different longitudinal modes;

Characters of ultrafast laser pulses: ultrashort

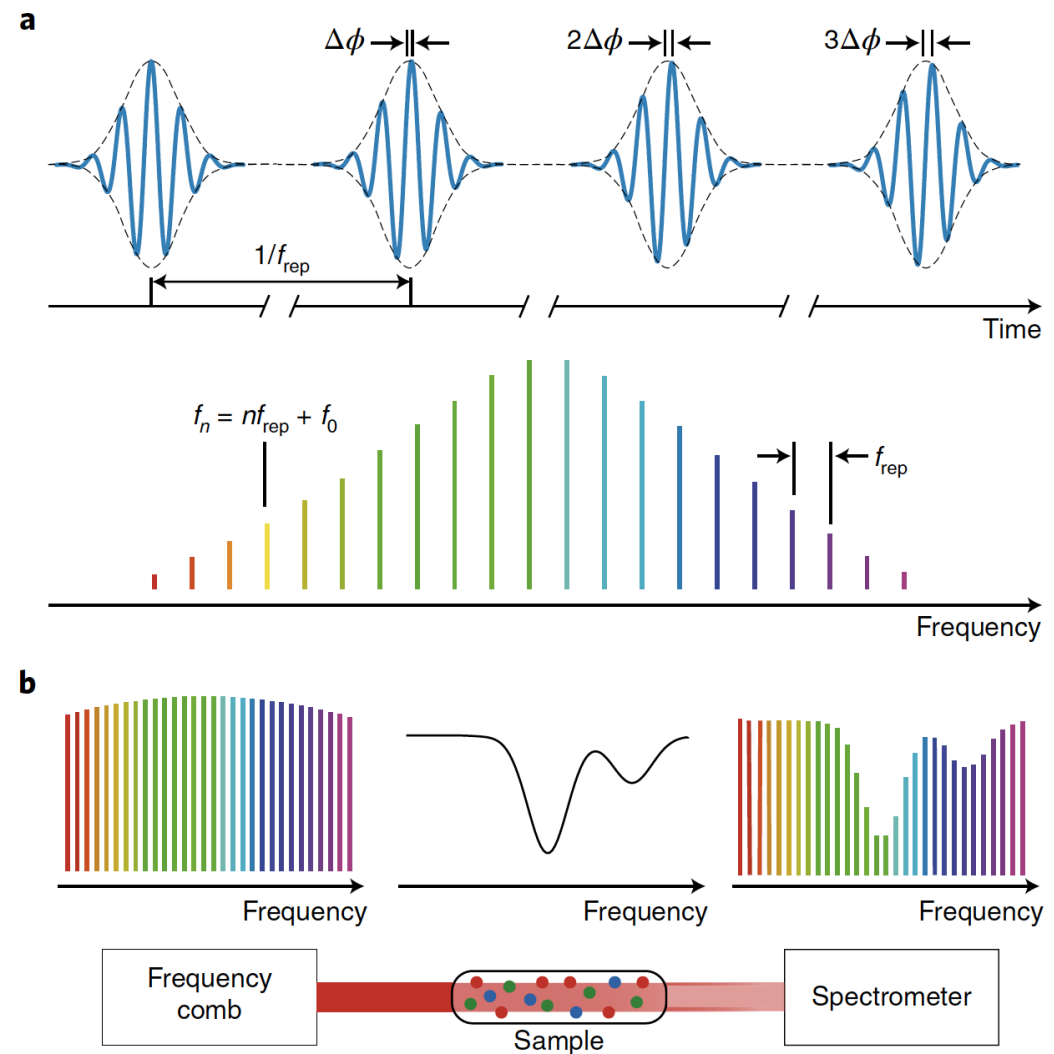


**Characters of ultrafast
laser pulses:
ultrahigh peak power**



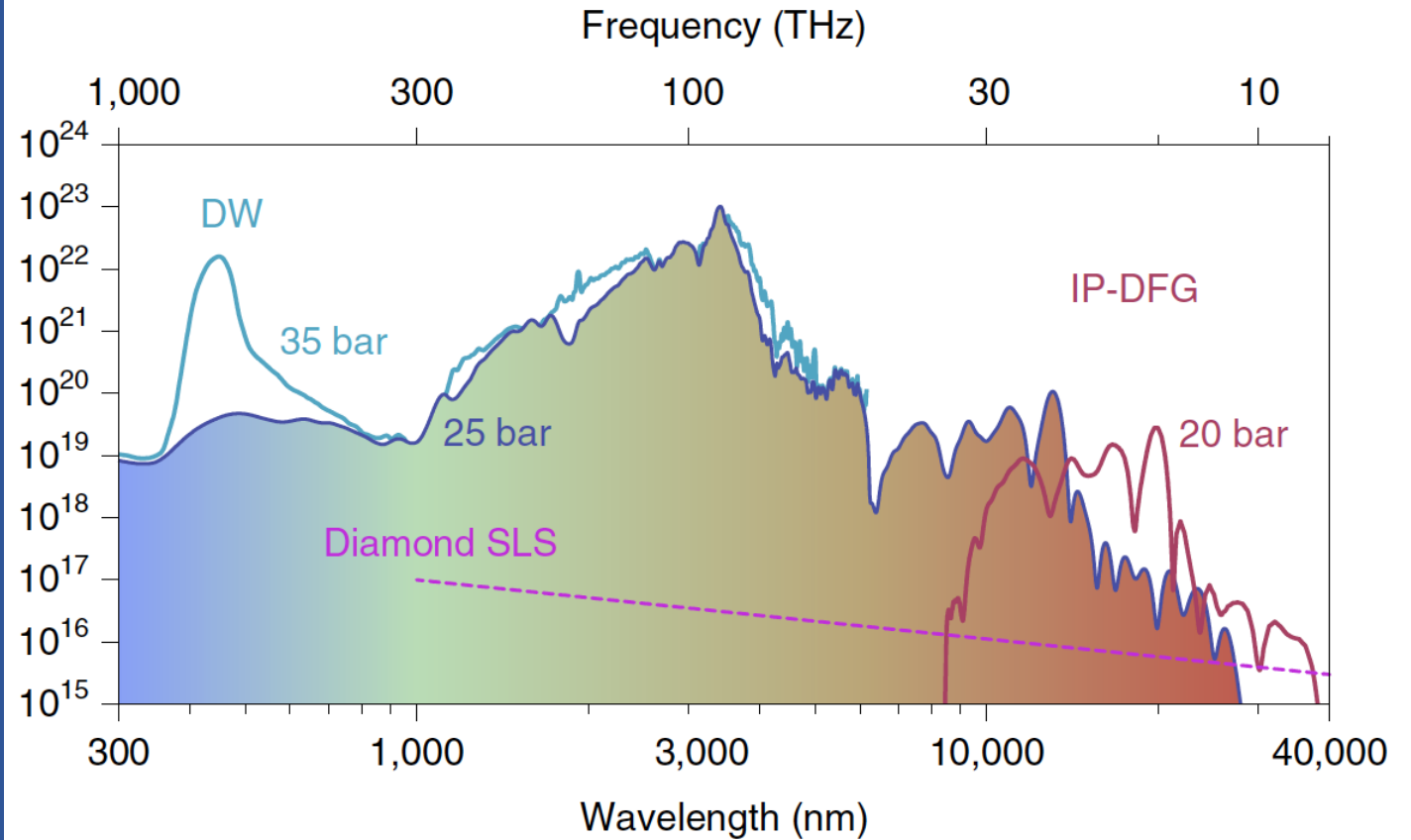
**SULF petawatt laser from SIOM: $10 \text{ PW} = 1 \times 10^{16} \text{ W}$,
(Generation capacity of the Three Gorges
hydroelectric station: $23 \text{ GW} = 2.3 \times 10^{10} \text{ W}$)**

Characters of ultrafast laser pulses: coherence



Nature Photonics **13**, 146 (2019)

Characters of ultrafast laser pulses: broad spectrum



Nature Photonics **15**, 277 (2021)

Femtochemistry: molecular “movie”



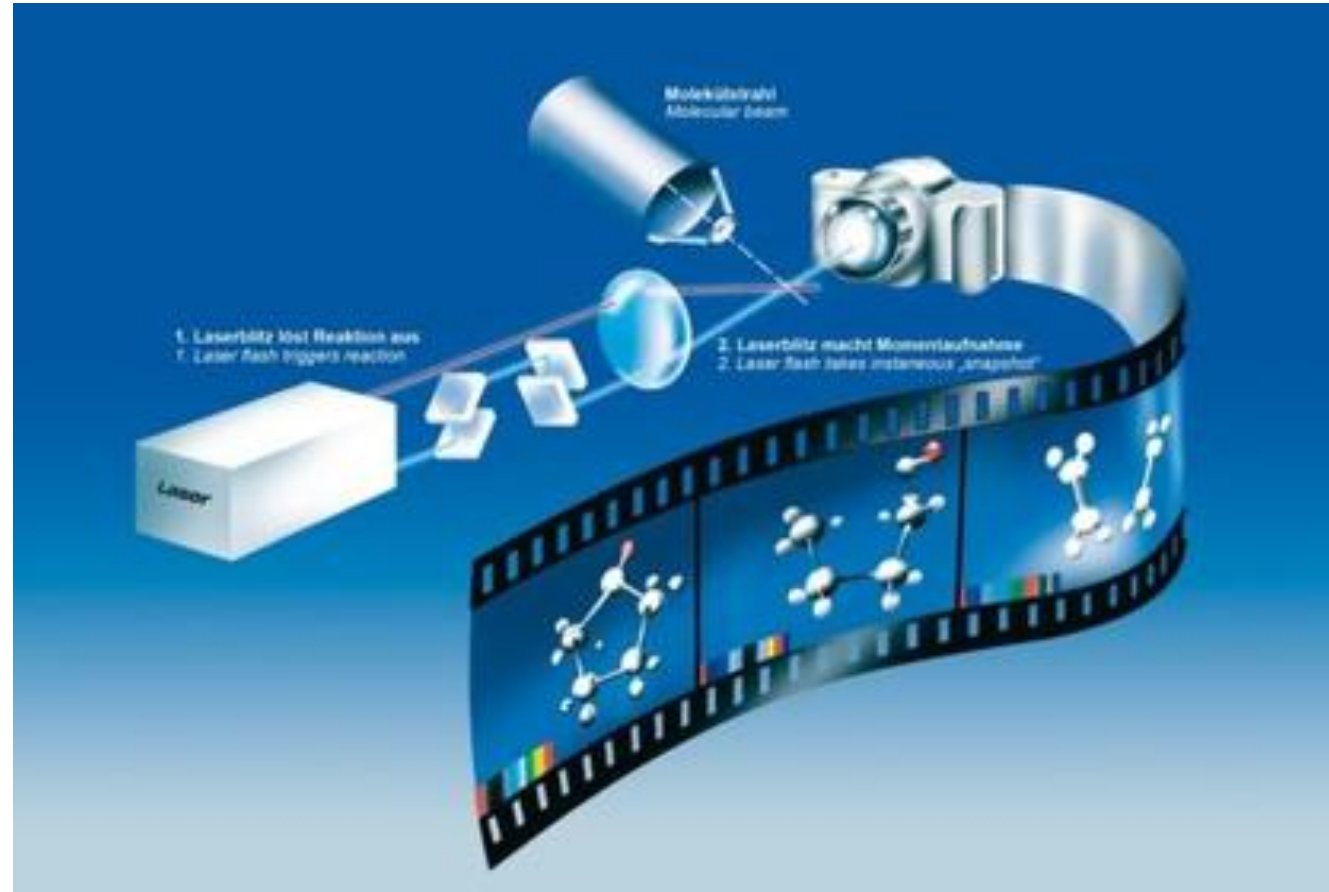
The Nobel Prize in Chemistry 1999

"for his studies of the transition states of chemical reactions using femtosecond spectroscopy"



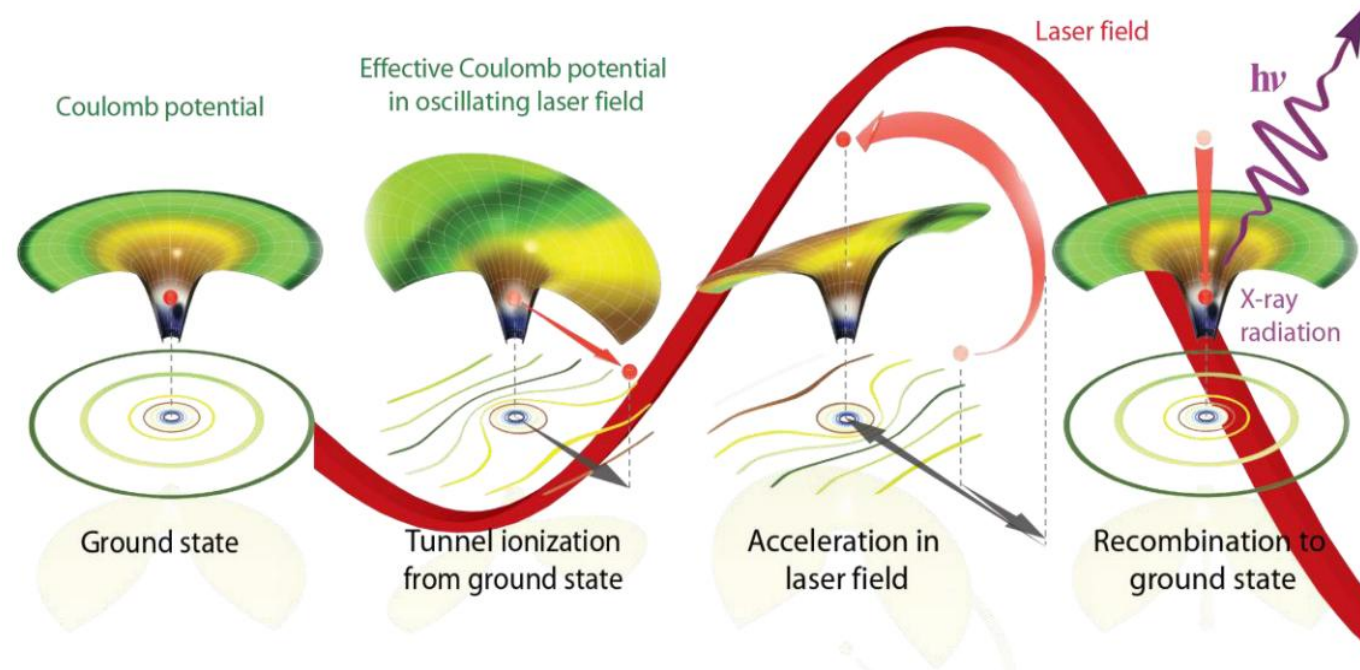
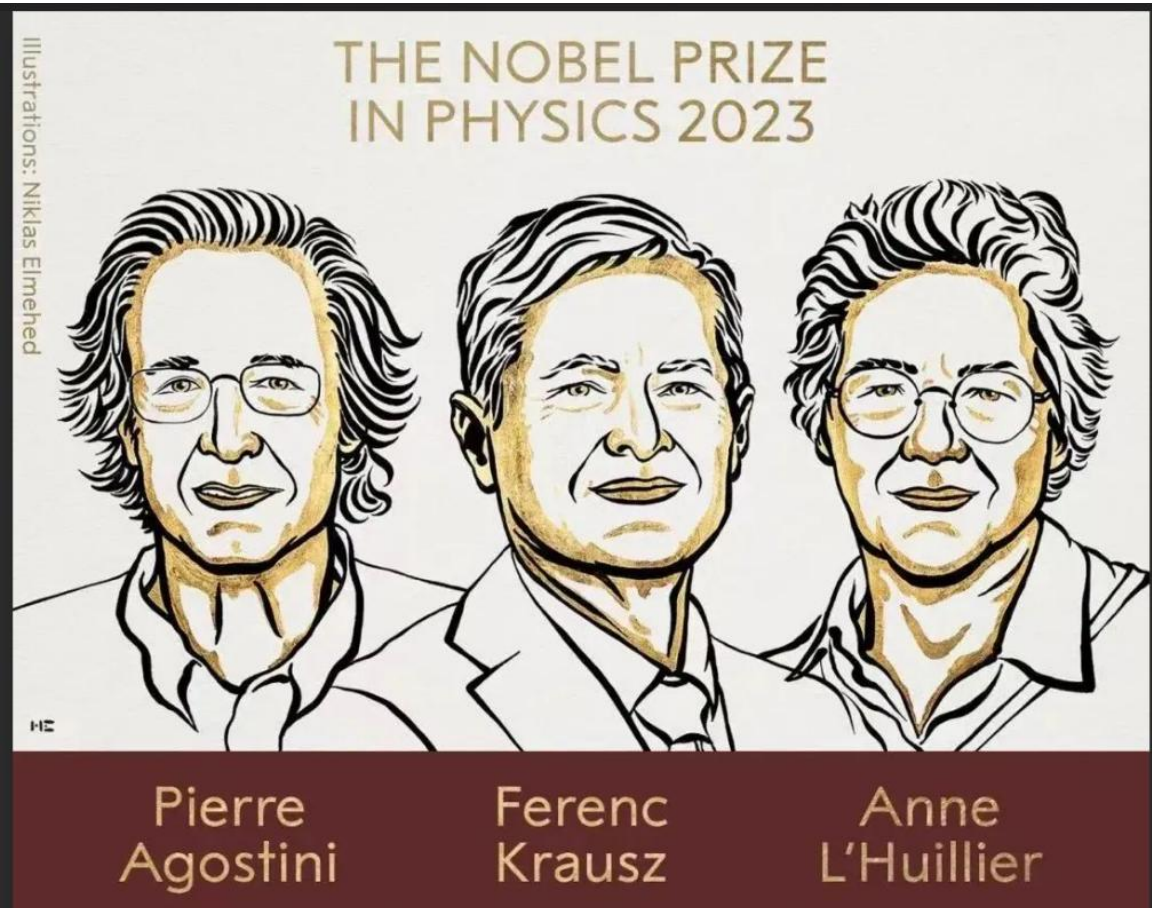
Ahmed H. Zewail

Egypt and USA



Exploiting the ultrashort pulse duration of fs laser pulses

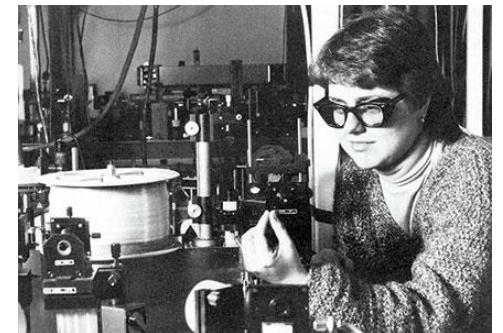
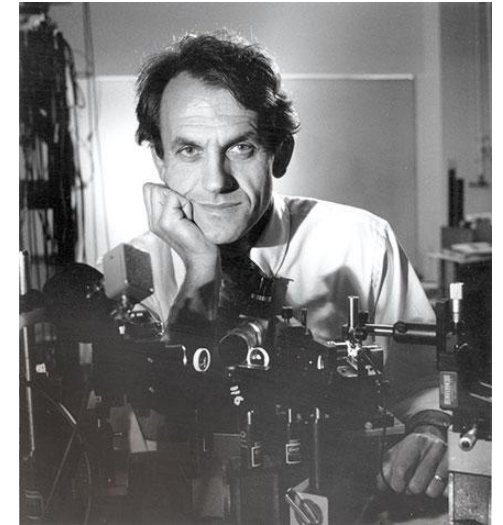
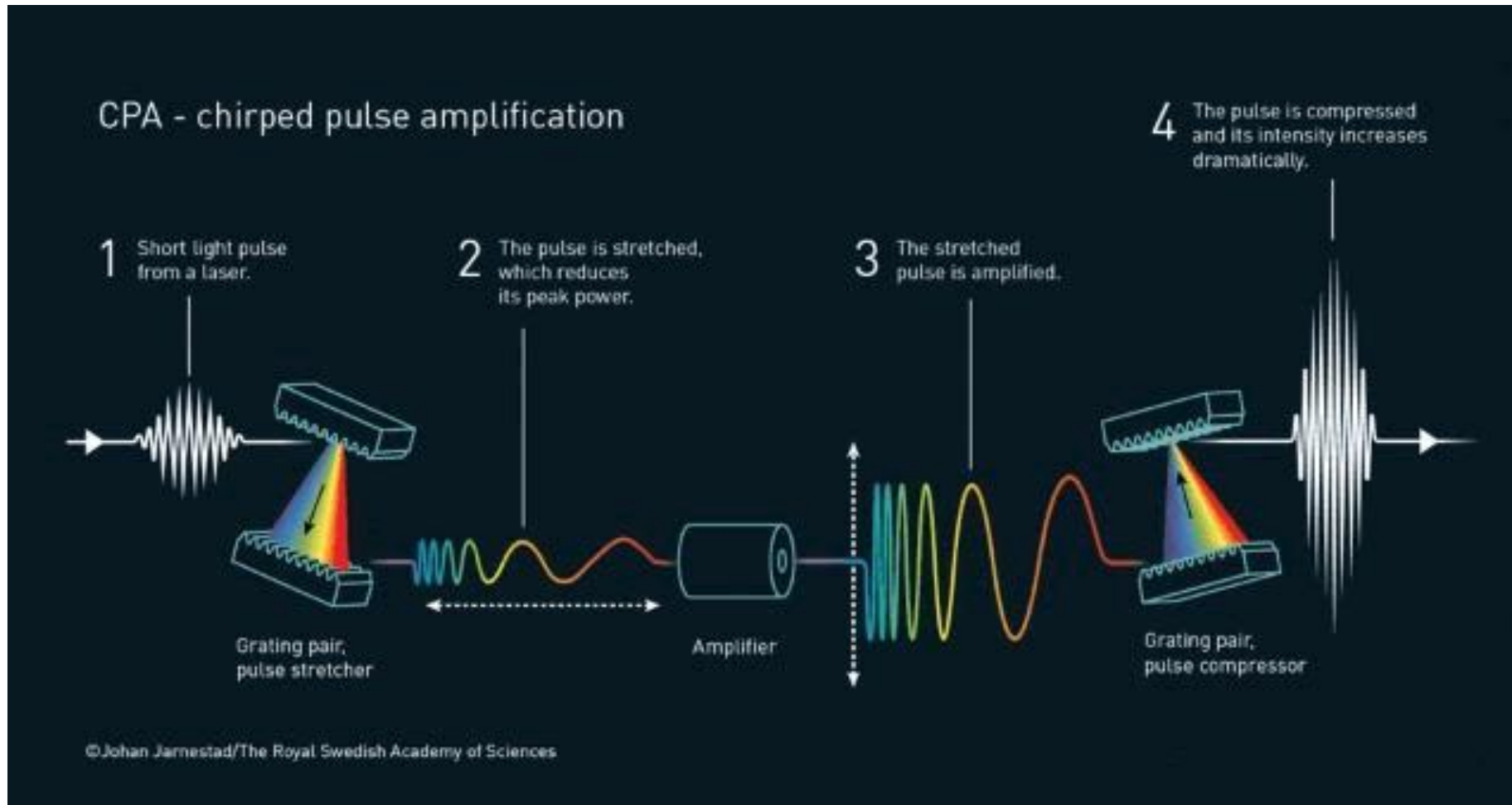
attosecond optics



Nat. Photonics **4**, 822 (2010)

the shortest optical pulse ever generated is only 43 attosecond

Pulse energy amplification technology: chirped pulse amplification (CPA)



Nobel prize in physics in 2018

Optical frequency comb



The Nobel Prize in Physics 2005

"for their contributions to the development of laser-based precision spectroscopy, including the optical frequency comb technique."

<https://www.nobelprize.org/prizes/physics/2005/summary/>



Photo: Sears.P.Studio

John L. Hall

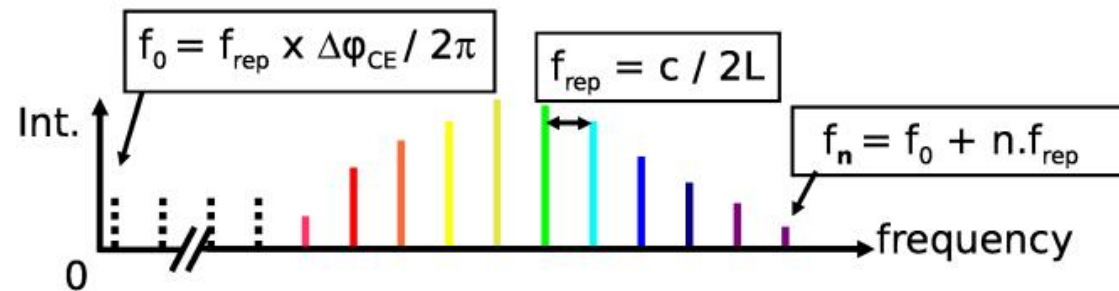
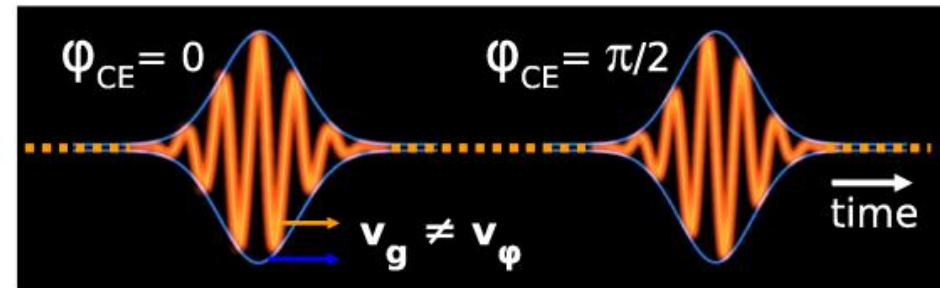
Prize share: 1/4



Photo: F.M. Schmidt

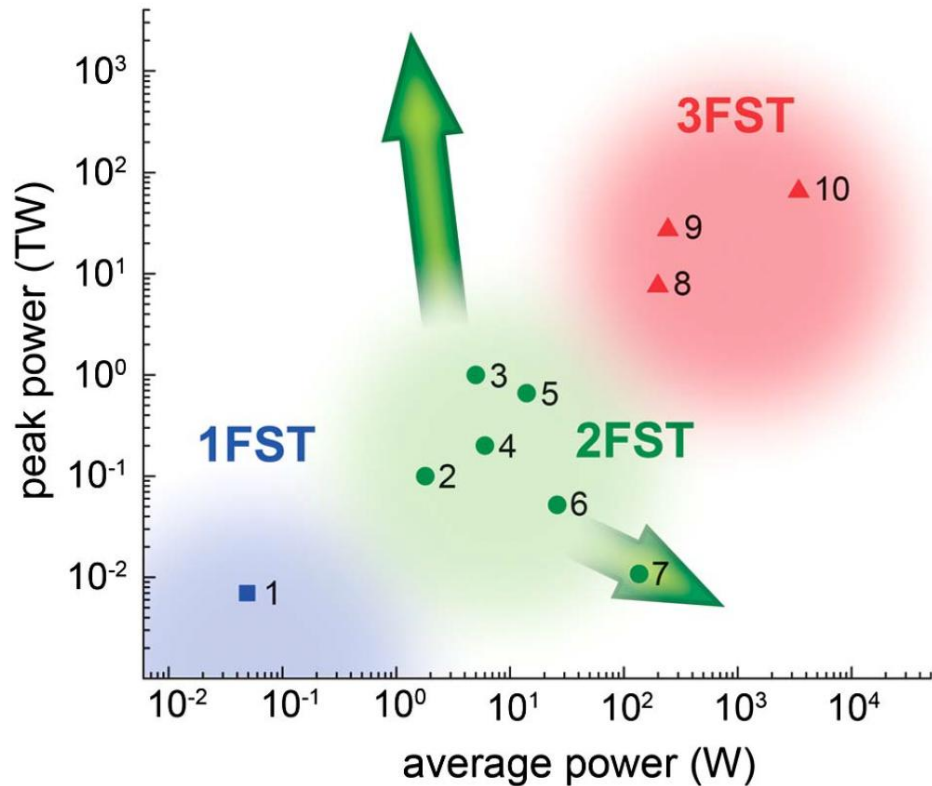
Theodor W. Hänsch

Prize share: 1/4



Exploiting the full coherence of fs laser pulses

Development of ultrafast laser technology



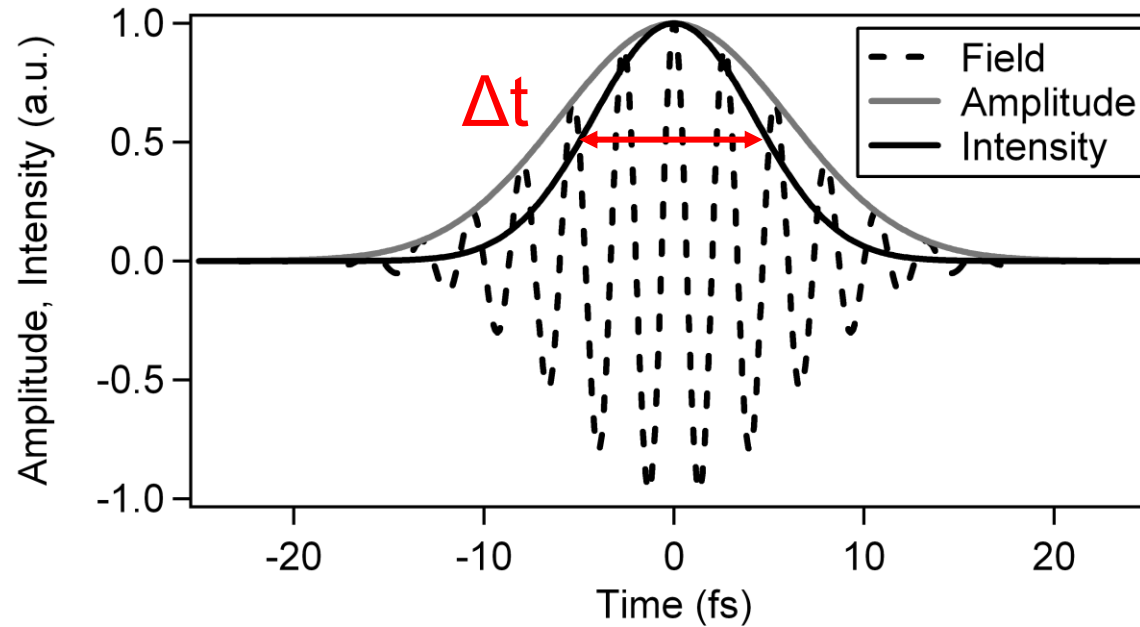
Optica 1, 45 (2014)

1 st	dye laser μJ pulse energy, MW peak power
2 nd	CPA based Ti:Sa laser, Yb-doped fiber, slab, disk lasers 10 W average power, PW peak power, or 100 W average power, GW peak power
3 rd	Yb-doped lasers + OPCPA or post compression KW average power, TW peak power

outline

- What is a laser?
- What is an ultrafast laser?
- **Mathematic description of an ultrafast laser pulse;**
- Ultrafast lasers for S³FEL;
- Femtosecond timing and synchronization

An ultrafast laser pulse: electric field and intensity distributions



$$E(t) = \sqrt{I(t)} \cos(\omega_0 t - \phi(t)) \quad \text{pulse energy} \propto \int |E(t)|^2 dt$$

$I(t)$: Intensity profile;

ω_0 : carrier angular frequency;

800 nm: $2\pi \times 3.7 \times 10^{14}$ Hz, $T = 2.67$ fs;

Δt : FWHM of intensity distribution;

$\phi(t)$: phase in time-domain;

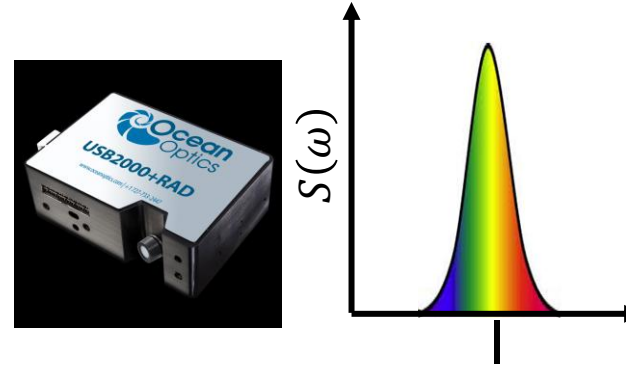
Electric field and spectrum amplitude: the Fourier Transform

Electric field

$$E(t) = \sqrt{I(t)} \cos(\omega_0 t - \phi(t))$$

Spectrum amplitude

$$E(\omega) = \sqrt{S(\omega)} \exp\{-i\phi(\omega)\}$$



- Both electric field and spectrum amplitude contains the complete information of the pulse. However, to think about ultrashort laser pulses, the Fourier Transform is essential: always consider things from both perspectives.

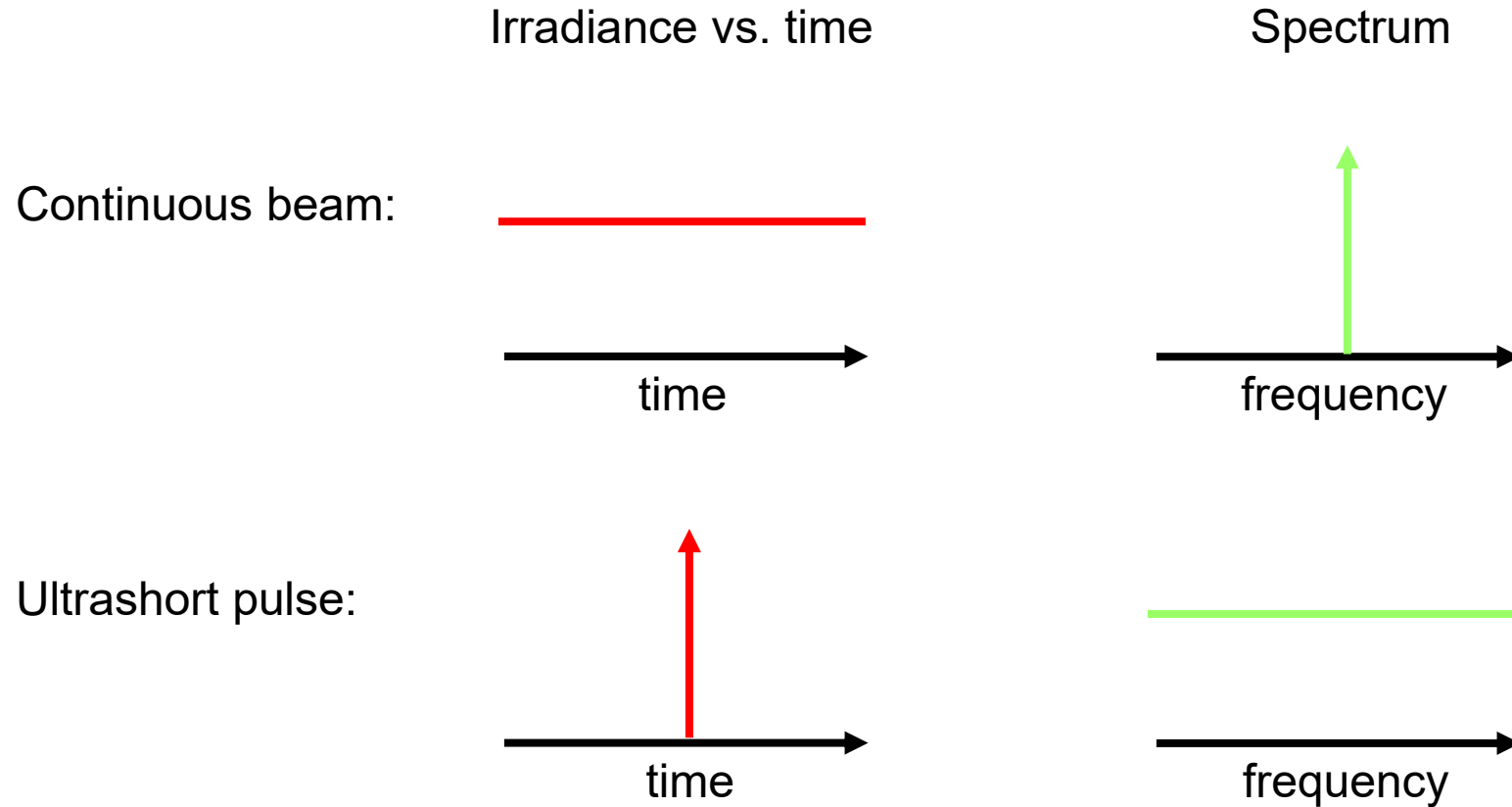
$$\widetilde{E}(\omega) = \int_{-\infty}^{\infty} E(t) \exp(-i\omega t) dt$$

$$E(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \widetilde{E}(\omega) \exp(i\omega t) d\omega$$

We always perform Fourier transforms on the real or complex pulse electric field.

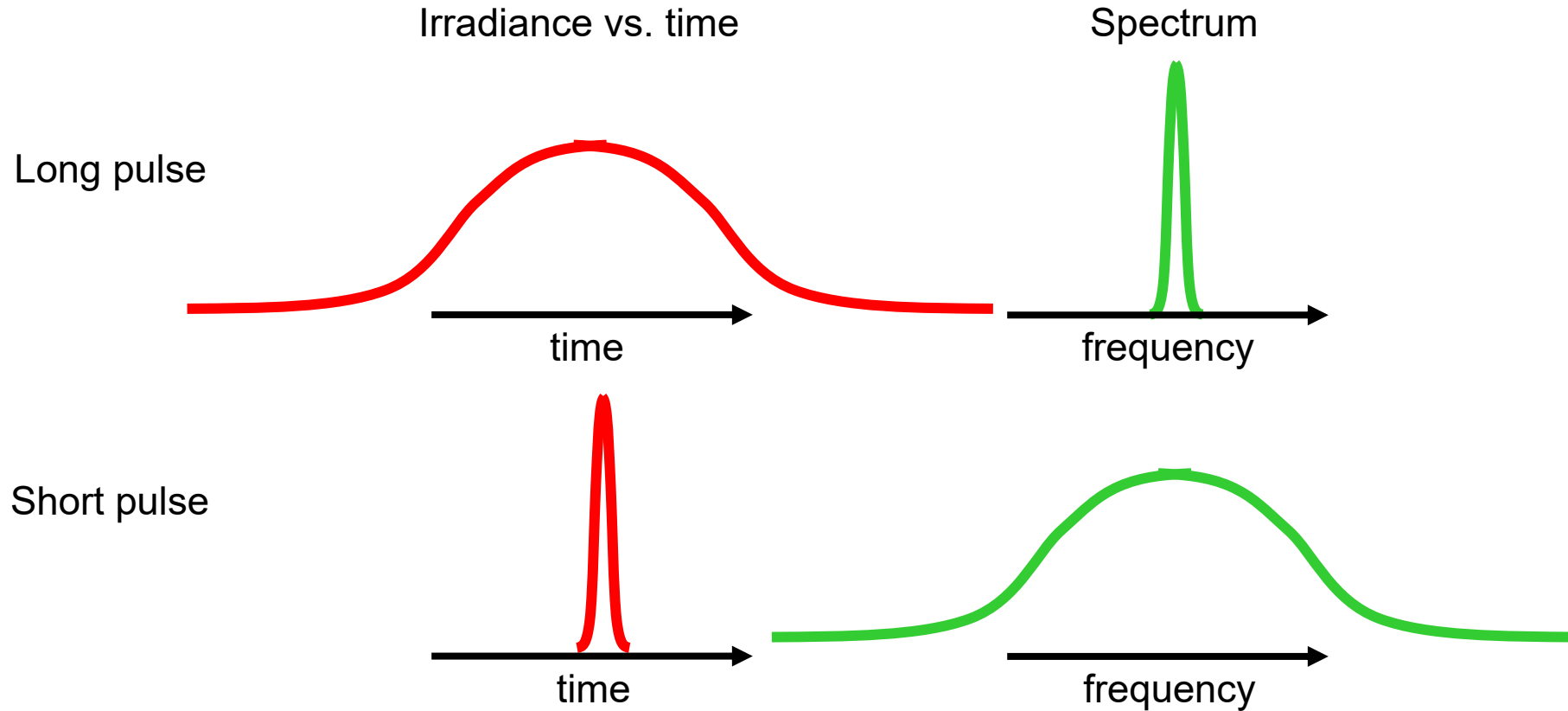
Continuous vs. ultrashort pulses of light

A constant and a delta-function are a Fourier-Transform pair.

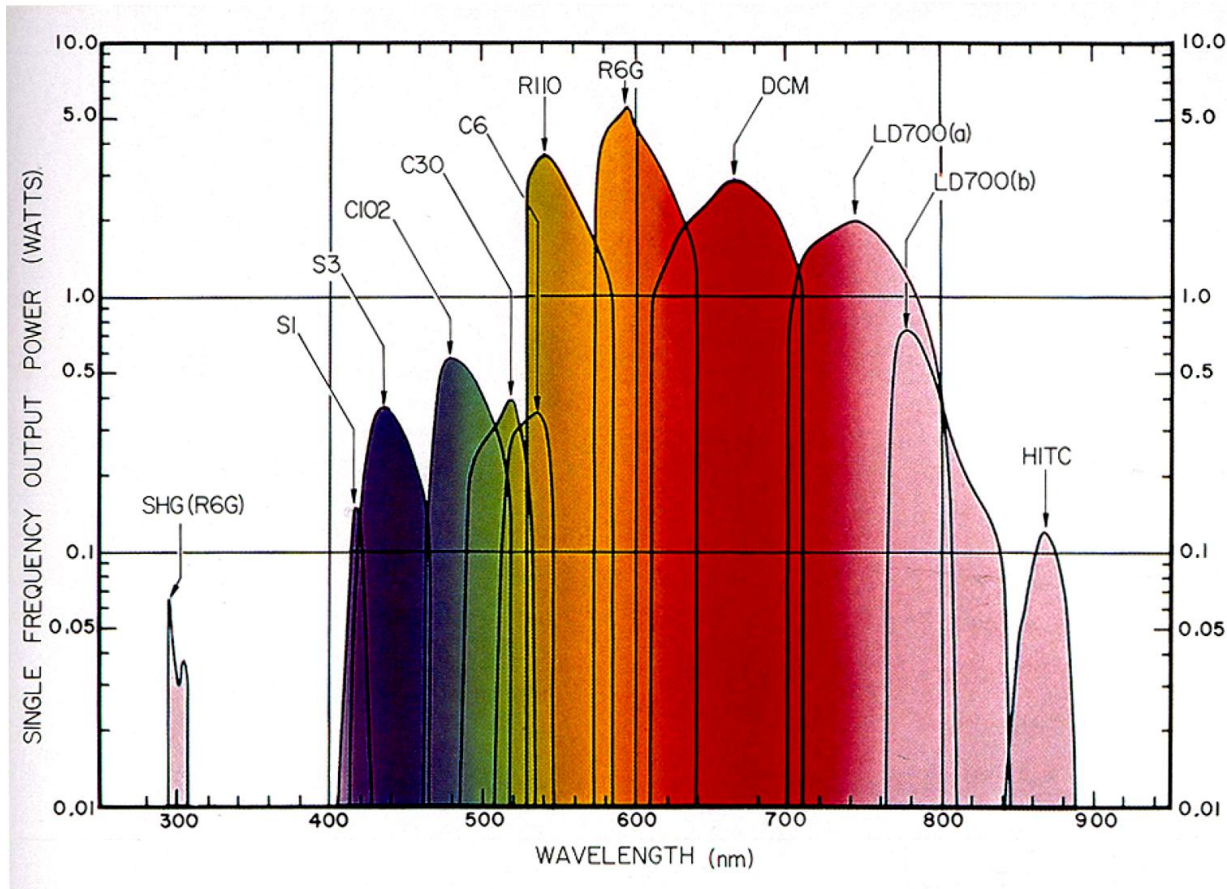


Long vs. short pulses of light

The uncertainty principle says that the product of the temporal and spectral pulse widths is greater than ~ 1 .



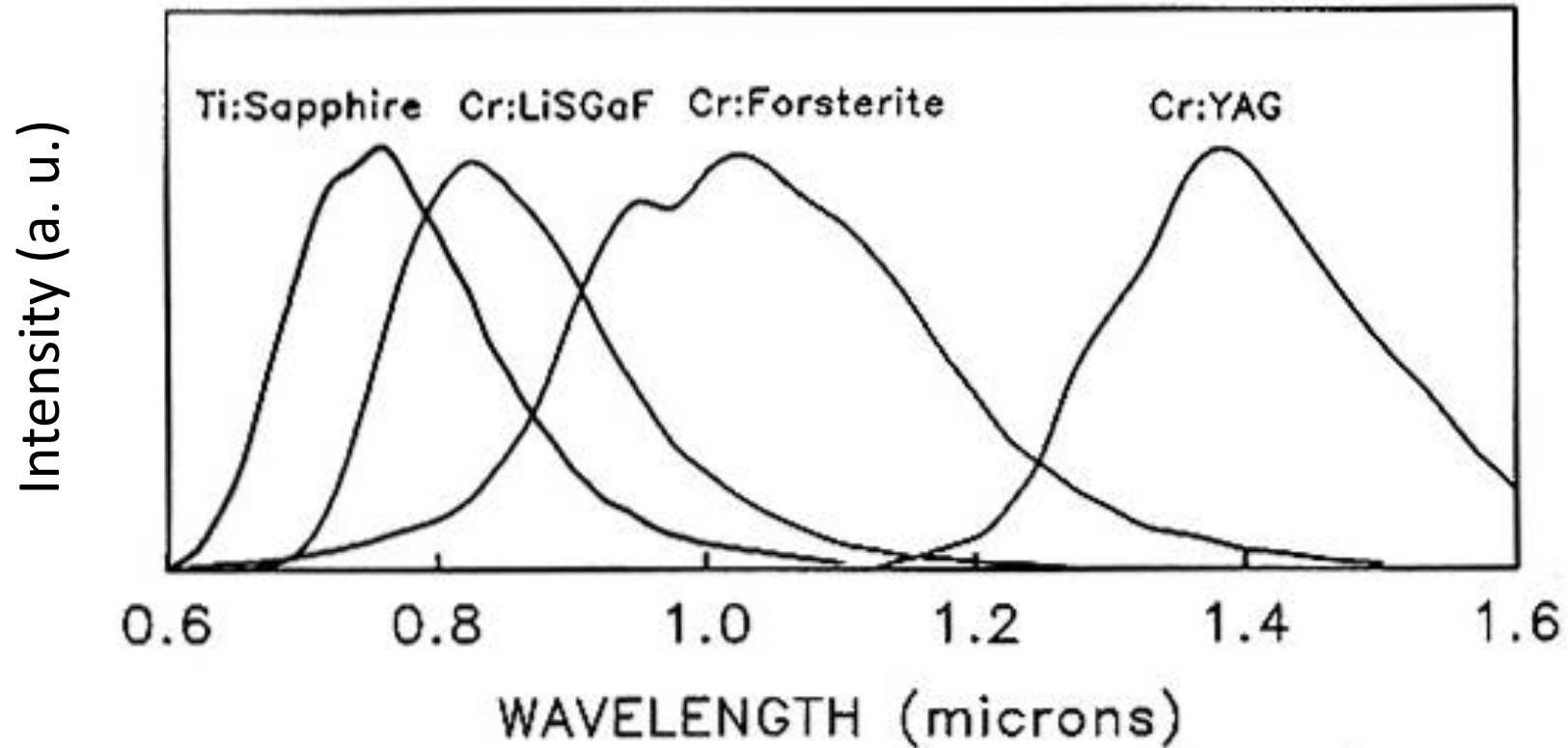
broadband gain media



For many years, dyes have been the broadband media that have generated ultrashort laser pulses.

broadband gain media

Solid-state laser media have broad bandwidths and are convenient.



A little bit more on phase

Electric field and phase

$$\mathcal{E}(t) = \frac{1}{2} \sqrt{I(t)} \cos(\omega_0 t - \phi(t))$$

Phase Taylor Series expansions:

$$\phi(t) = \phi_0 + \phi_1 \frac{t}{1!} + \cancel{\phi_2 \frac{t^2}{2!}} + \cancel{\phi_3 \frac{t^3}{3!}} + \dots$$

Spectrum and spectral phase

$$\mathcal{E}(\omega) = \sqrt{S(\omega)} \exp\{-i\varphi(\omega)\}$$

Phase Taylor Series expansions:

$$\varphi(\omega) = \varphi_0 + \varphi_1 \frac{\omega - \omega_0}{1!} + \varphi_2 \frac{(\omega - \omega_0)^2}{2!} + \varphi_3 \frac{(\omega - \omega_0)^3}{3!} + \dots$$

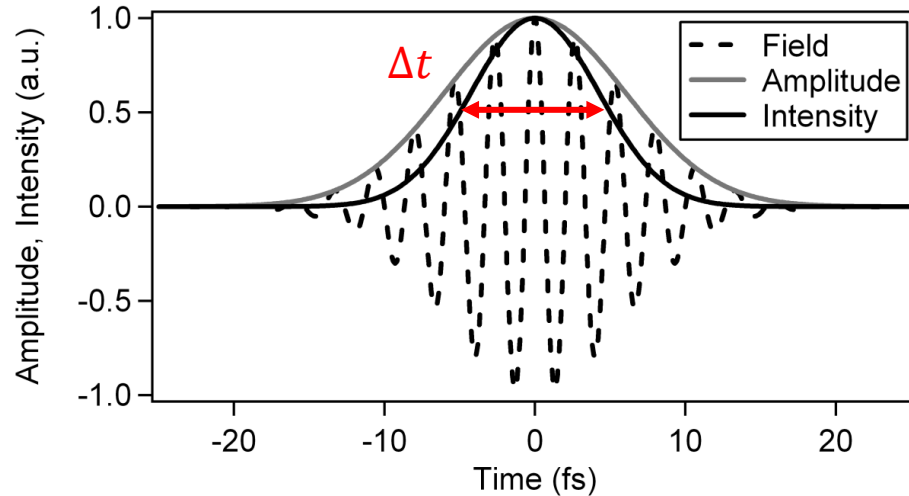
➤ If $\phi_n (n > 1) = 0$,

$$\mathcal{E}(t) = \frac{1}{2} \sqrt{I(t)} \cos(\omega_0 t - \phi(t)) = \frac{1}{2} \sqrt{I(t)} \cos\{(\omega_0 - \phi_1)t - \phi_0\},$$

electric field oscillates with a fixed frequency, this pulse is called transform limited, bandwidth-limited, or unchirped.

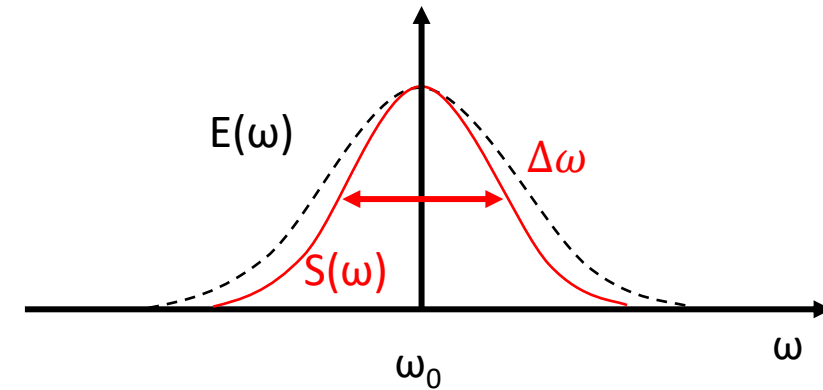
➤ If the second or higher orders derivative of phase are non-zero, oscillation frequency changes within the pulse. In analogy to bird sounds, this pulse is called a "chirped" pulse.

A transform limited Gaussian pulse and its spectrum



$$\mathcal{E}(t) = \frac{E_0}{2} e^{-2\ln 2 \frac{t^2}{\Delta t^2}} \cos(\omega_0 t - \phi_0)$$

Δt : FWHM of intensity distribution;
 ϕ_0 : carrier to envelop phase;
 ω_0 : carrier angular frequency;
 $\Delta \omega$: FWHM of the spectrum;



$$E(\omega) = \mathcal{F}\{E(t)\} = \frac{E_0 \Delta t}{2} \sqrt{\frac{\pi}{2 \ln 2}} e^{-\frac{\Delta t^2}{8 \ln 2} (\omega - \omega_0)^2}$$

$$\Delta t \frac{\Delta \omega}{2\pi} = 0.441$$

For a FT limited pulse of 100 fs, the FWHM of spectrum is 147 cm⁻¹.

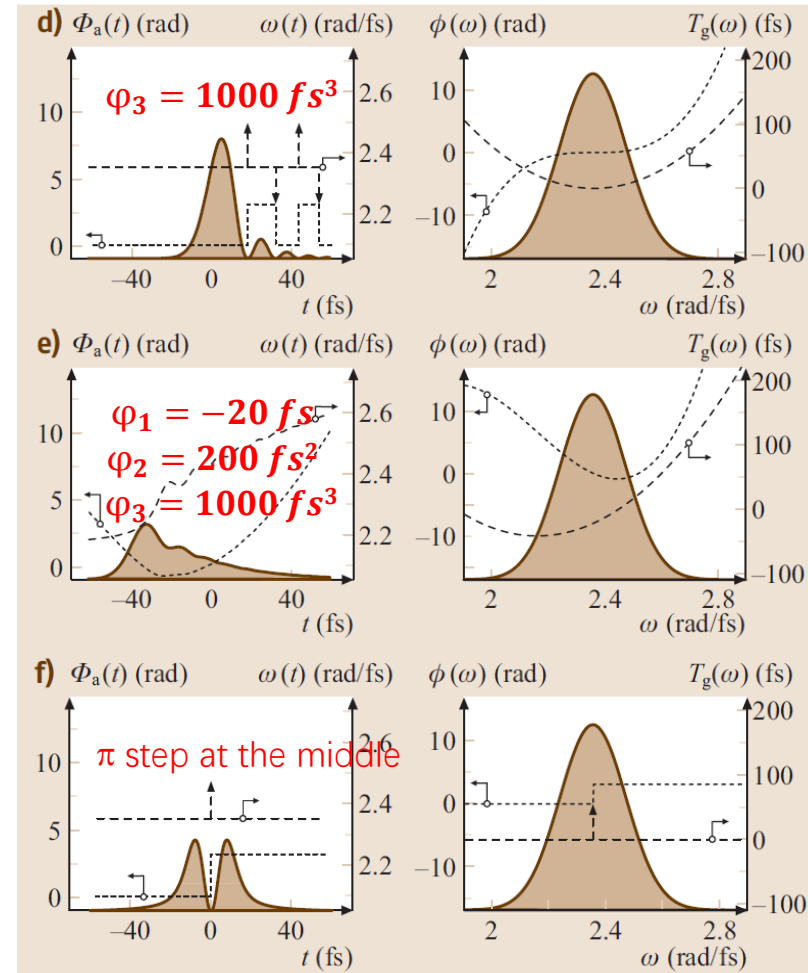
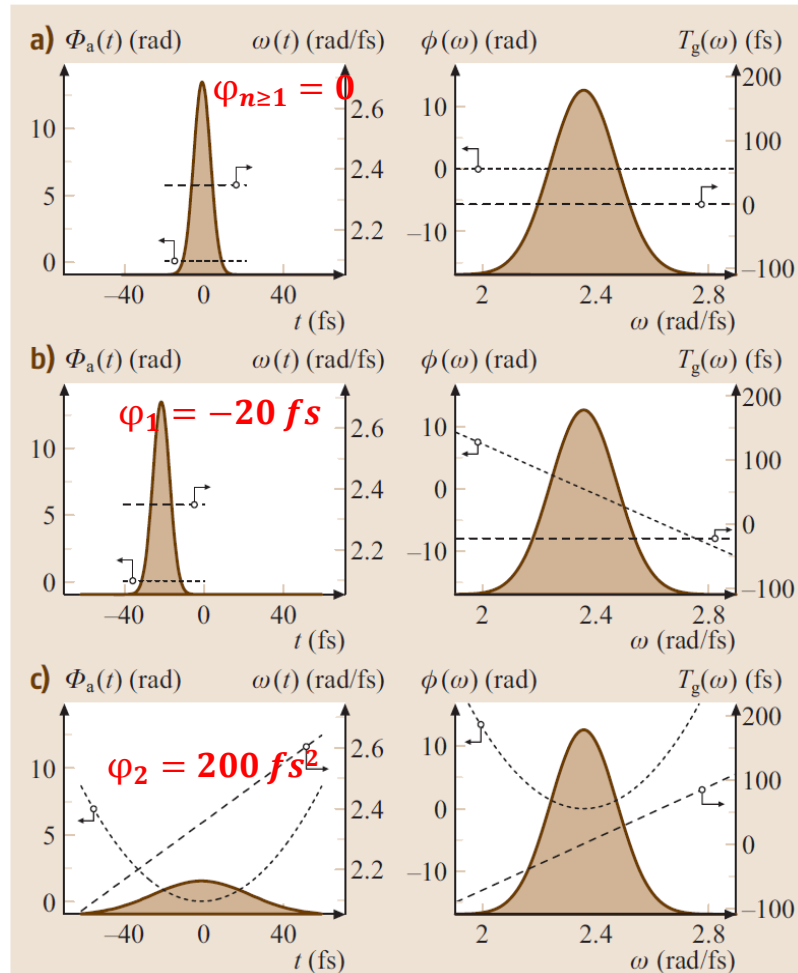
Temporal and spectral intensity profiles and time bandwidth products ($\Delta\nu\Delta t \geq K$) of various pulse shapes

Shape	$I(t)$	$I(\omega)$	$\Delta\nu\Delta t$	
Gaussian			0.441	$\mathcal{E}(t) = \frac{E_0}{2} e^{-2\ln 2 \frac{t^2}{\Delta t^2}}$ $\mathcal{E}(\omega) = \frac{E_0 \Delta t}{2} \sqrt{\frac{\pi}{2\ln 2}} e^{-\frac{\Delta t^2}{8\ln 2} (\omega - \omega_0)^2}$
Hyperbolic sechant			0.315	$\mathcal{E}(t) = \frac{E_0}{2} \operatorname{sech}[2\ln(1 + \sqrt{2}) \frac{t}{\Delta t}]$ $\mathcal{E}(\omega) = E_0 \Delta t \frac{\pi}{4\ln(1 + \sqrt{2})} \operatorname{sech}(\frac{\pi \Delta t}{4\ln(1 + \sqrt{2})} \omega)$
Square			0.886	$\mathcal{E}(t) = \frac{E_0}{2} t \in \left[-\frac{\Delta t}{2}, \frac{\Delta t}{2}\right]$ $\mathcal{E}(\omega) = \frac{E_0 \Delta t}{2} \operatorname{sinc}(\frac{\Delta t}{2} \omega)$
Single sided exponential			0.110	$\mathcal{E}(t) = \frac{E_0}{2} e^{-\frac{\ln 2}{2} \frac{t}{\Delta t}} t \in [0, \infty]$ $\mathcal{E}(\omega) = \frac{E_0 \Delta t}{2i\Delta t \omega + \ln 2}$
Symmetric exponential			0.142	$\mathcal{E}(t) = \frac{E_0}{2} e^{-\ln 2 \frac{t}{\Delta t}} t \in [0, \infty]$ $\mathcal{E}(\omega) = \frac{E_0 \Delta t \ln 2}{\Delta t^2 \omega^2 + (\ln 2)^2}$

Effects of the spectral phase on ultrashort pulses

$$\mathcal{E}(\omega) = \sqrt{S(\omega)} \exp\{-i\phi(\omega)\}$$

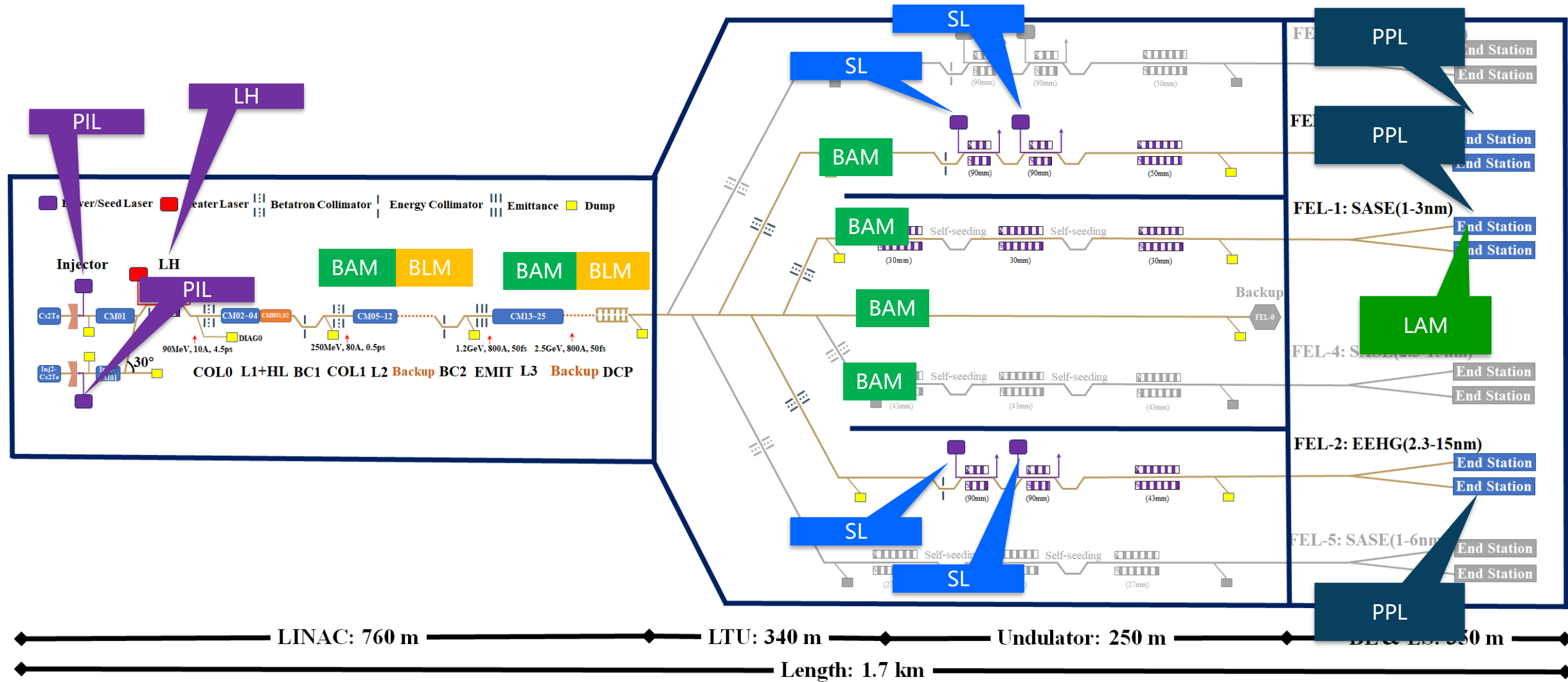
$$\phi(\omega) = \phi_0 + \phi_1 \frac{\omega - \omega_0}{1!} + \phi_2 \frac{(\omega - \omega_0)^2}{2!} + \phi_3 \frac{(\omega - \omega_0)^3}{3!} + \dots$$



outline

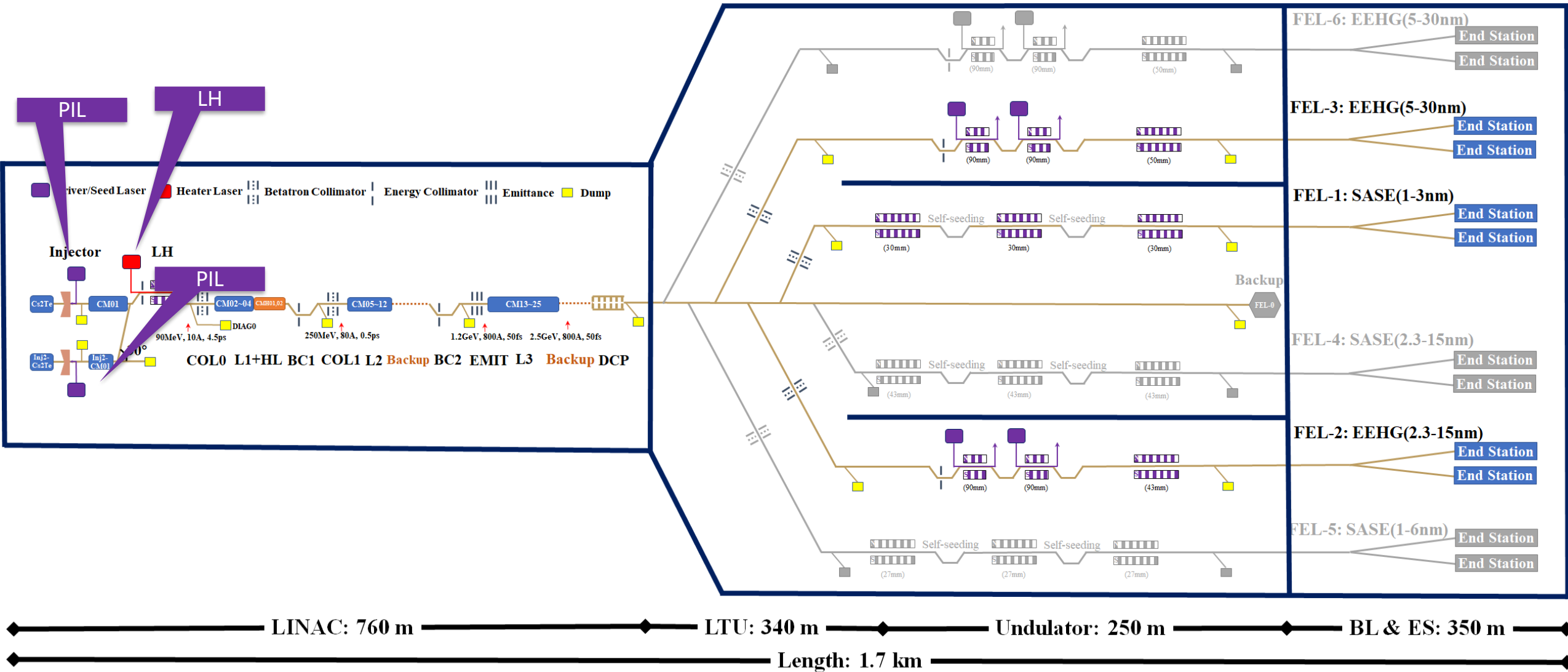
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Ultrafast laser systems for Shenzhen Superconducting Soft X-ray FEL (S³FEL)



photoinjector laser (PIL), laser heater (LH), seed laser (SL), pump-probe laser (PPL), beam length monitor (BLM), beam arriving-time monitor (BAM), X-ray arriving-time monitor (LAM), etc.

Photoinjector laser and laser heater



Photoinjector laser and laser heater

Photoinjector laser

Parameter	Value	Unit	Notes
repetition rate	1 – 1 M	Hz	——
Wavelength	257.5±2	nm	——
power	>200	mW	on photocathode
pulse energy stability	<1.5%	--	rms
pulse duration	20 – 60	ps	——
pulse shape	flat-top or expanded Gaussian	——	
beam size	0.2 – 2 (Hard-edge Dia.)	mm	nominal 1.6
beam shape	flat-top		
pointing stability	<10	um	rms

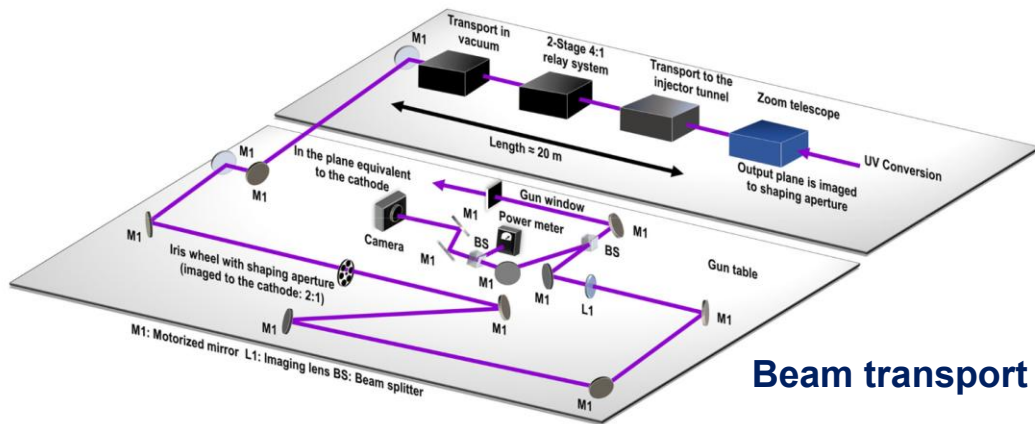
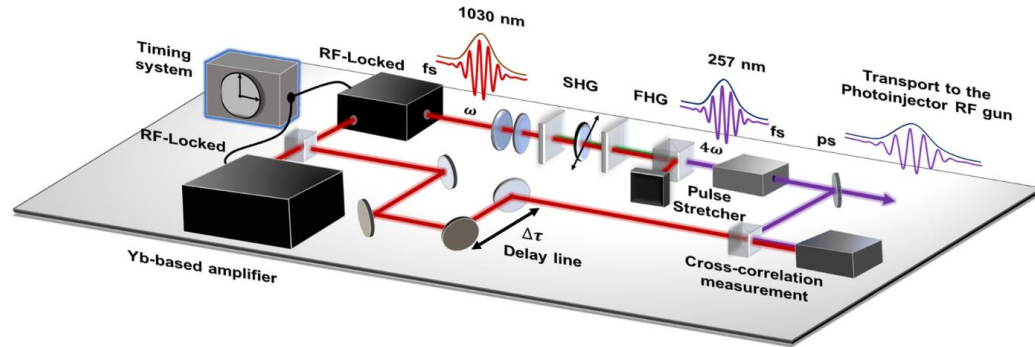
Laser heater

Parameter	Value	Unit
Repetition rate	1 – 1 M	Hz
Wavelength	1030	nm
Pulse energy	Up to 15	μJ
Pulse duration	10 – 30	ps
Beam size	0.2 – 0.8 (FWHM)	mm

PIL and LH share large similarities.

PIL at high rep. rate FEL facilities

LCLS II

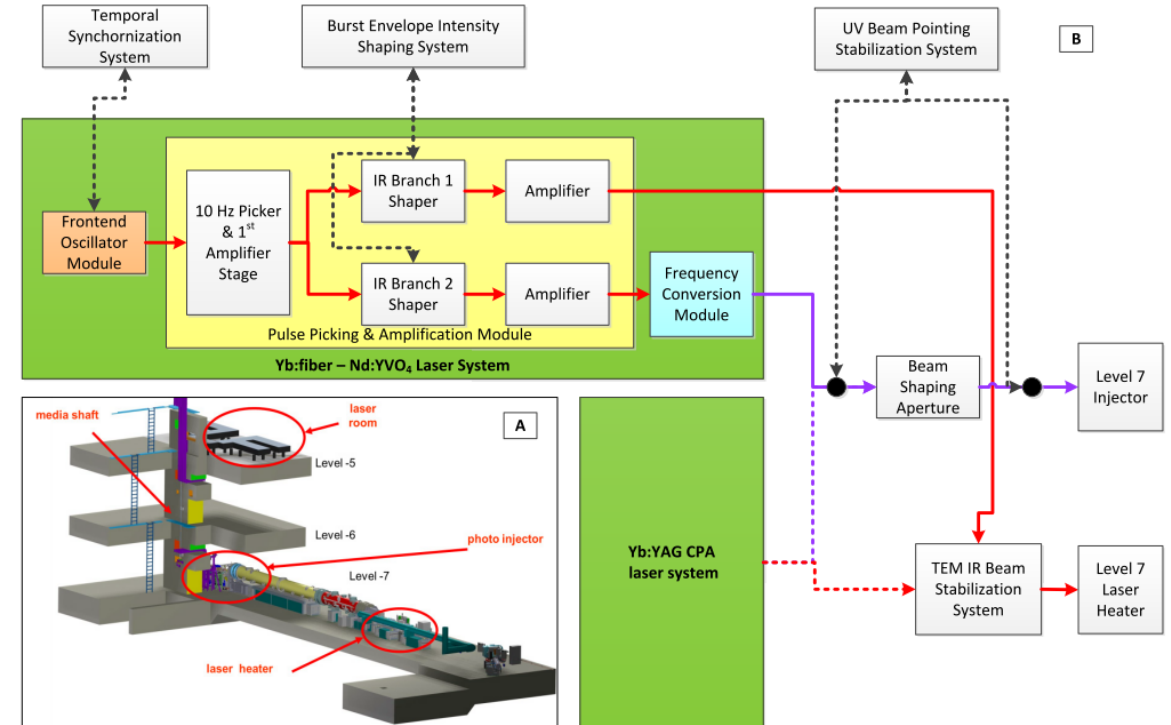


Beam transport

LCLS-II photoinjector laser based on Yb fiber laser

arXiv preprint:2307.12030 (2023).

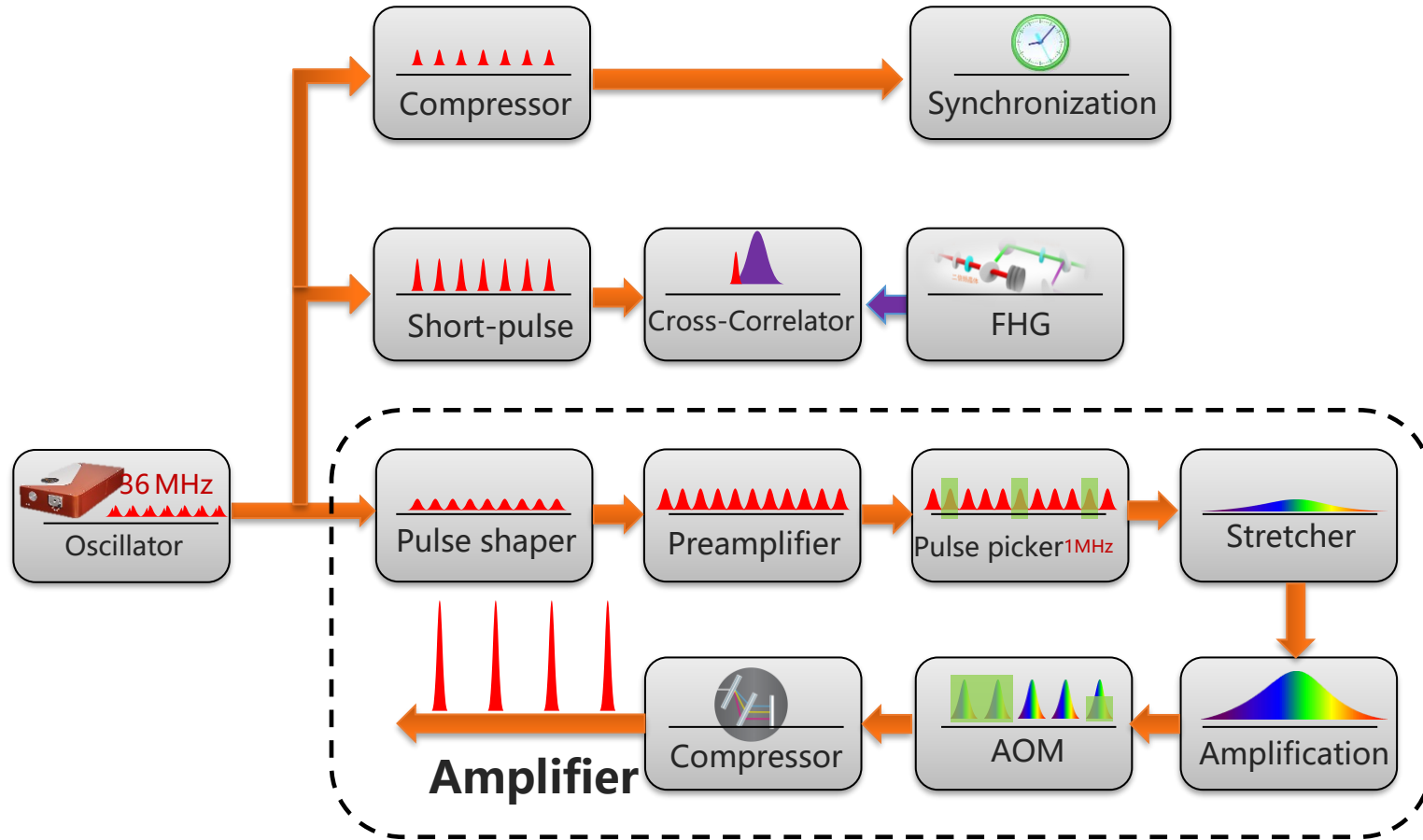
European XFEL



Two laser systems, Nd:YVO₄ and Yb:YAG, acting as “hot swap” backups for each other

Proceedings of the International Free-Electron Laser Conference (FEL'19), Hamburg, Germany. 2019.

50 μ J/50 W is sufficient for PIL & LH



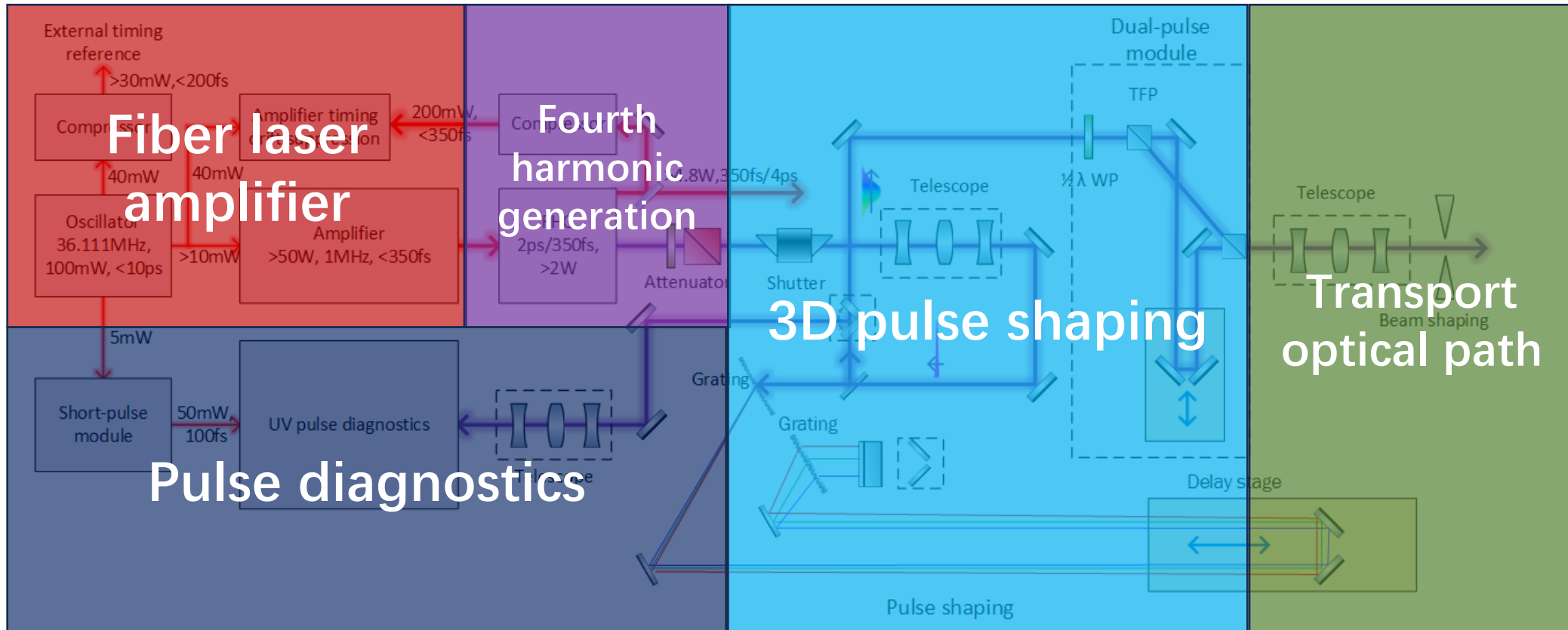
Commercial available ultrafast laser

- **Laser synchronization**
- **Slow drift compensation**
- **Temporal/spatial pulse shaping**
- **Improved pulse energy stability**
- **Stable transportation**
- ...

- **Multiple industry-grade choices**
 - **But a lot of customized upgrades need to be done**
- 

Photoinjector laser system for S³FEL

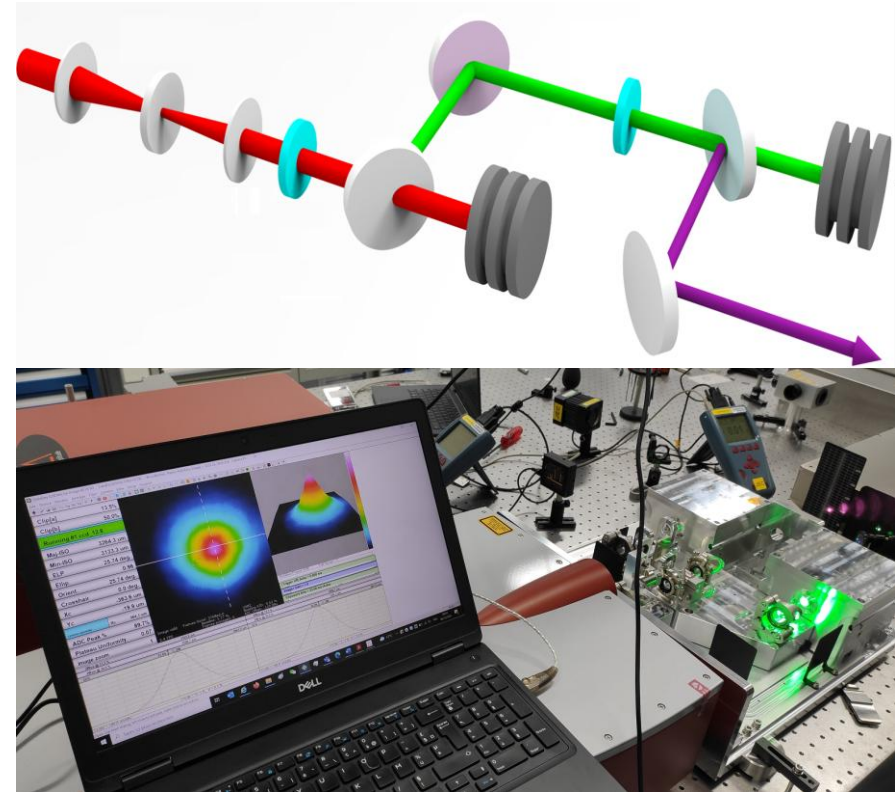
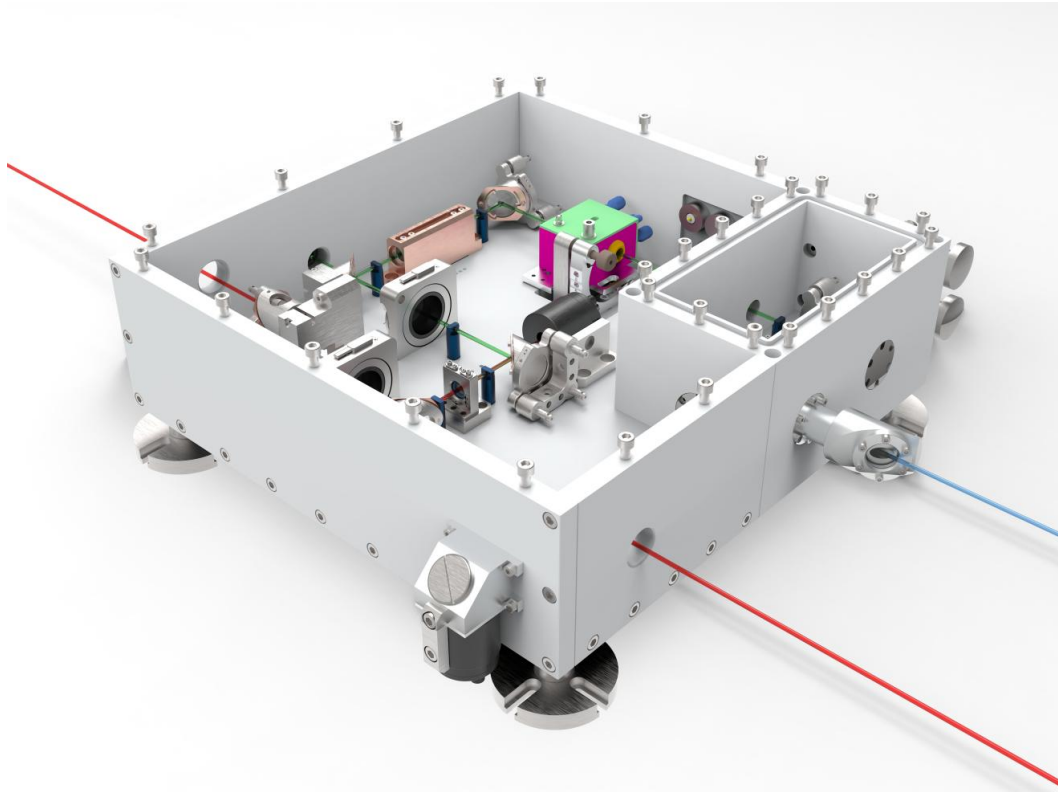
- High rep. rate(1MHz): matching FEL CW operating mode
- High power(>2W): generating sufficient electron beam charge
- 3D shaping: reducing emittance of electron beams, accelerating the saturation process of FEL



Fourth harmonic generation

Converts infrared fundamental to DUV pulses (257.5 nm)

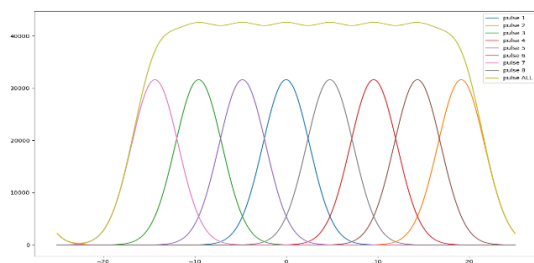
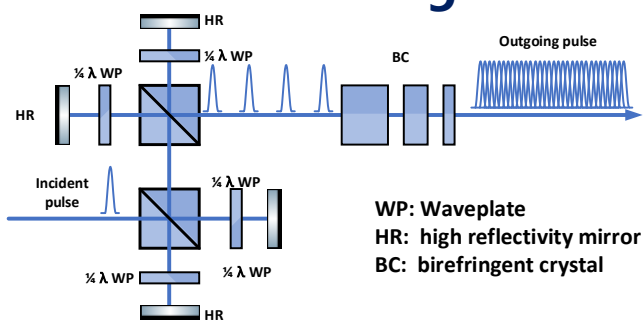
- 257.5 nm ultraviolet photon effectively extracts electrons from the Cs_2Te photocathode
- Nonlinear conversion efficiency continuously adjustable
- High-quality beam provides a good foundation for beam shaping



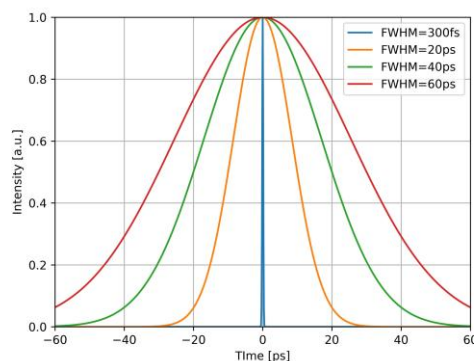
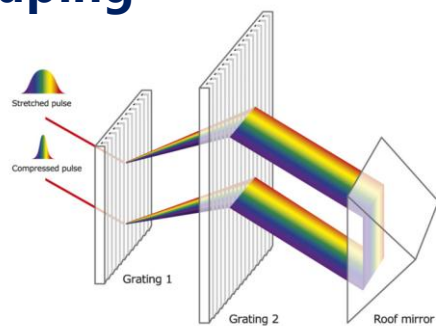
Sophisticated 3D pulse shaping

Perform longitudinal and transverse shaping of DUV laser pulses to improve the quality of electron bunches.

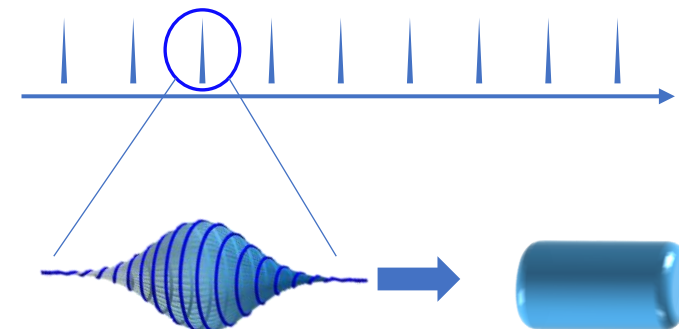
Longitudinal Shaping



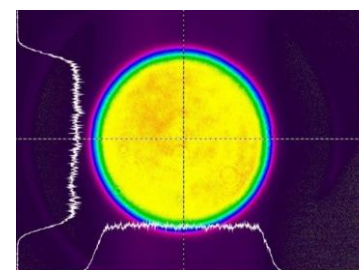
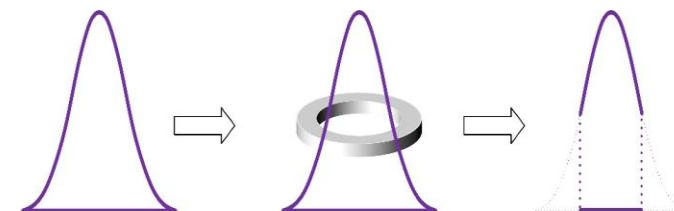
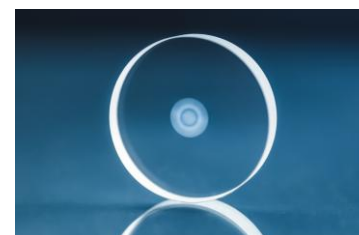
Pulse stacking



40 Grating pair



Transverse Shaping



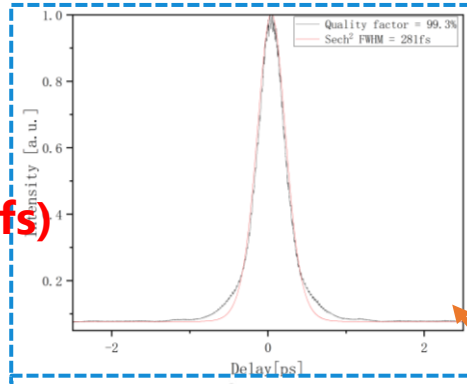
Flat-Top converter



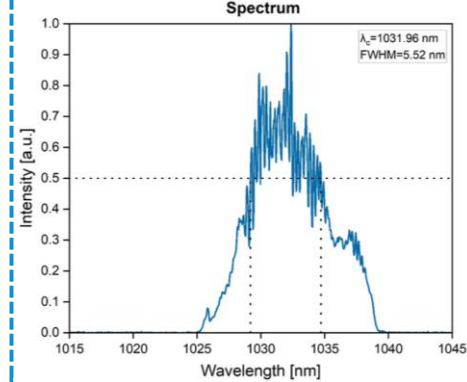
Aperture truncation

PIL for S³FEL

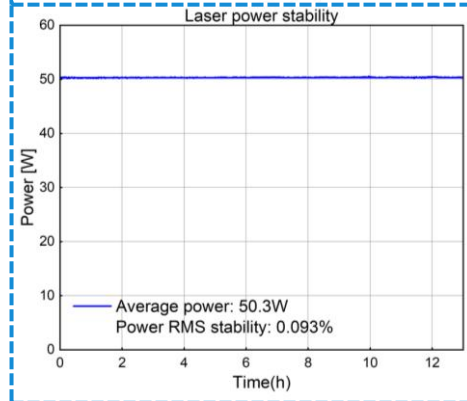
IR pulse duration (280 fs)



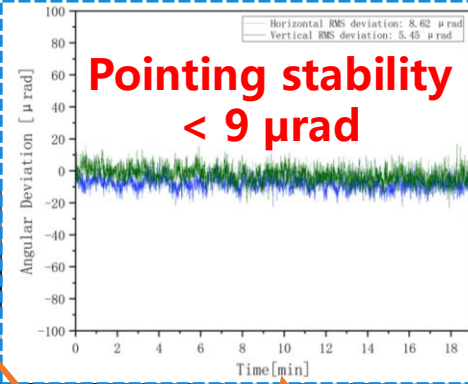
IR spectrum



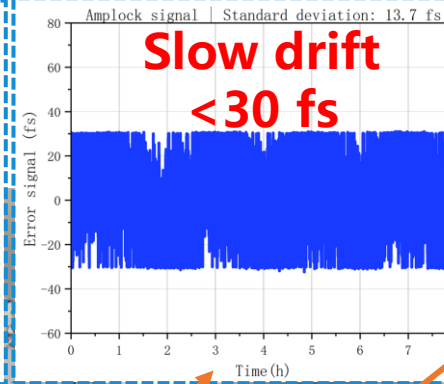
Long term power stability



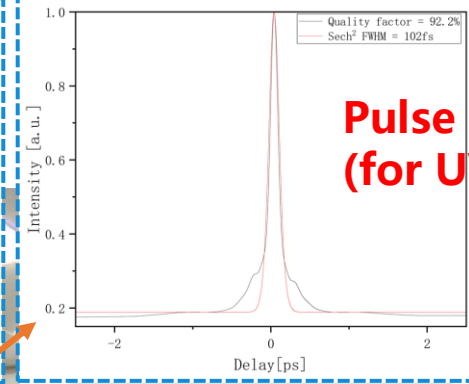
Pointing stability
< 9 μ rad



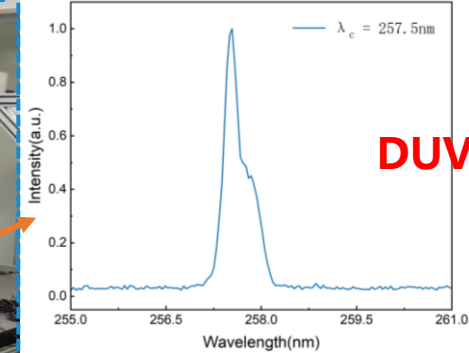
Slow drift
< 30 fs



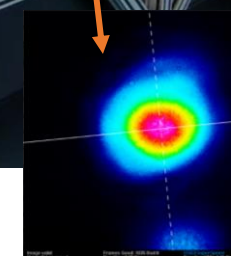
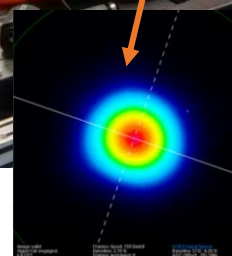
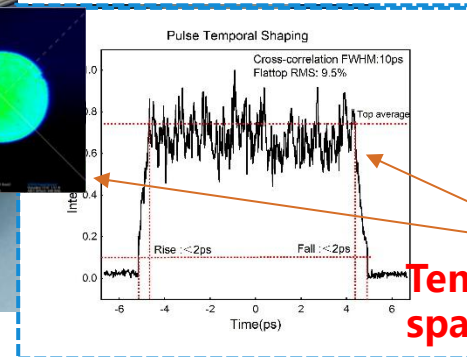
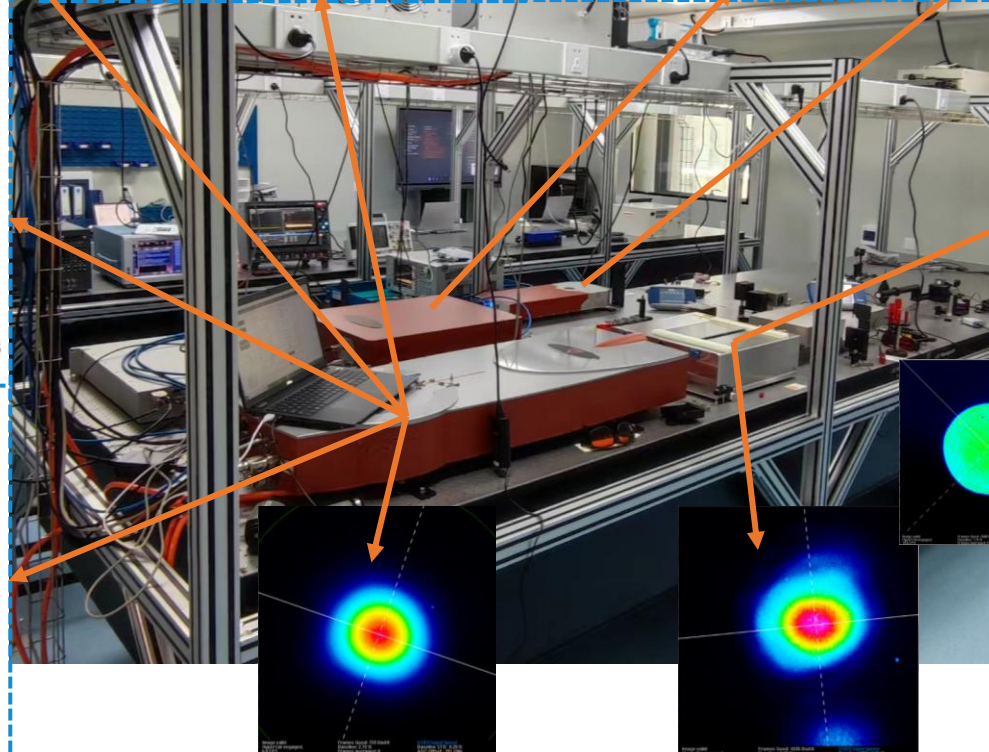
Pulse shortening
(for UV diagnosis)



DUV spectrum



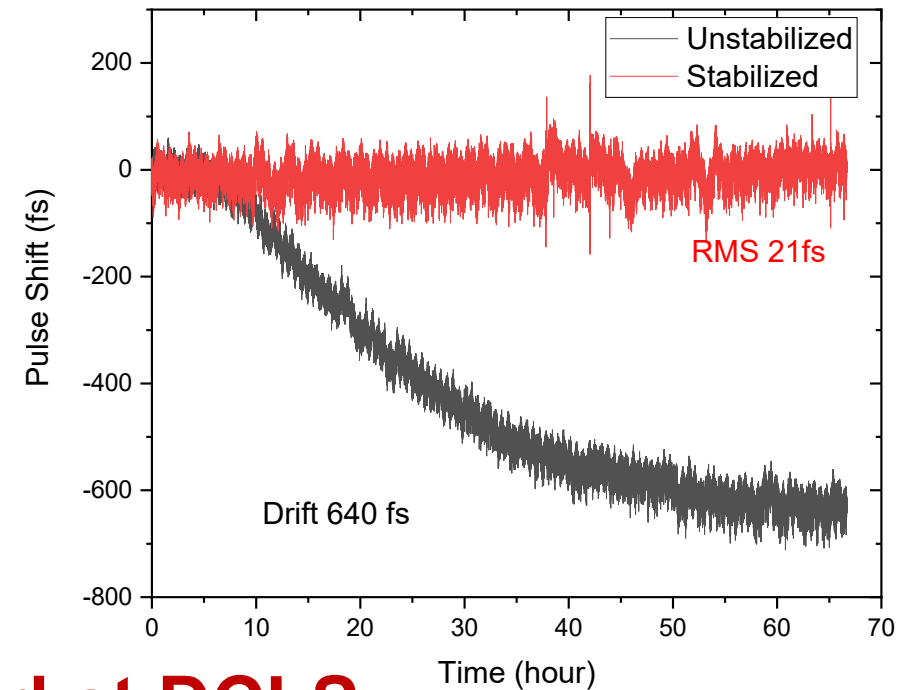
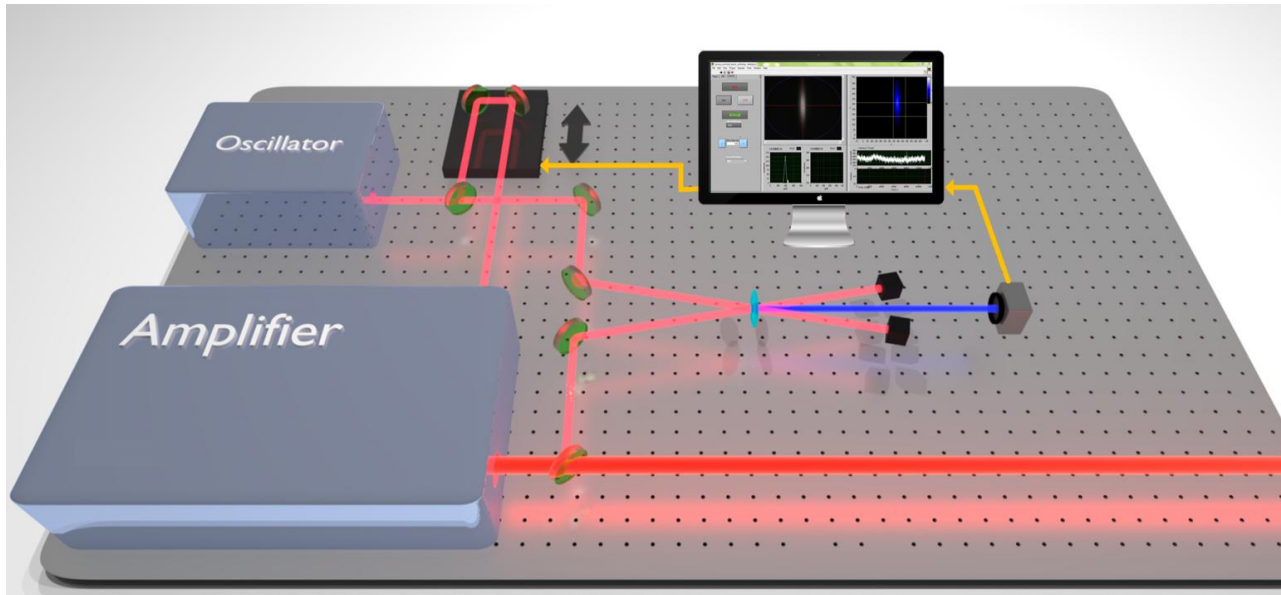
Temporal and spatial shaping



Amplifier timing drift suppression

- Suppress the long-term timing drift

- Single-shot cross-correlation technique for measuring temporal jitter and drift;
- Butterworth filter to remove high-frequency noise components;
- PID feedback loop to eliminate long-term timing drift introduced by amplifiers

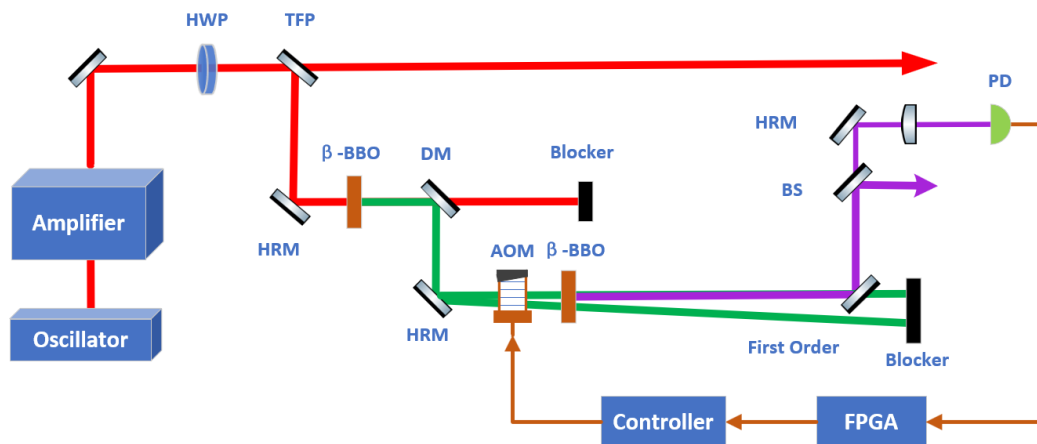


It has been successfully implemented at DCLS.

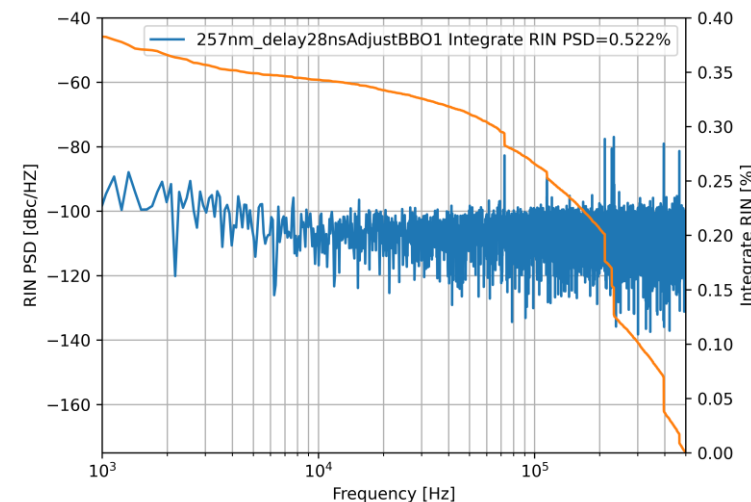
Pulse energy stabilization: from ~2% to 0.5%

Pulse energy stability of PIL is critical to many aspects of S³FEL. We have put lots of efforts to improve it, utilizing both active feedback and passive control techniques.

Active and passive noise suppression



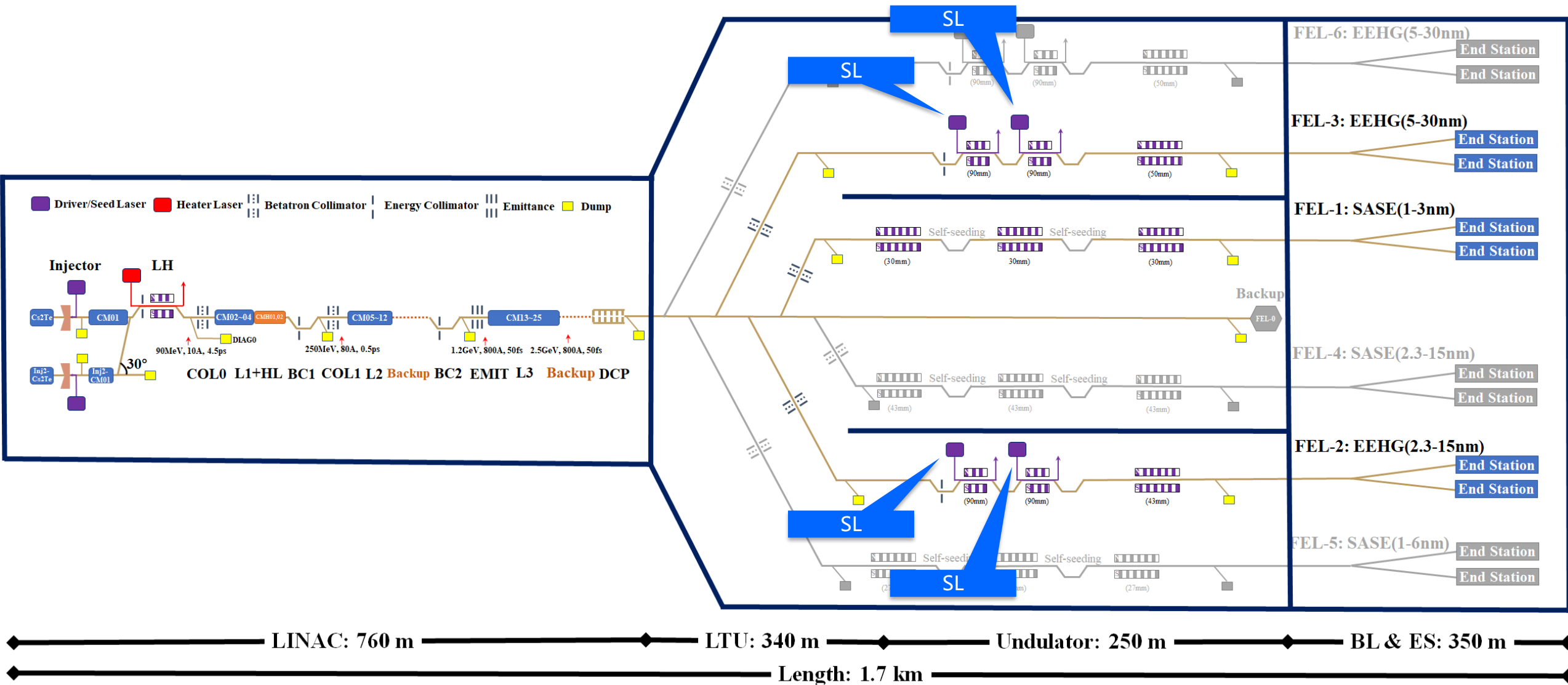
Relative intensity noise of UV laser



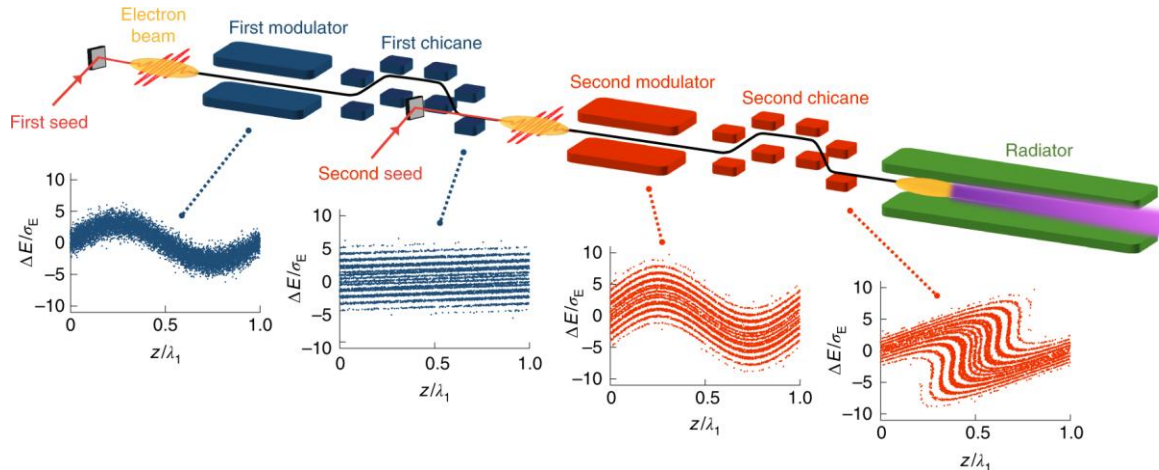
- Optimizing the Amplifier system parameters
- Optimizing the nonlinear frequency conversion process
- Active feedback bandwidth of 500 kHz;
- Combining feedforward active noise canceling (ANC) with the PID algorithm to maximize control effectiveness

First experimental result of ~ 0.5% rms was just demonstrated.

Seed lasers



EEHG seed lasers requirements for S³FEL

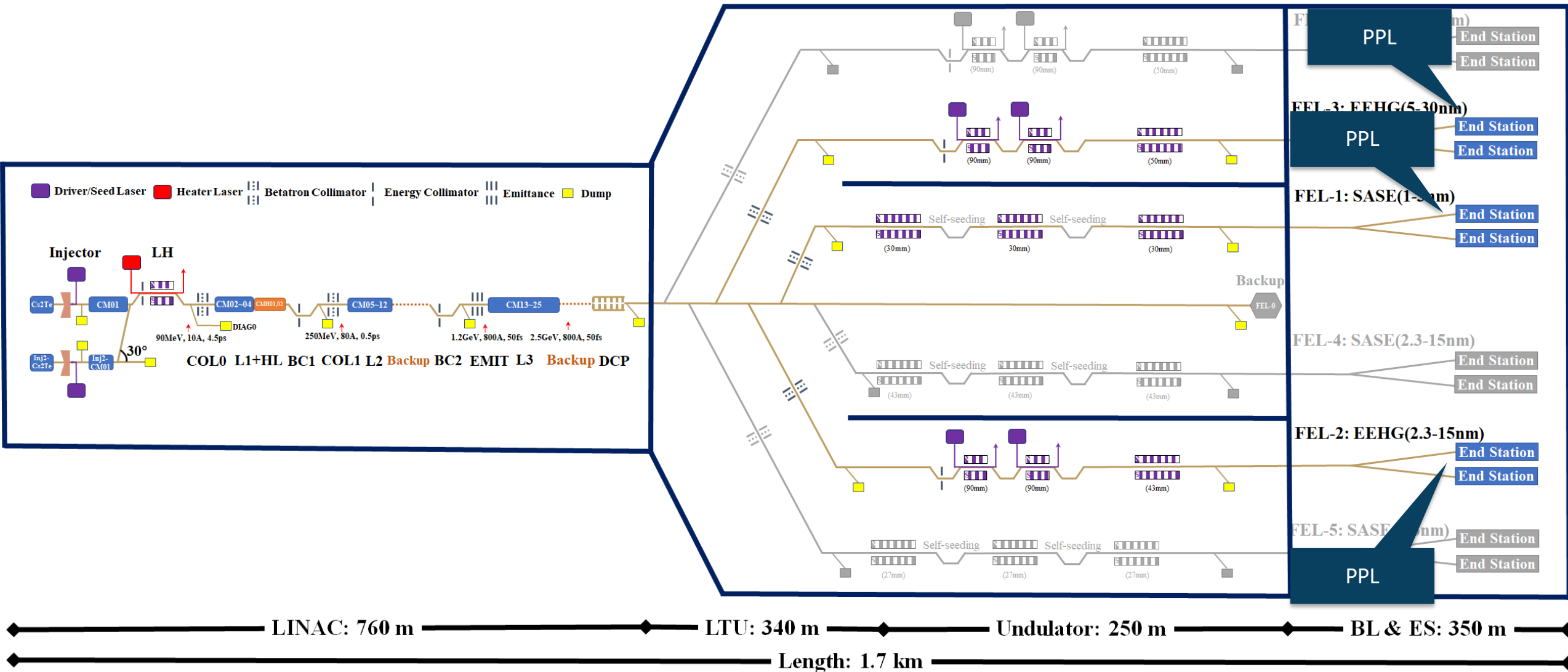


Rebernik Ribič, Primož, et al. "Coherent soft X-ray pulses from an echo-enabled harmonic generation free-electron laser." *Nature Photonics* 13.8 (2019): 555-561.

	FEL-3 & 4	
Rep. rate (kHz)	1	
wavelength (nm)	λ_1 267	λ_2 266-295
Pulse energy (μJ)	20	20
Beam radius (mm)	0.5	
Pulse duration (fs)	100	

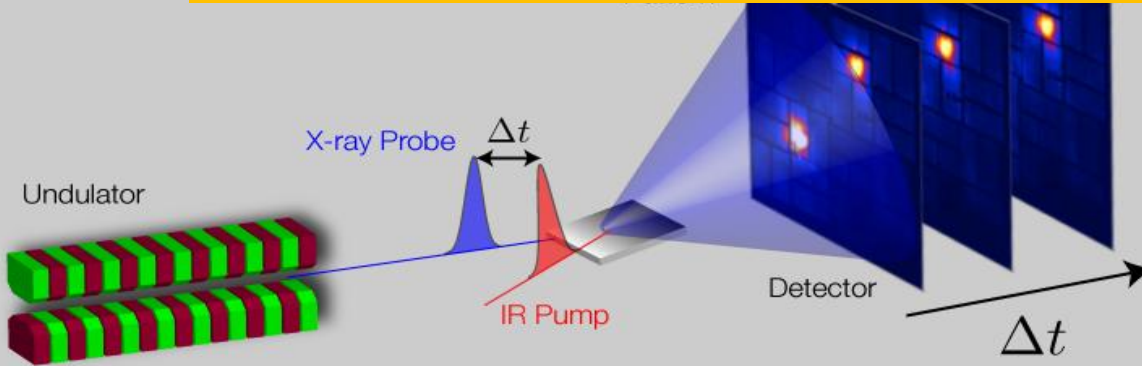
- **Nonlinear frequency conversion to the required wavelength range, the pump pulse energy should be in multi-mJ level;**
- **The bottom line: Seed lasers will be running at 1 kHz with mature Ti : Sapphire technology;**
- **Advanced seeding schemes are under development with which one to two order-of-magnitudes less pulse energy is required and seeding at 100 kHz is feasible.**

Pump-probe lasers



Pump-probe laser requirements for S³FEL

Time-resolved experiments



- Wavelength: EUV, VUV, UV, visible, IR, THz
- Repetition rate: up to 100 kHz;
- Pulse energy: multi-mJ;
- Average power: 100-1000 W
- Pulse duration: < 50 fs
- Tightly synchronize with X-ray (< 10 fs)

75%

100
kHz

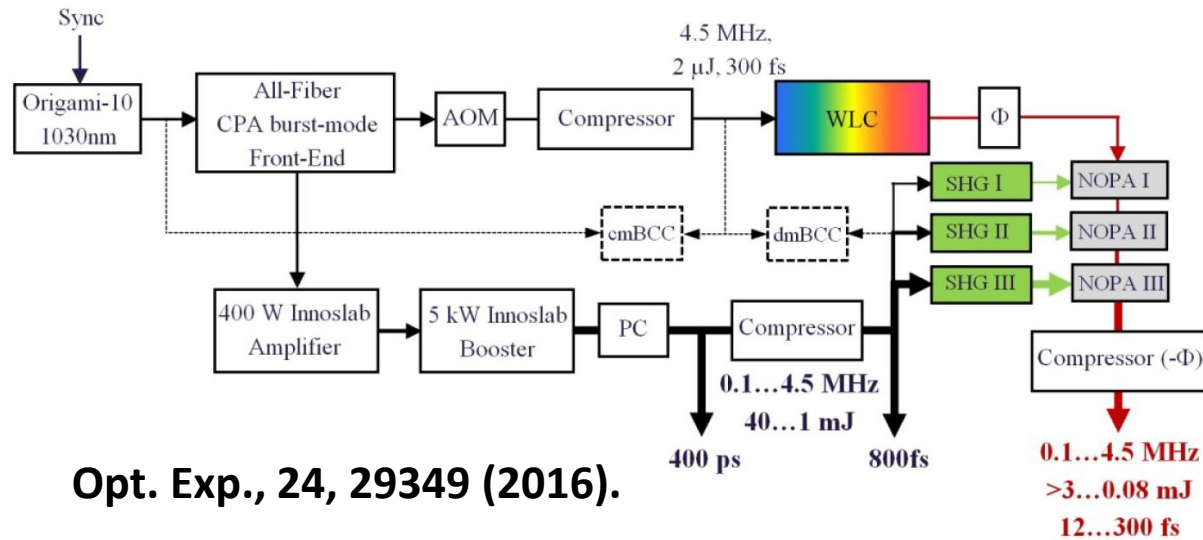
1
mJ

50
fs

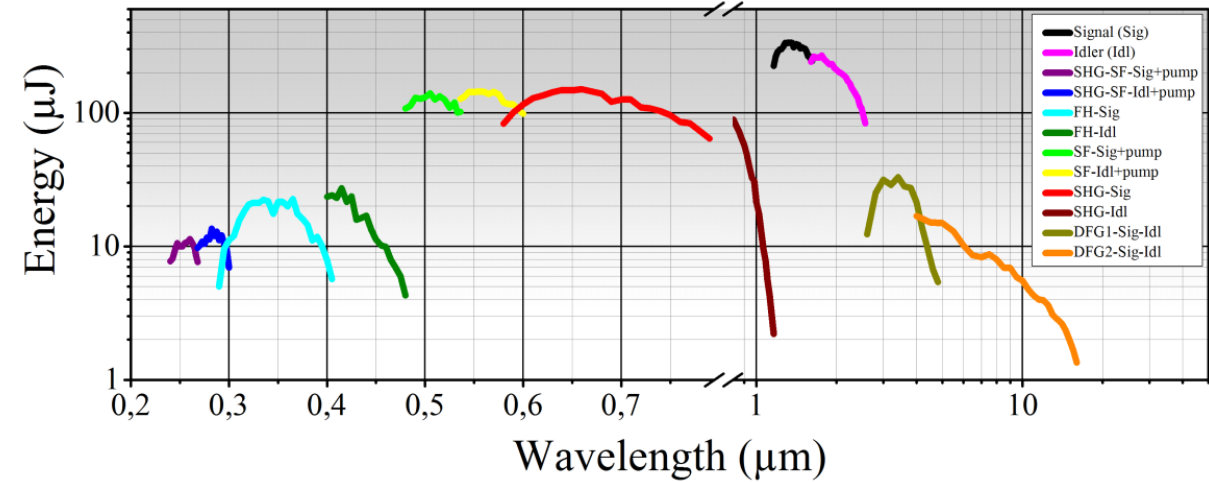
10
fs

Eu-XFEL pump-probe laser (OPCPA)

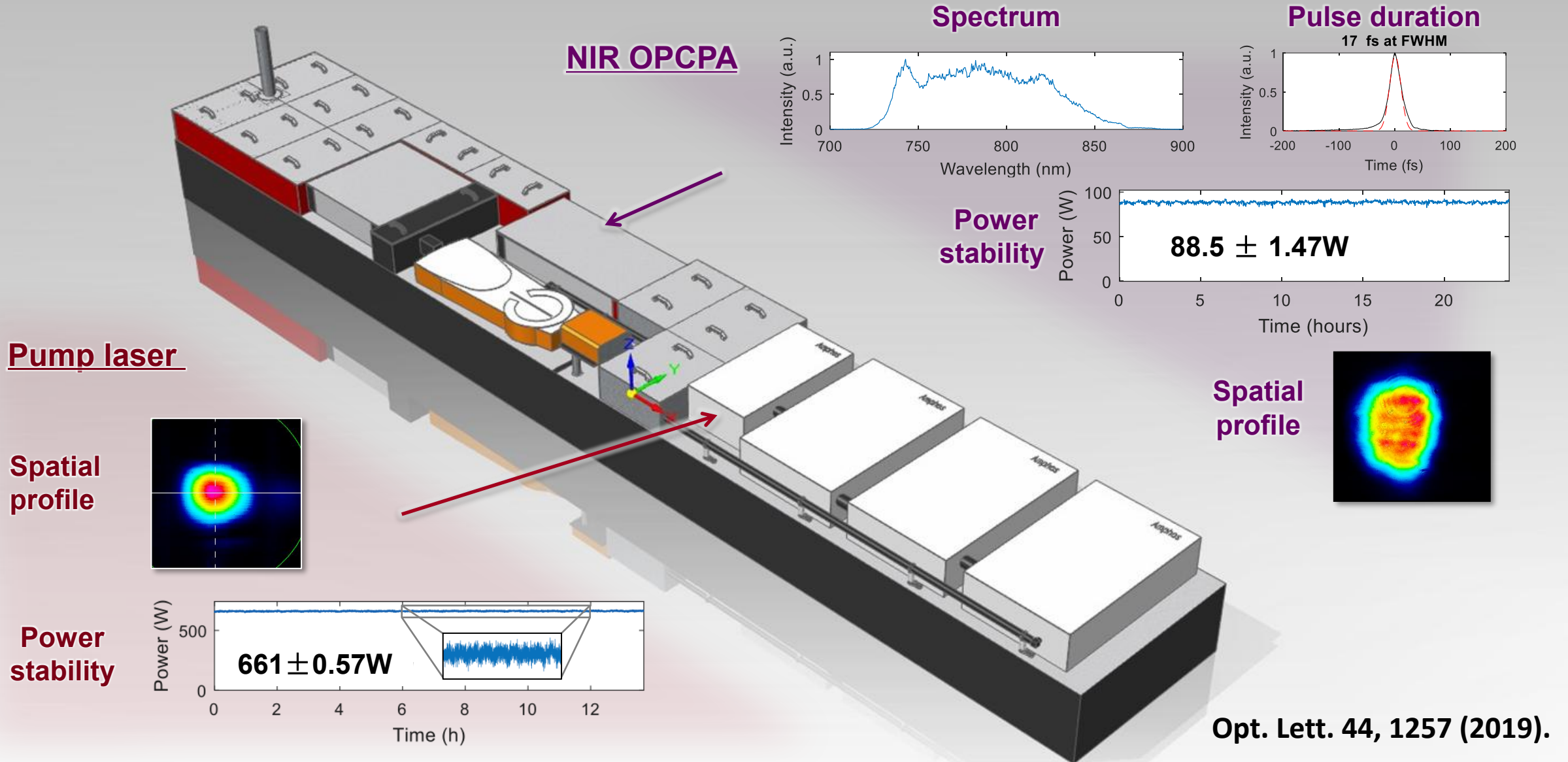
OPCPA: optical parametric chirped-pulse amplification



Opt. Exp., 24, 29349 (2016).

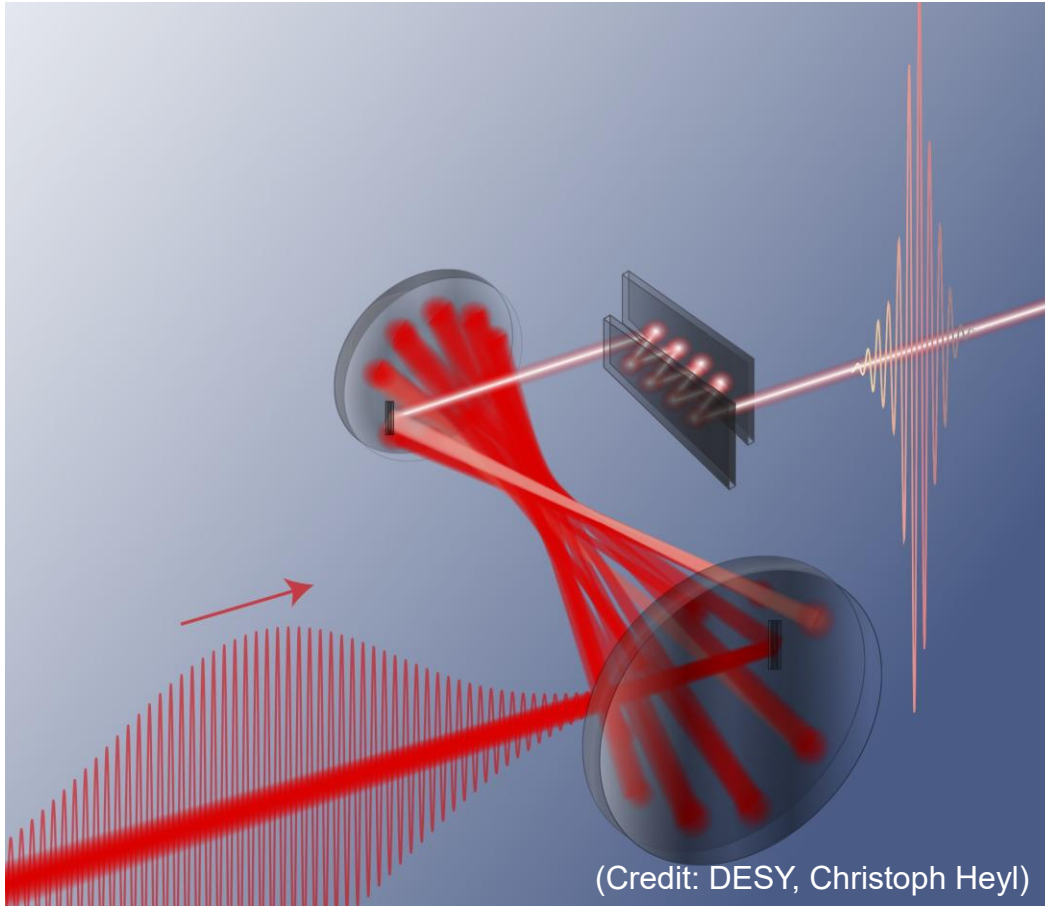


LCLS-II pump-probe laser (OPCPA)



A disruptive technology is on the way...

MPC: multi-pass compression with about 10 times higher efficiency



able (Amplitude Systems 2019). We will start a research program exploring these options. If this is successful, we plan to immediately implement these schemes in both laser systems (FLASH1 and FLASH2). This would allow us to deliver multi mJ-level 100 fs level pulses at MHz repetition rates centered at 1030 nm to the MODs and instruments, opening up new opportunities for both experiments and wavelength conversion stages. At the MOD stations then

FLASH 2020+

Making FLASH brighter, faster and more flexible
Conceptual Design Report

Deutsches Elektronen-Synchrotron DESY
A Research Centre of the Helmholtz Association



The map displays the layout of DESY facilities, with various laser and accelerator locations marked. Key areas include:

- FLASH2 FEL pump-probe lasers** and **FLASH1 FEL pump-probe lasers** at the top center.
- FEL Pump-probe laser R&D** and **Laser-plasma acceleration (collaborators: DESY accelerator division)** in the upper middle.
- Comb-laser for spectroscopy** in the center.
- Lasers for attosecond physics (collaborators: FS-Atto)** and **Lasers for HHG (collaborators: UFOX)** on the right side.
- XFEL FEL pump-probe lasers (collaborators: XFEL laser group)** on the left side.
- FLASH pulsed laser sync.** and **XFEL pulsed laser sync.** in the lower left.
- FLASH gun lasers** and **XFEL gun lasers** in the lower left.
- LPA: KALDERA Laser** and **LPA: Angus Laser 300TW** in the lower right.
- gun laser R&D** and **ARES gun laser** in the center-right.
- precision spectroscopy** and **THz electron acceleration** in the center-right.
- FLASH pp-lasers** and **AMO** in the upper right.
- waveform synthesizer**, **HHG driver**, and **THz acc. R&D** in the upper right.
- THz sampling** and **THz streaking** in the upper center.
- Pump-probe laser R&D** in the center.
- AMO** and **chemistry** in the center-right.
- electron acc.** in the center-right.
- CLASS5** logo in the center-right.


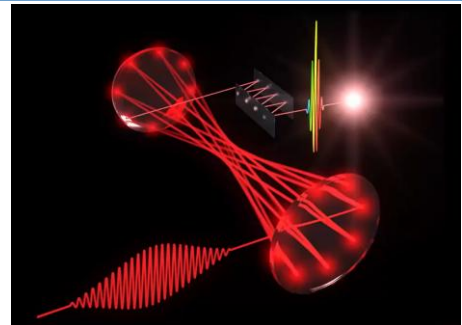
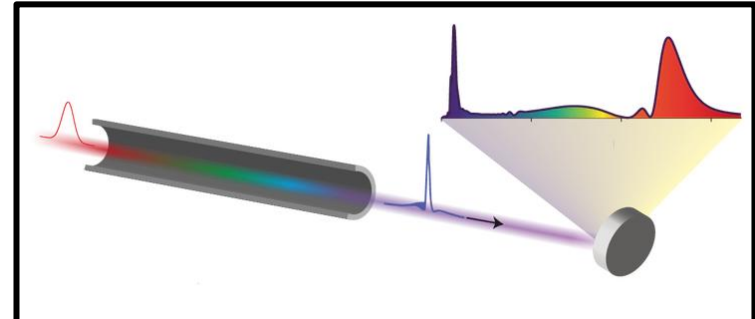
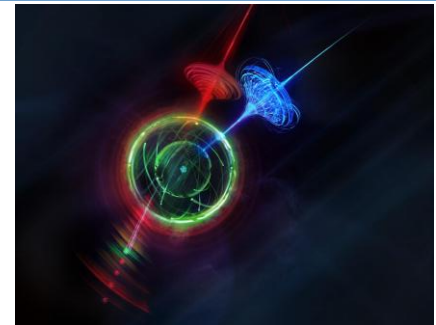
```
graph TD; A[Picosecond laser] --> B[Nonlinear spectral broadening]; B --> C[Optional wavelength conversion]; C --> D[User experiment];
```

The flowchart illustrates the experimental setup for generating a quantum state. It consists of four sequential steps, each represented by a green rectangular box with a dark blue border. The steps are connected by red arrows pointing downwards. The steps are: 1. Picosecond laser, 2. Nonlinear spectral broadening, 3. Optional wavelength conversion, and 4. User experiment.

Our design choice and R&D Strategy

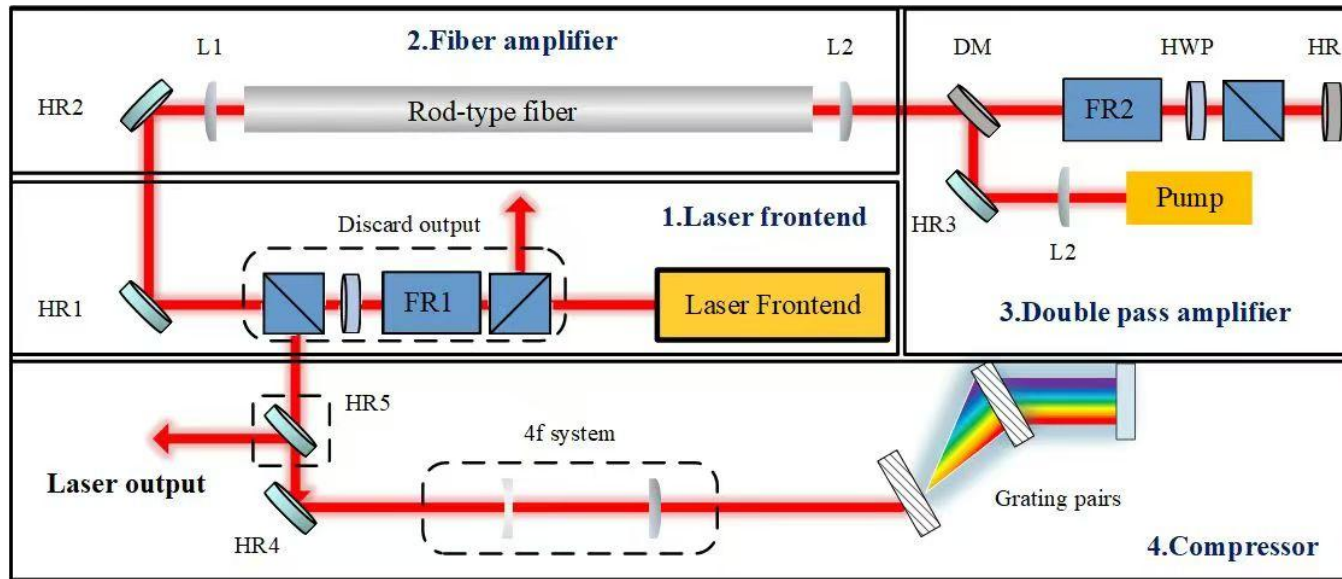
Multi-mJ, high average power: 100 – 1000 W	Short pulse duration: < 50 fs FWHM	Broad λ tunability: THz \rightarrow EUV	Tight sync.: < 10 fs timing jitter
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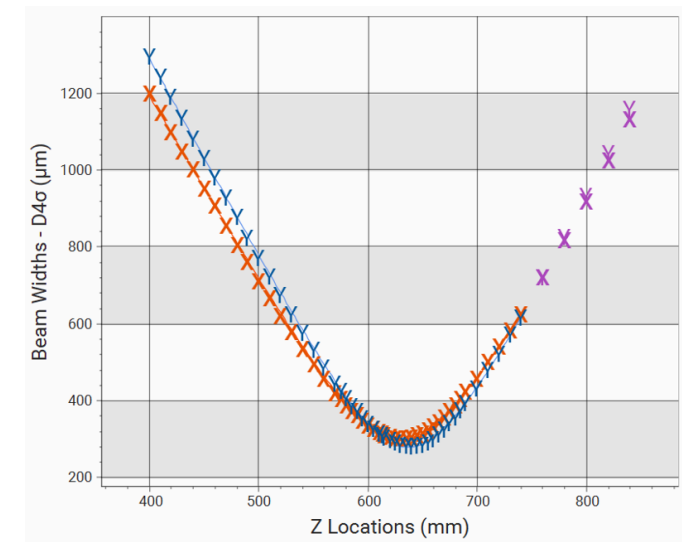
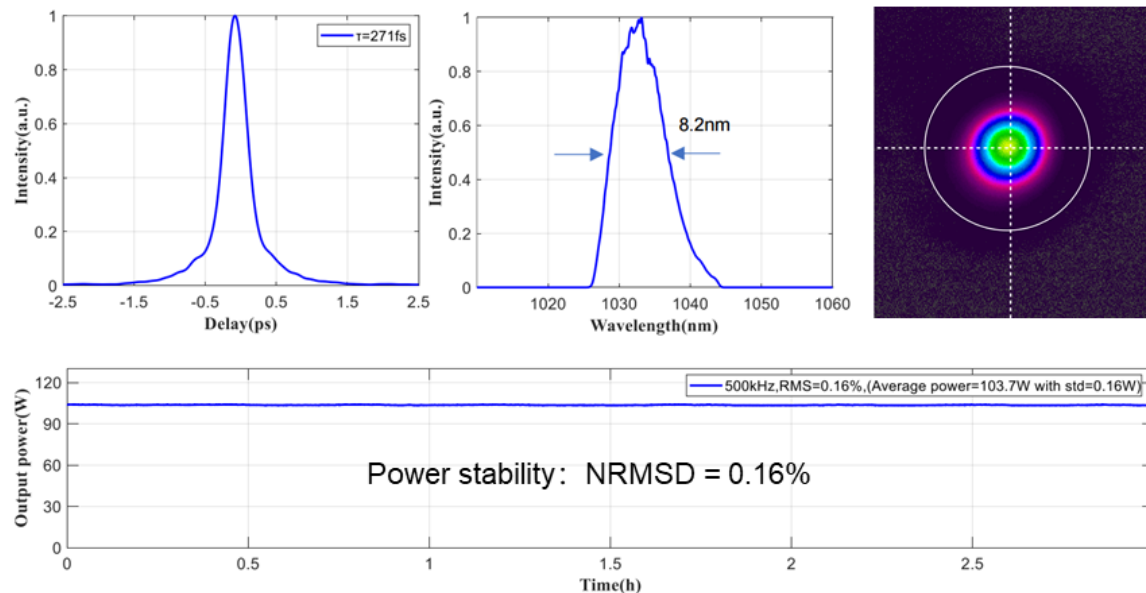
Yb platform	post compression: multi-pass cell & cascade compression	Commercial OPA, Home-built DWG, HHG	Pulsed laser-based synchronization
			

Since most of the solution is not commercially available, we need to put tons of efforts on self-development and engineering

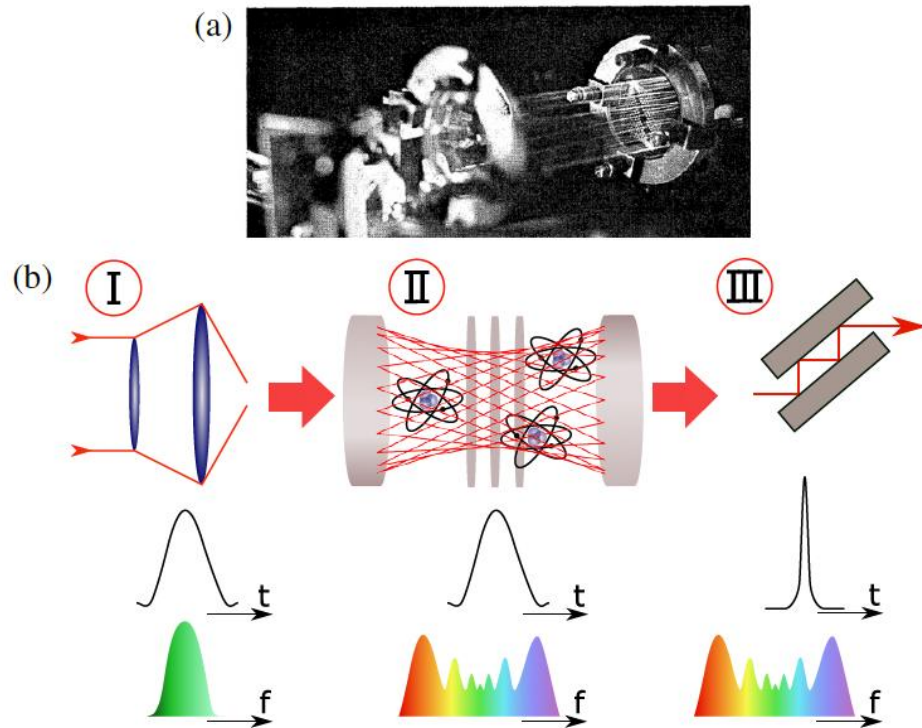
~200W Yb fiber amplifier



Center wavelength	1034 ± 5 nm
Maximum output power	180 W
Maximum pulse energy	400 μ J
Pulse duration	271 fs
Repetition rate	100 kHz - 1 MHz
Long-term power stability	0.16%
Beam quality, M^2	< 1.2

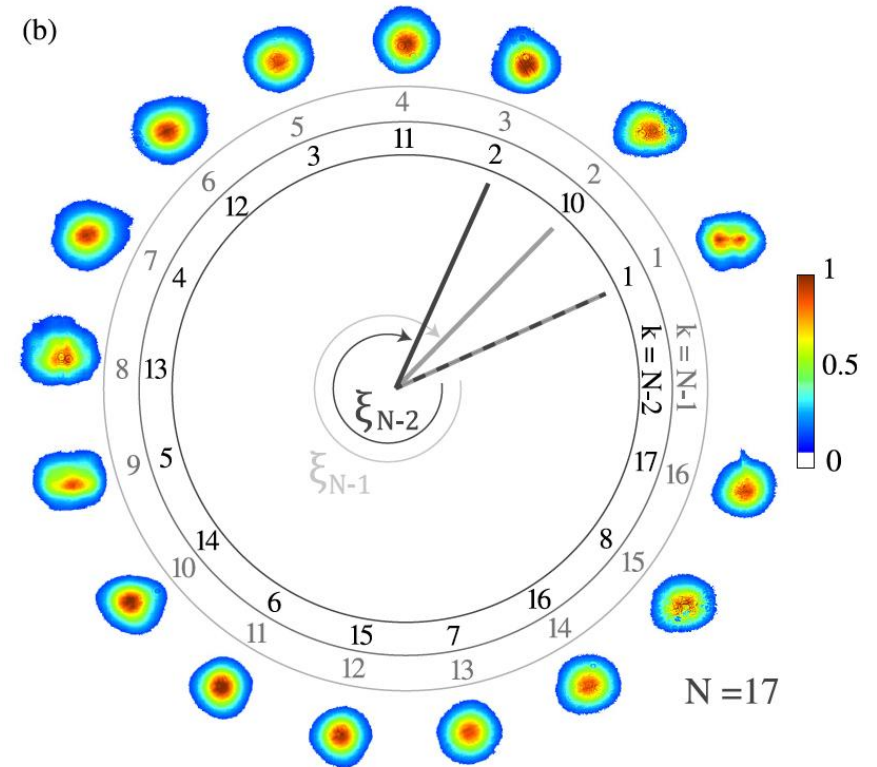


Post compression: multi-pass cell (MPC)



Key features:

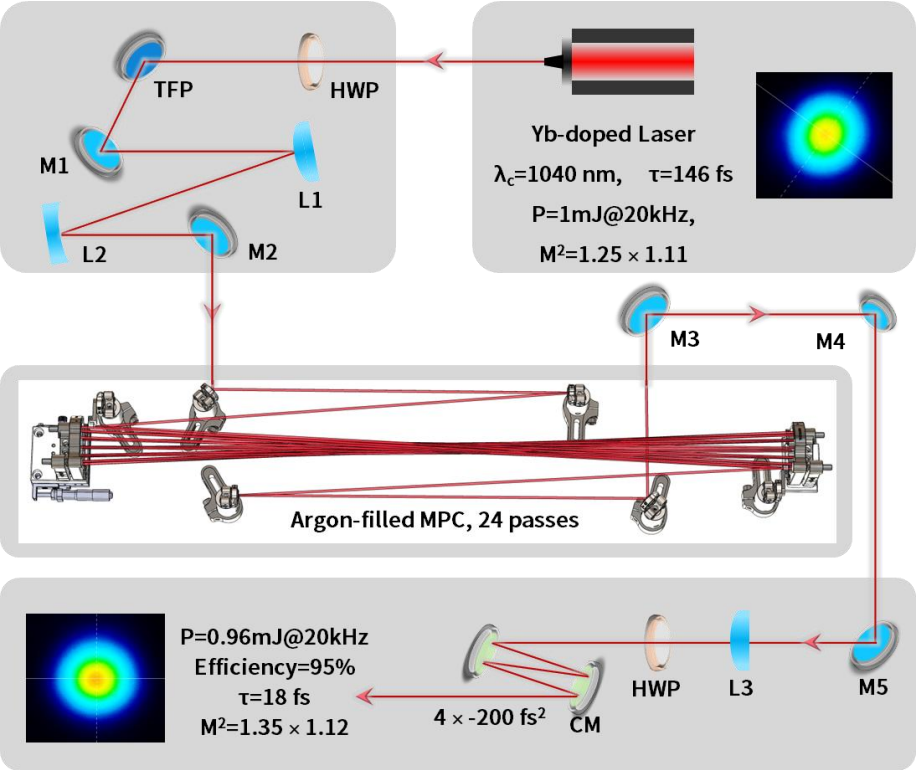
- High power compatibility
- Large compression factors
- High efficiency
- Excellent beam quality
- Compatible with large parameter range
-



Optica **9**, 197 (2022)

MPC for < 50 fs output

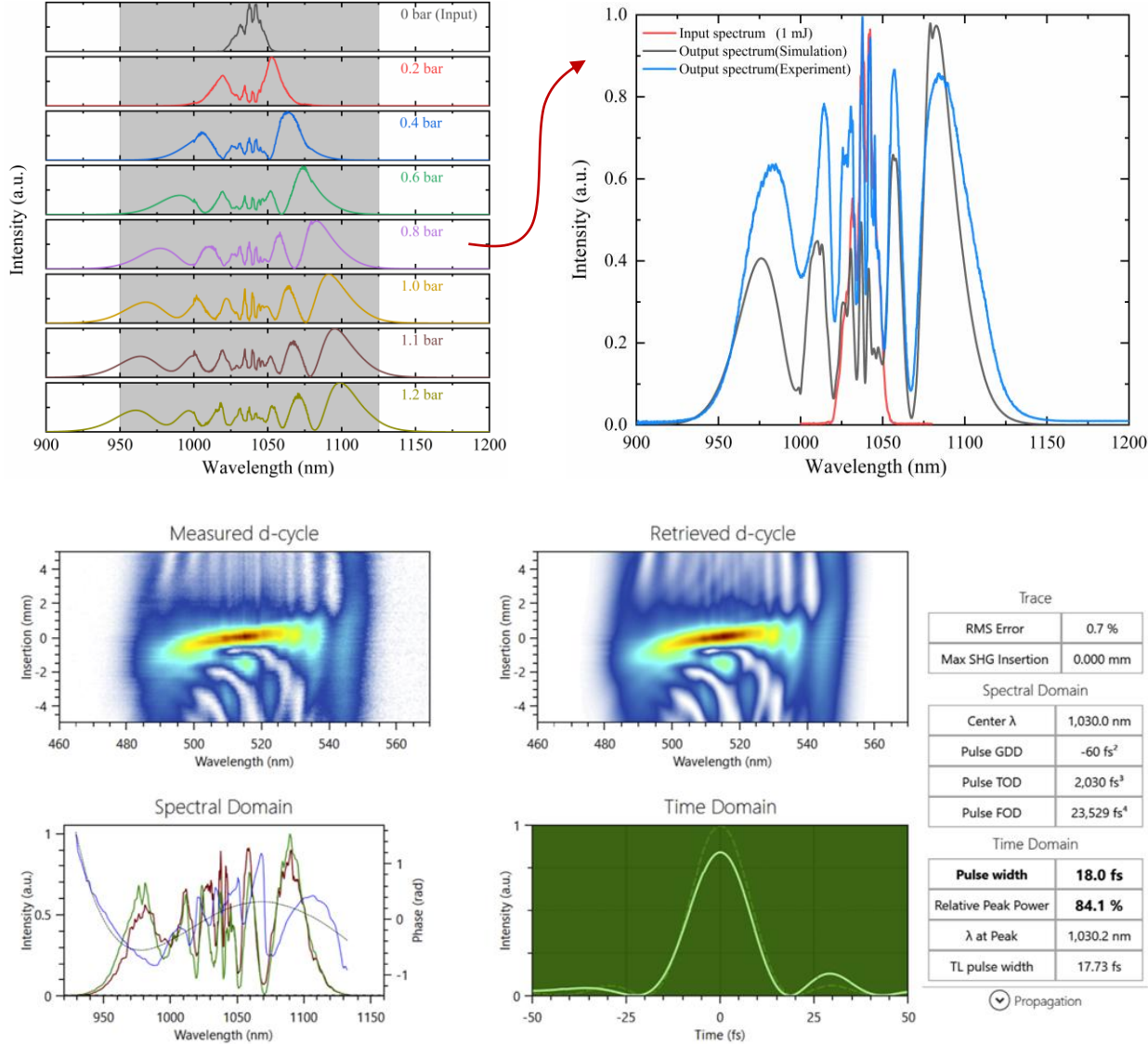
Setup



146 fs → 18 fs with 95% efficiency

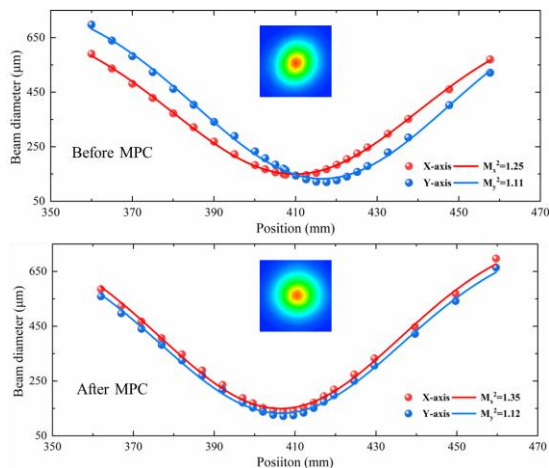
IEEE PHOTONICS TECHNOLOGY LETTERS, 37, 313 (2025)

Experimental results

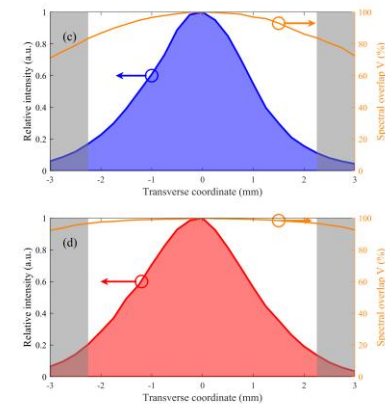
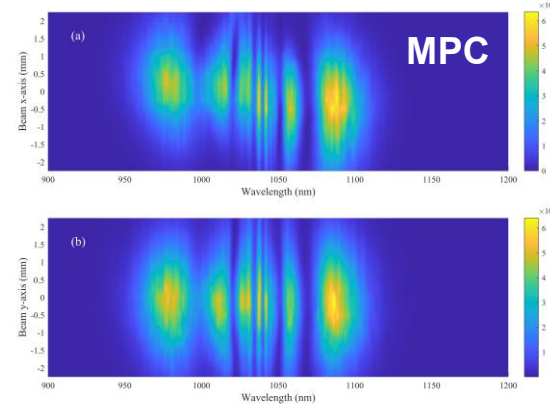
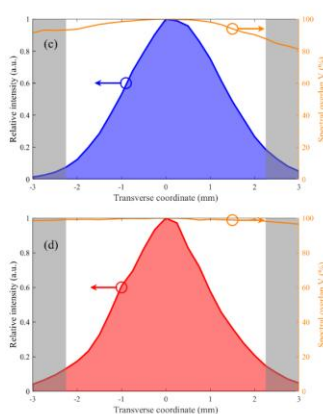
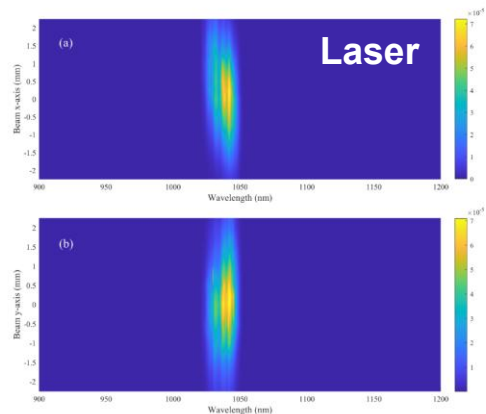


Performance of MPC output

Beam quality M^2

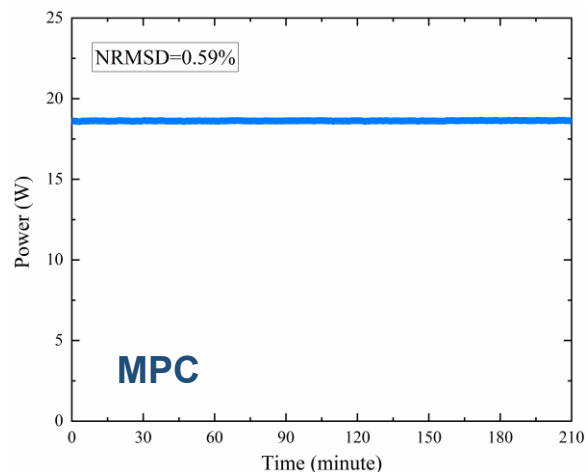
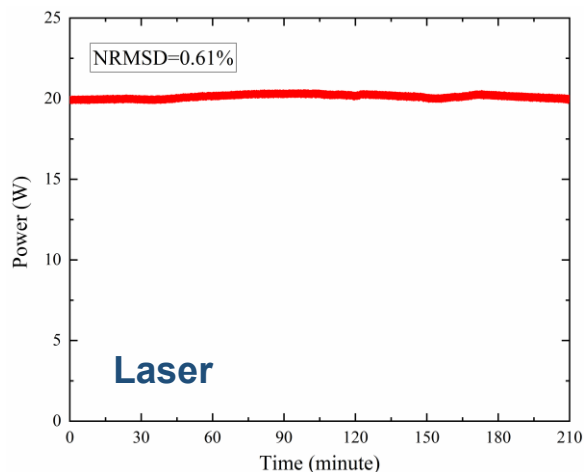


Spatial-spectral homogeneity

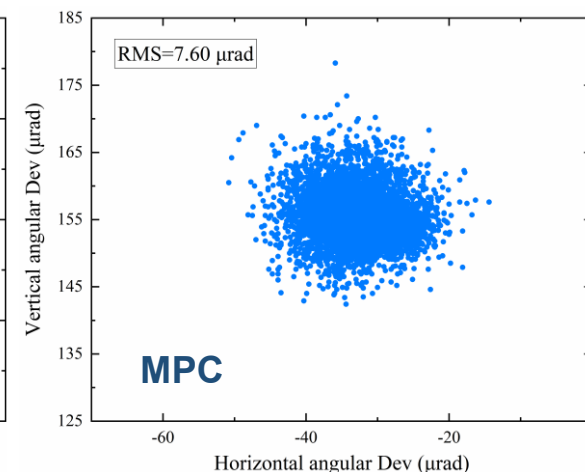
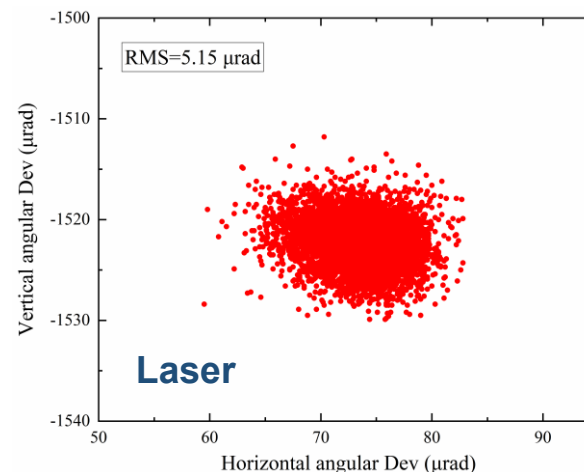


$$V_X = 95.9\%, V_Y = 98.9\%$$

Power stability

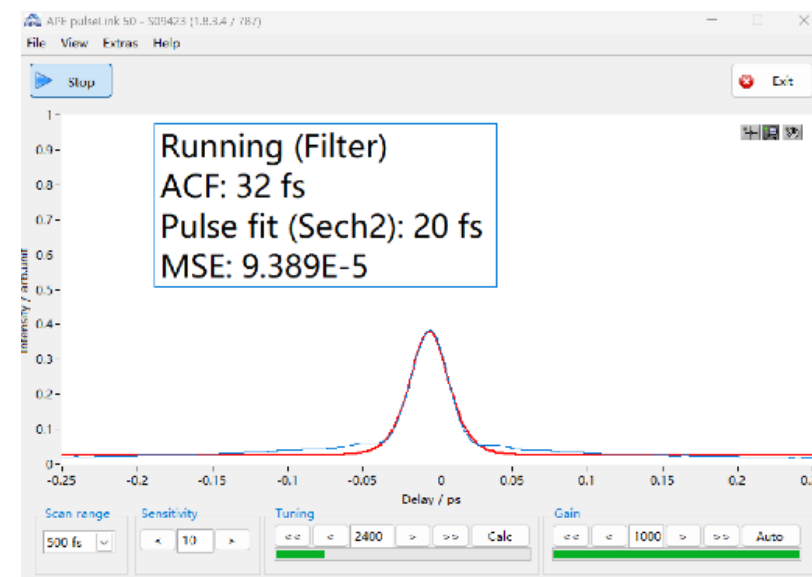
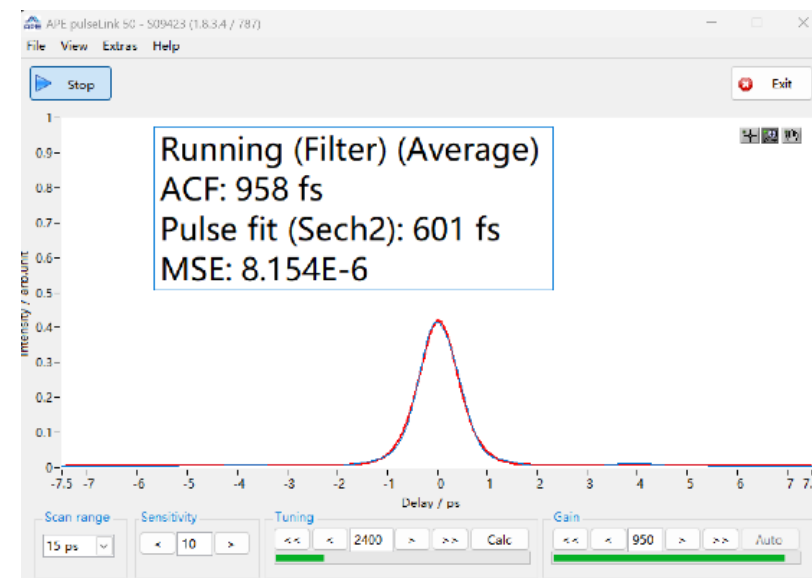
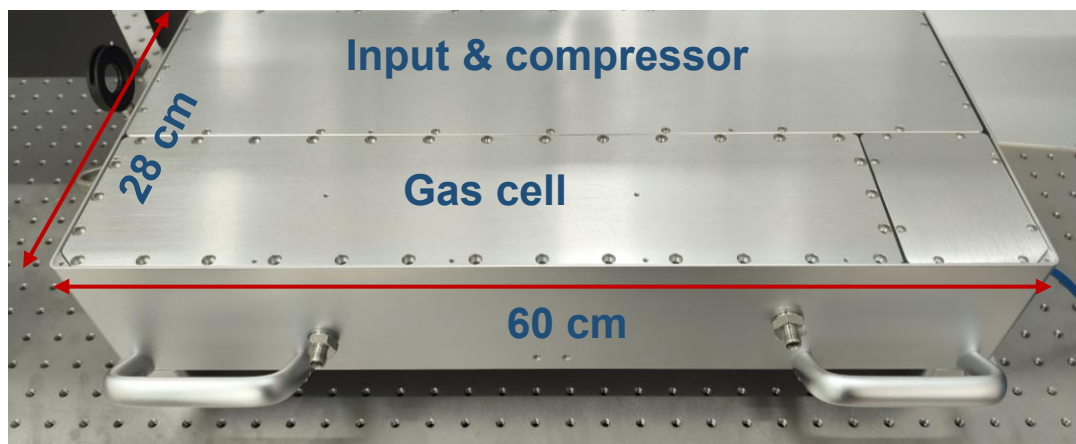
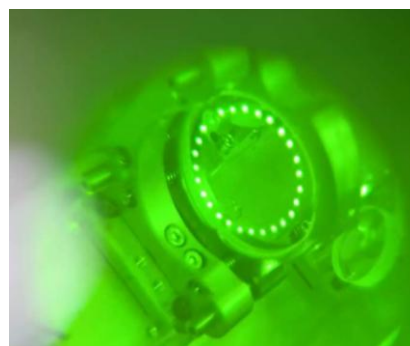
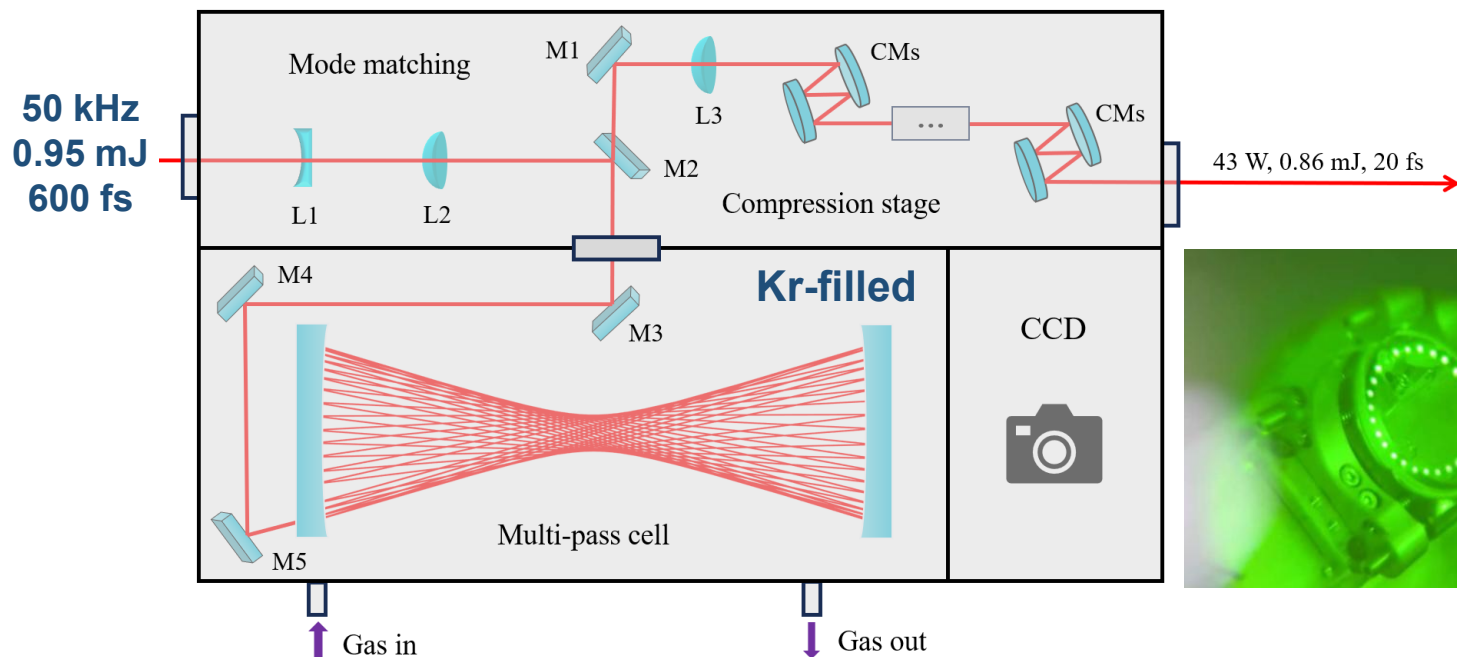


Beam pointing

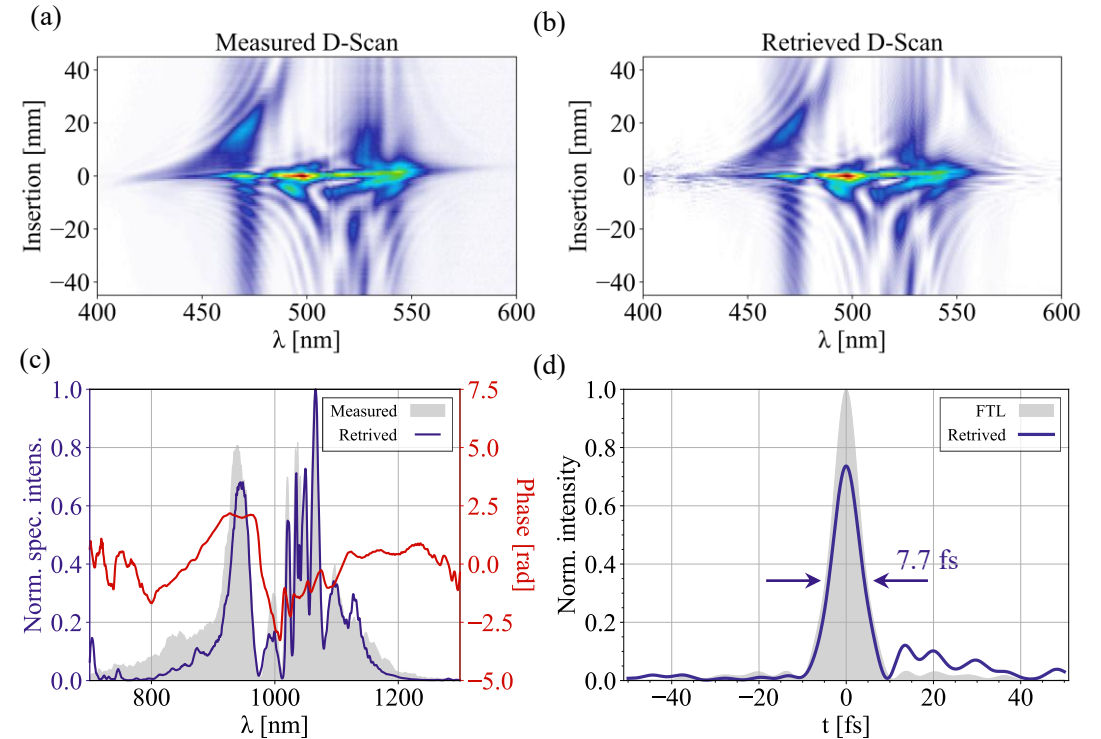
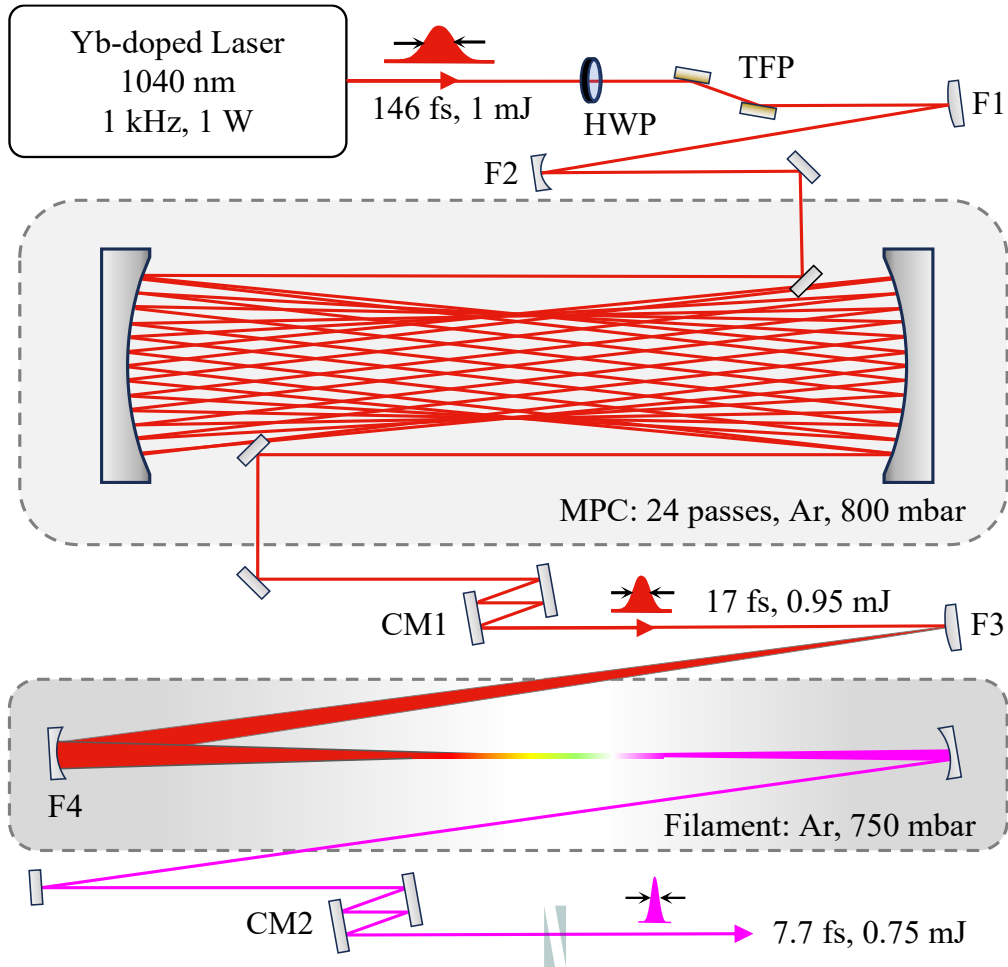


MPC maintains the power stability, pointing stability, beam quality, with homogeneous spectral broadening

30× compression single-stage MPC system



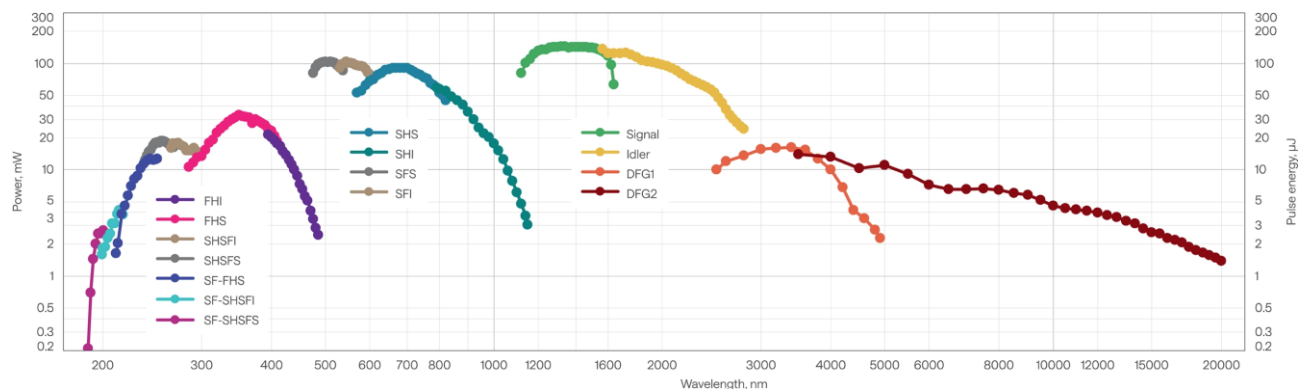
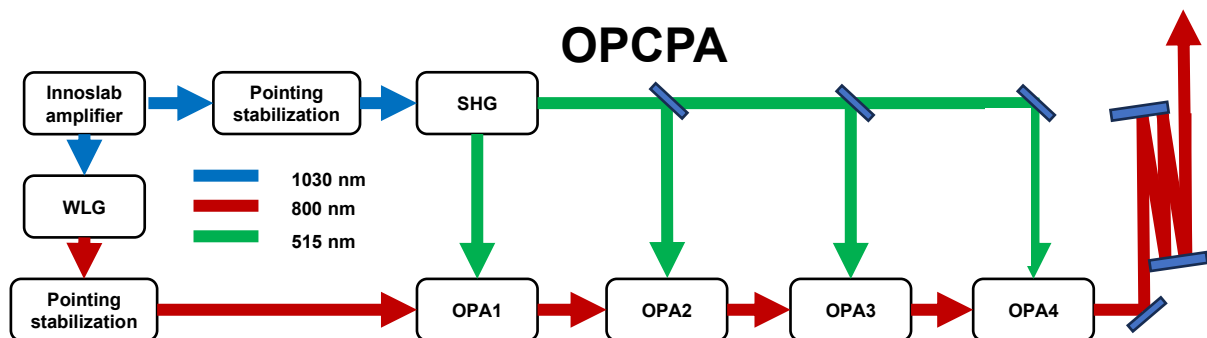
Cascaded post-compression: MPC + filamentation ----- sub 10 fs pulse generation



- **Pulse width:** from 146 to 7.7 fs;
- **Pulse energy:** 0.75 mJ;
- **Efficiency:** 75 % (total);

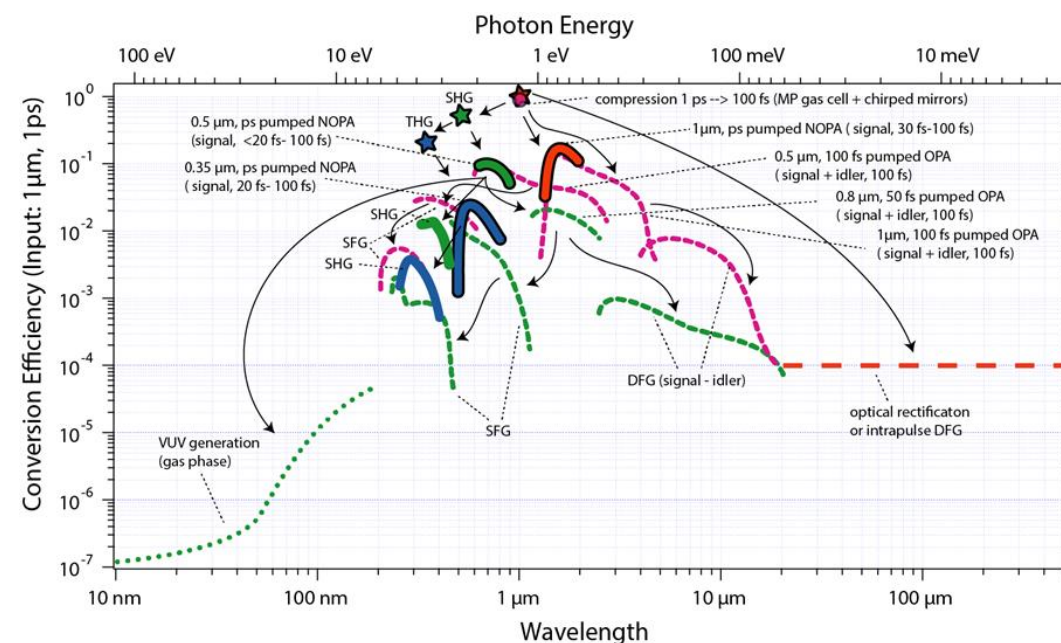
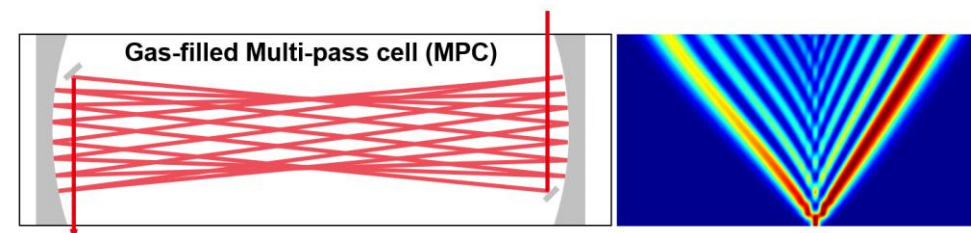
Wavelength extension: commercial OPA (down to 200 nm)

Yb-laser + OPCPA + OPA (800 nm)

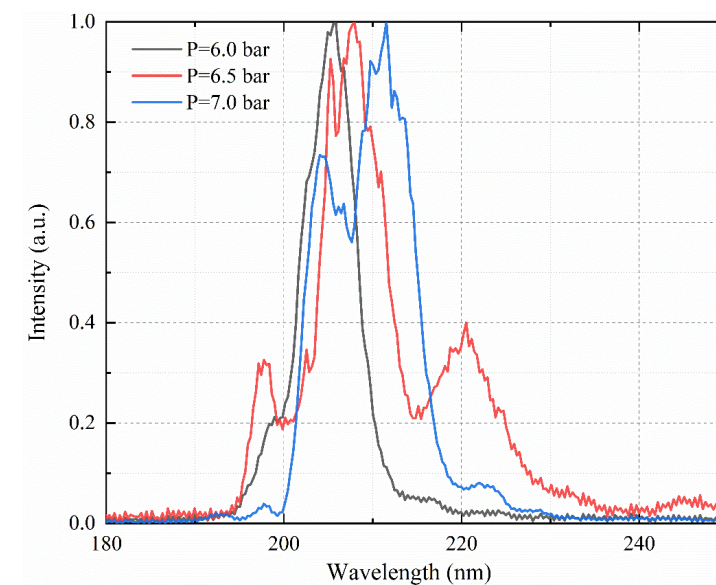
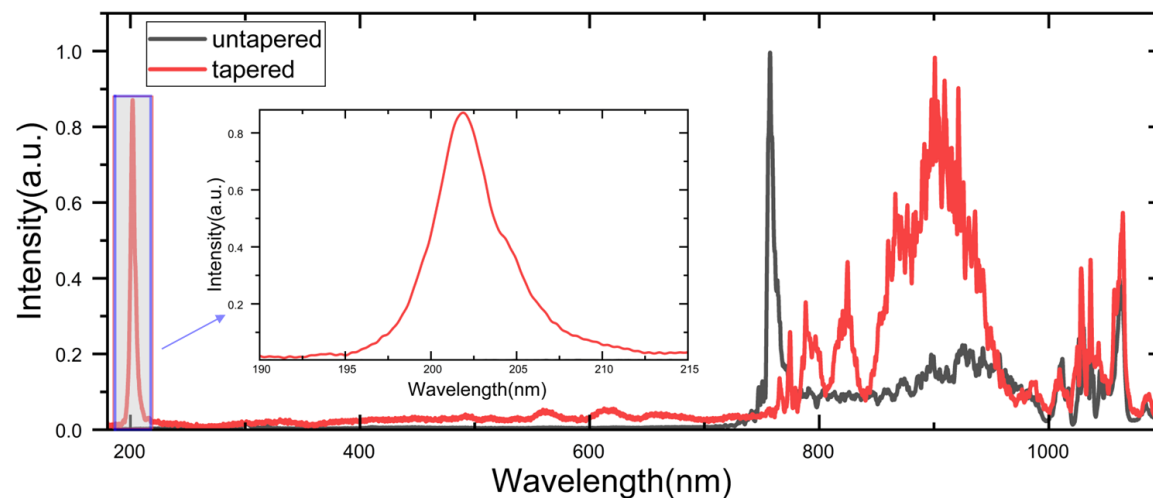
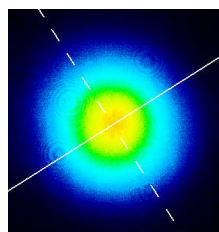
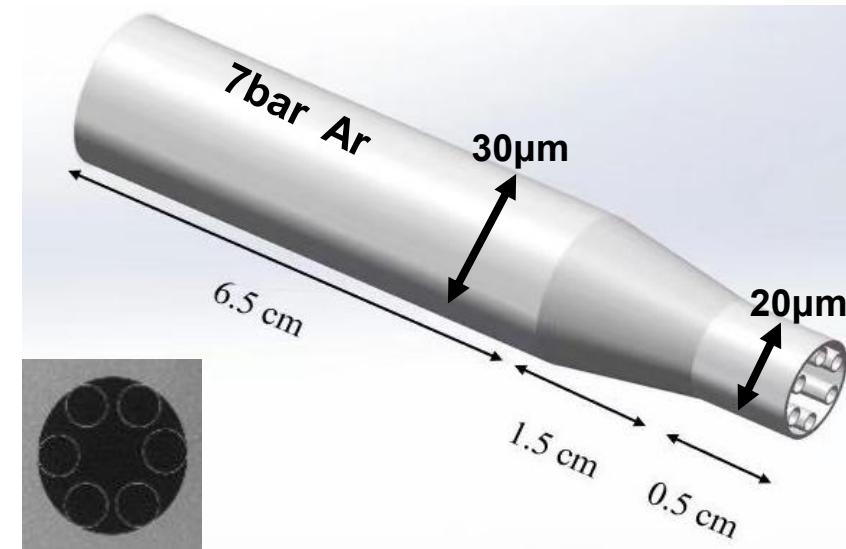
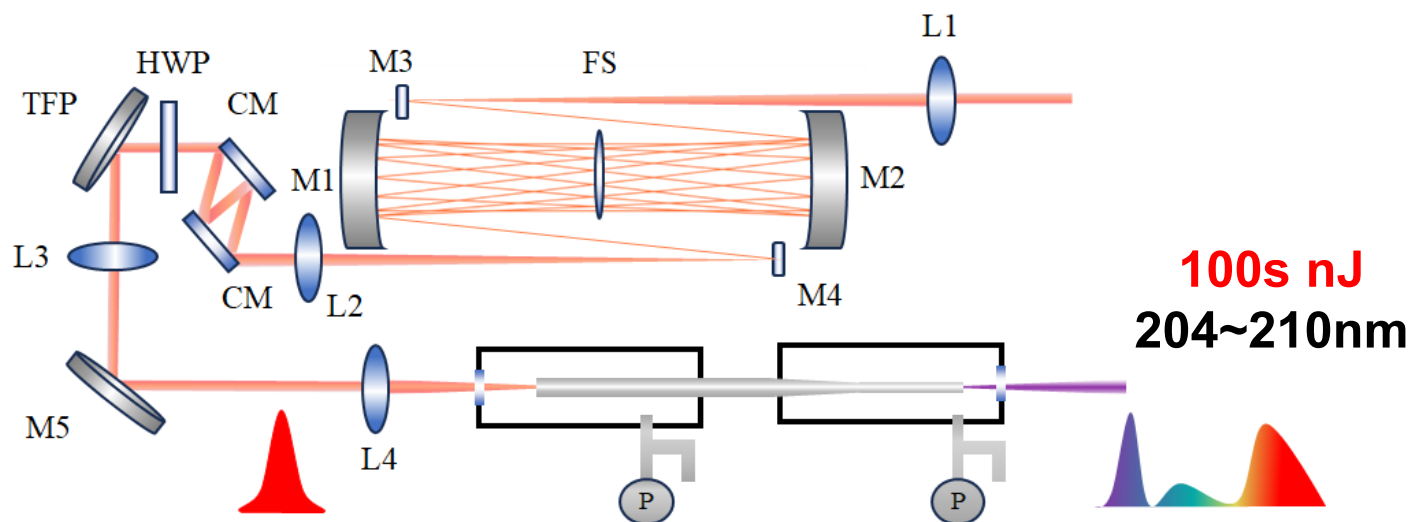


<https://lightcon.com/products/topas-prime-opa/>

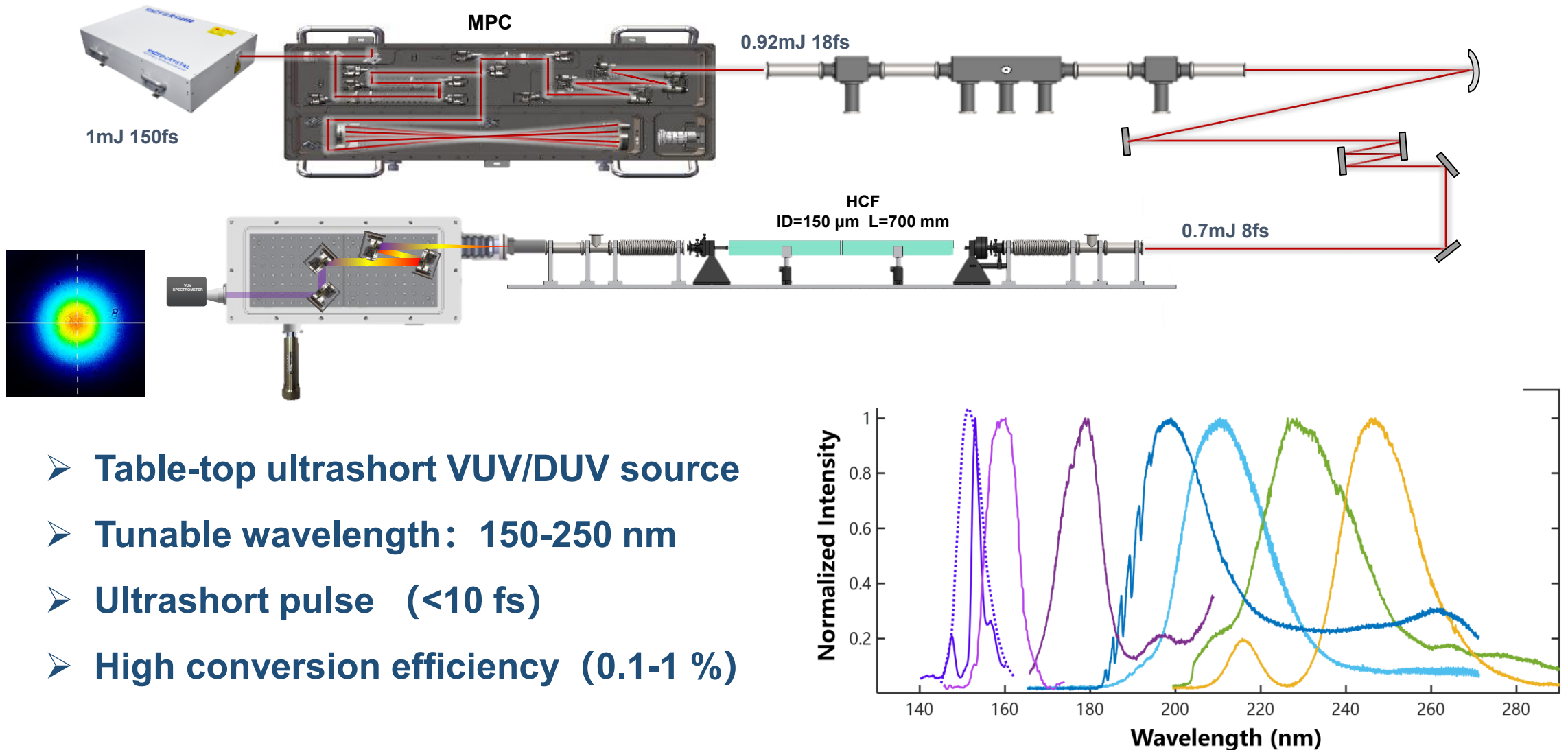
Yb-laser + MPC + OPA (1030 nm)



Wavelength extension to DUV (down to 200 nm): DWG with antiresonant PCF

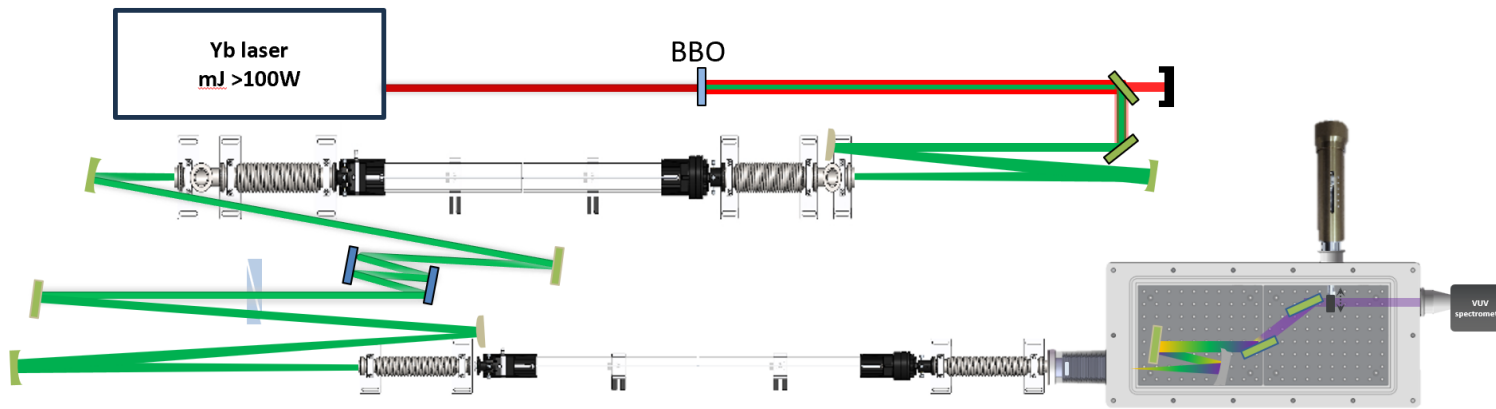


Wavelength extension to DUV & VUV (down to 150 nm): DWG with HCF



Wavelength extension to 120 nm: DWG with 515 nm pump laser

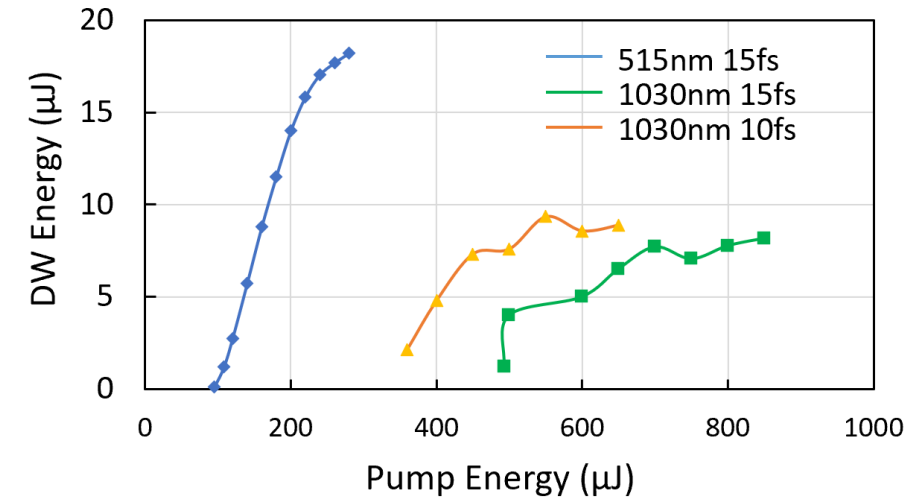
DWG pumped by 515 nm



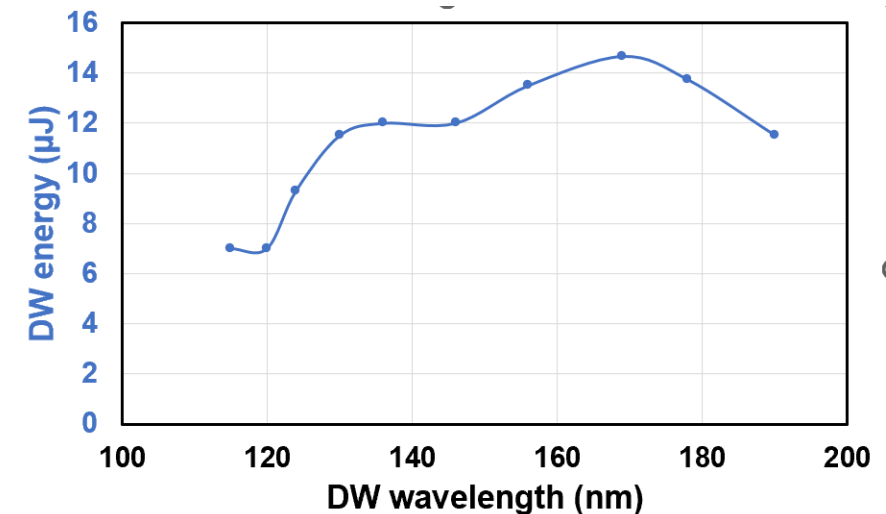
Pumping by green laser pulses:

- Shorter wavelength: ~120 nm
- Ultrashort pulse: < 10 fs
- Higher efficiency: 0.5-1.5%

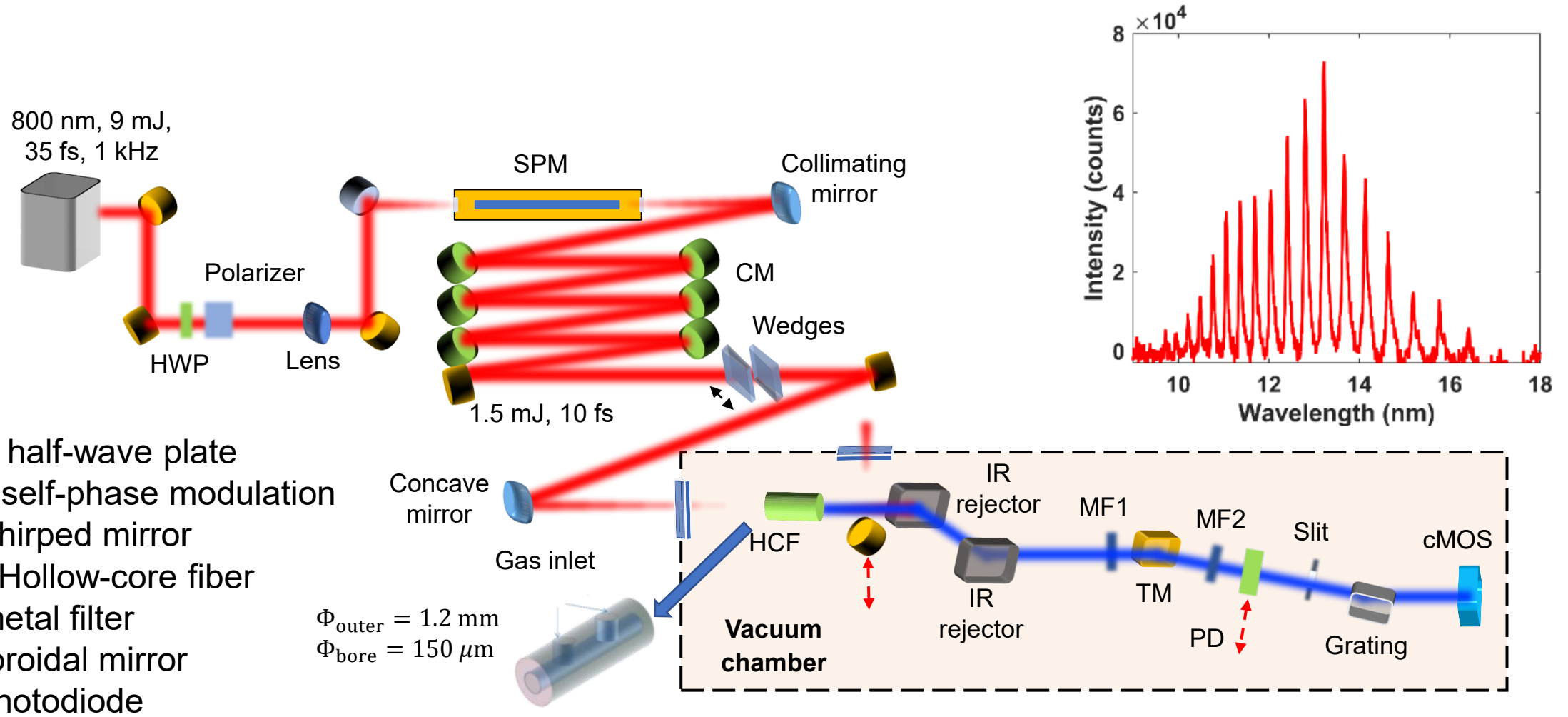
130 nm generation (simulation)



pulse energy of RDW (simulation)



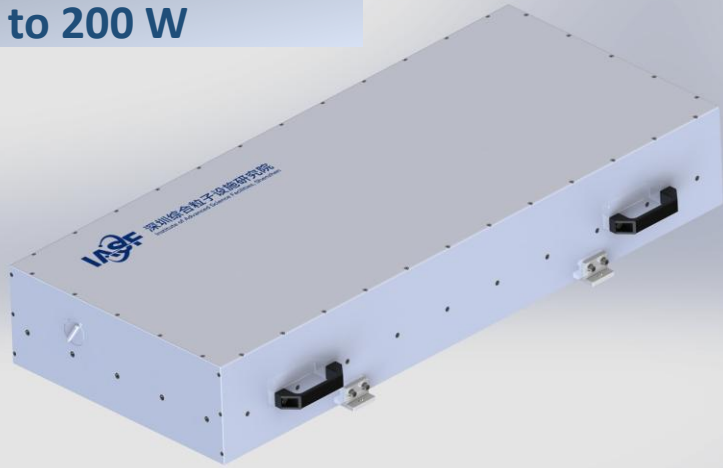
Wavelength extension to EUV: high-efficiency HHG system



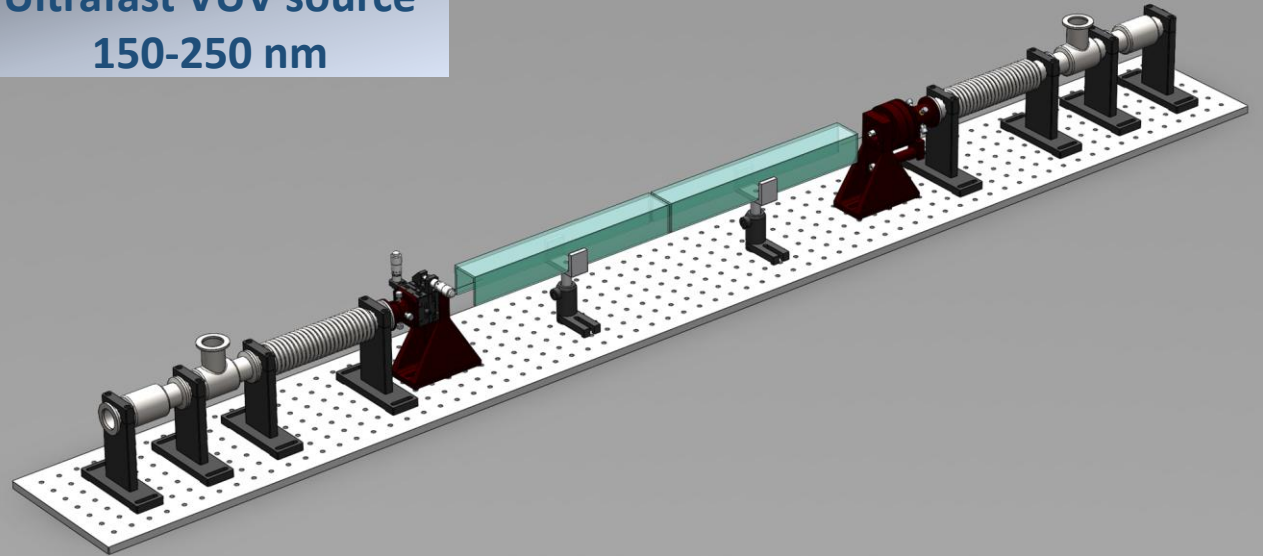
Post-compression & capillary HHG

Engineering & Commissioning Progress

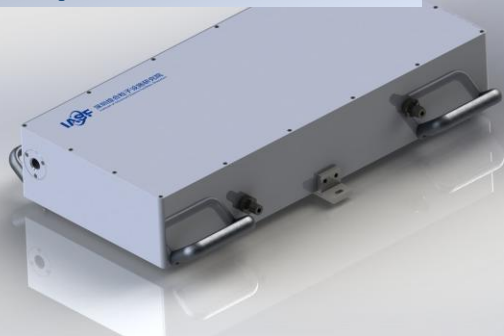
High power ultrafast laser
Up to 200 W



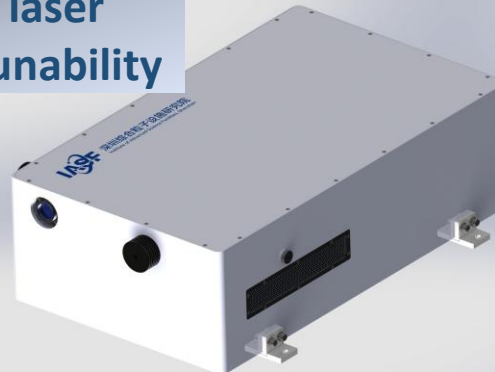
Ultrafast VUV source
150-250 nm



MPC compressor
*shortest pulse width < 10 fs



Mid-IR laser
6-20 μ m tunability



GMA
Pulse width < 40 fs



outline

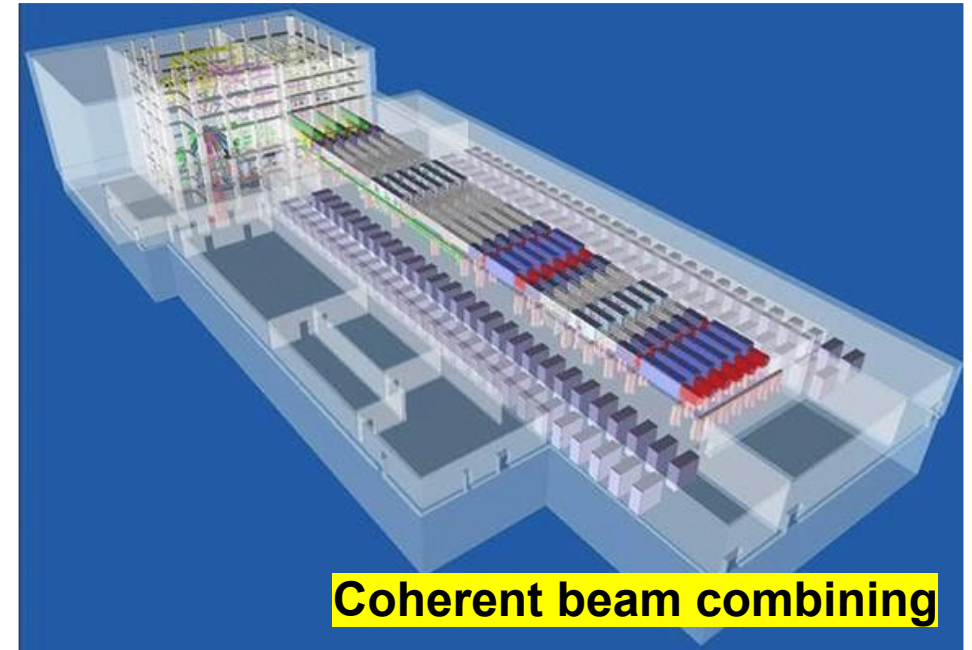
- What is a laser?
- What is an ultrafast laser?
- Mathematic description of an ultrafast laser pulse;
- Ultrafast lasers for S³FEL;
- **Femtosecond timing and synchronization**

Synchronization

**Synchronization is the
coordination of events to
operate a system in unison
[wiki]**



High-precision timing and synchronization applications

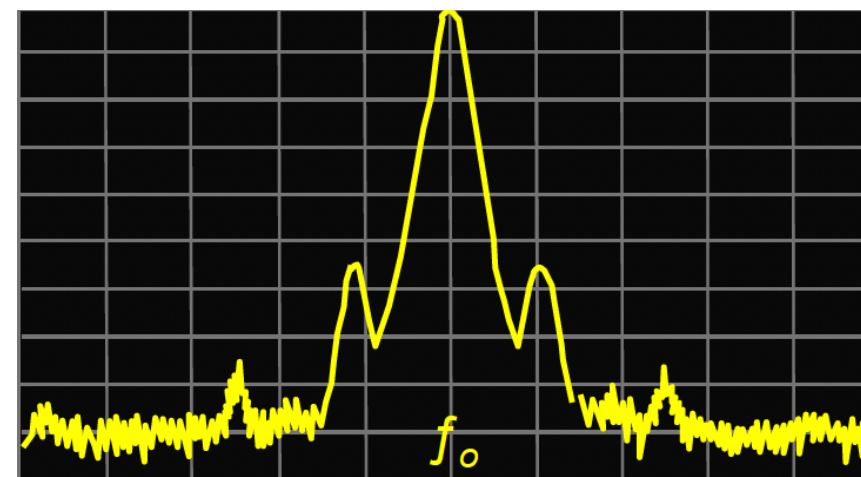
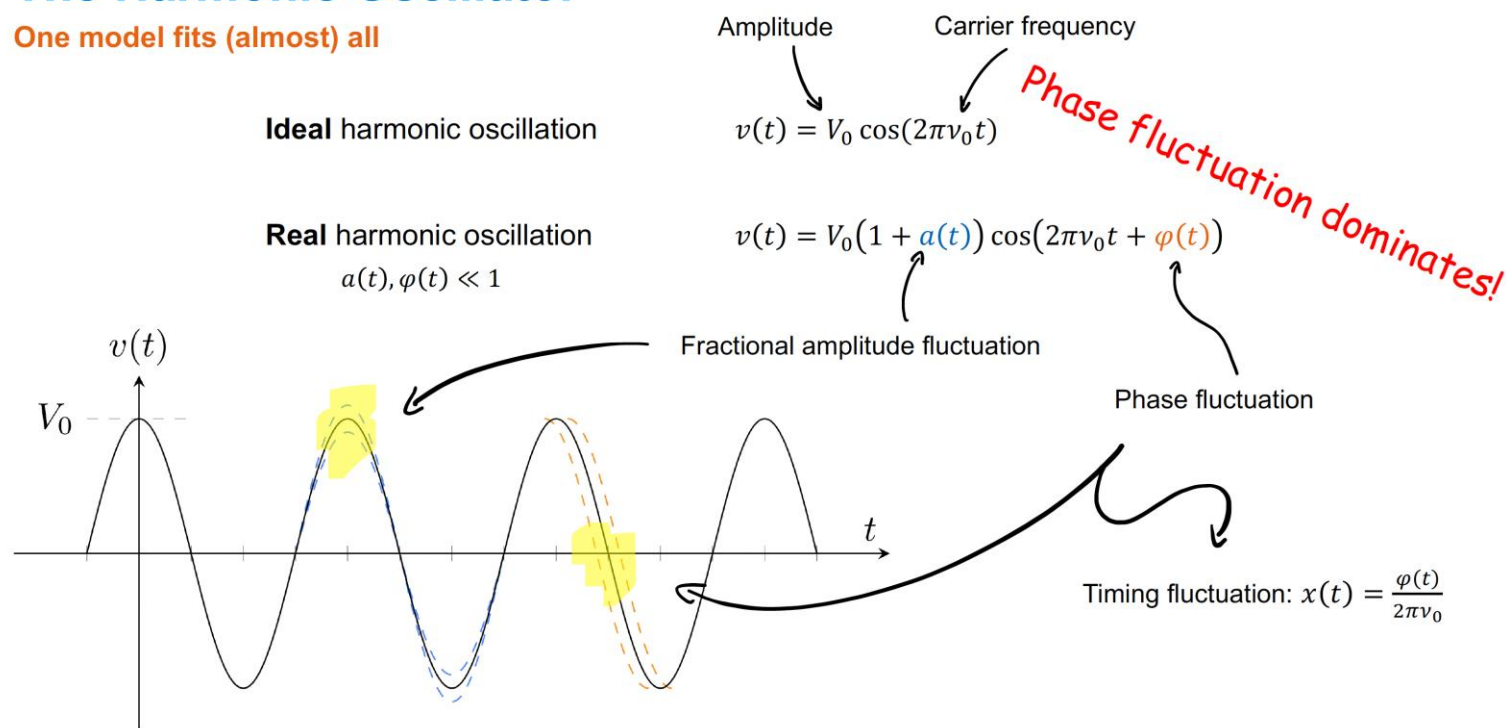


Phase noise of a real world harmonic signal

Phase noise is generally considered as the **short-term phase/frequency instability** of an oscillator or other RF/microwave component.

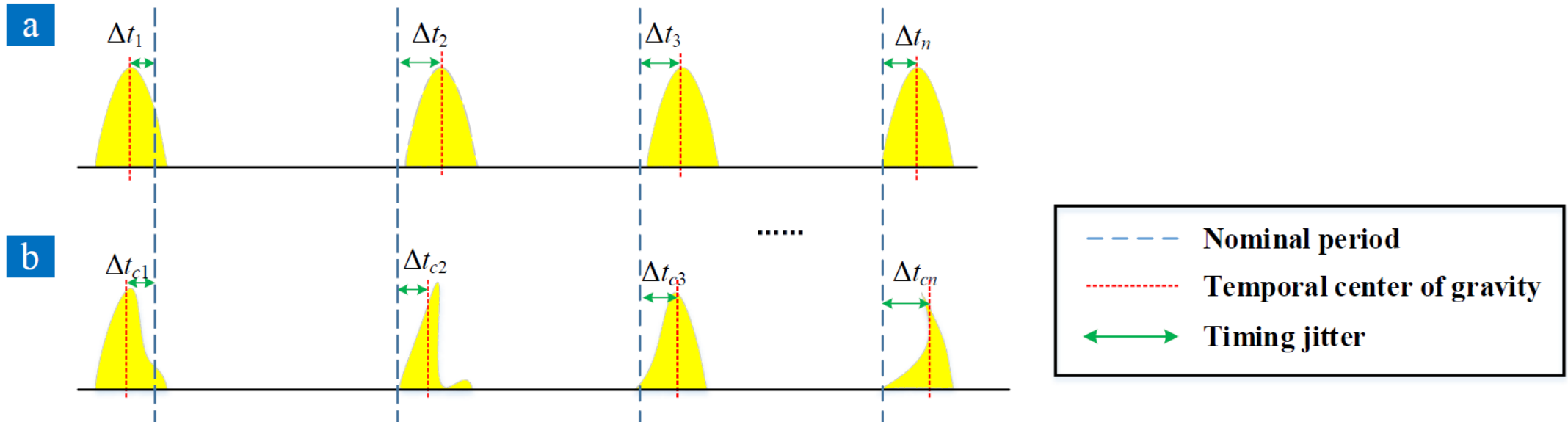
The Harmonic Oscillator

One model fits (almost) all



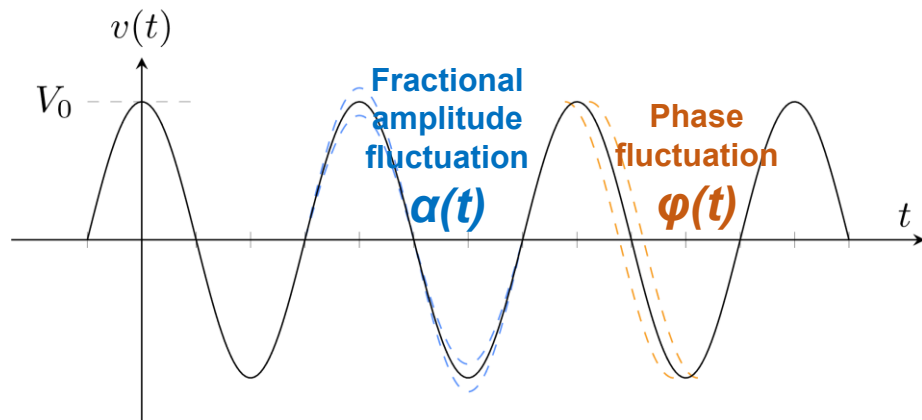
Timing jitter of pulsed signal

Timing jitter: deviation between the real arrival time of each pulse (timing) and its expected arrival time (nominal period)



$$T_{COG} = \int_{-\infty}^{+\infty} t |E(t)|^2 dt \bigg/ \int_{-\infty}^{+\infty} |E(t)|^2 dt$$

Quantitation of timing jitter and phase noise



$$v(t) = V_0(1 + \alpha(t))\cos(2\pi\nu_0 t + \varphi(t))$$

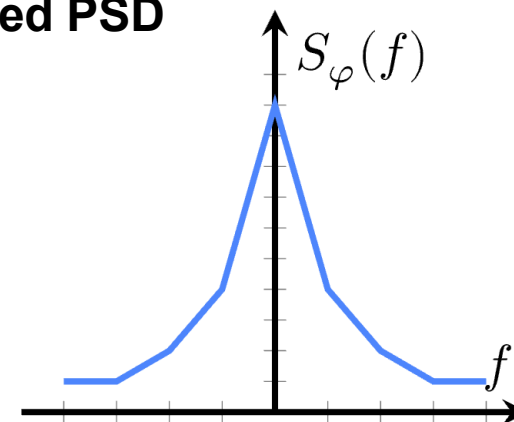
The phase noise power spectral density

$$S_\varphi(f) = \int_{-\infty}^{+\infty} e^{-j2\pi f\tau} \langle \varphi(t + \tau)\varphi(t) \rangle d\tau$$

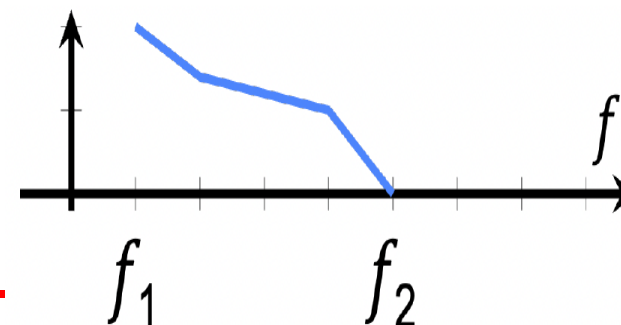
The timing jitter

$$t_{\text{jitter}} = \frac{1}{2\pi f_0} \sqrt{2 \cdot \int_{f_1}^{f_2} S_\varphi(f) df}$$

Two-Sided PSD



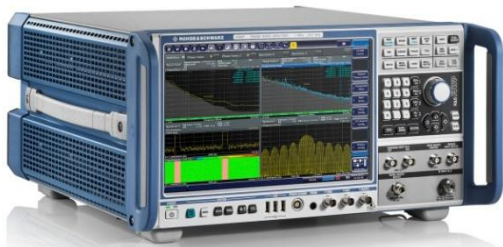
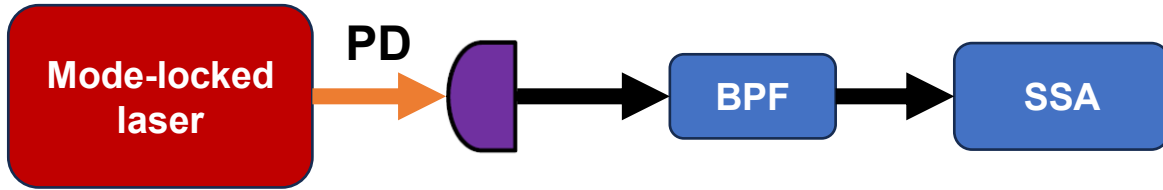
Integrated timing jitter (fs)



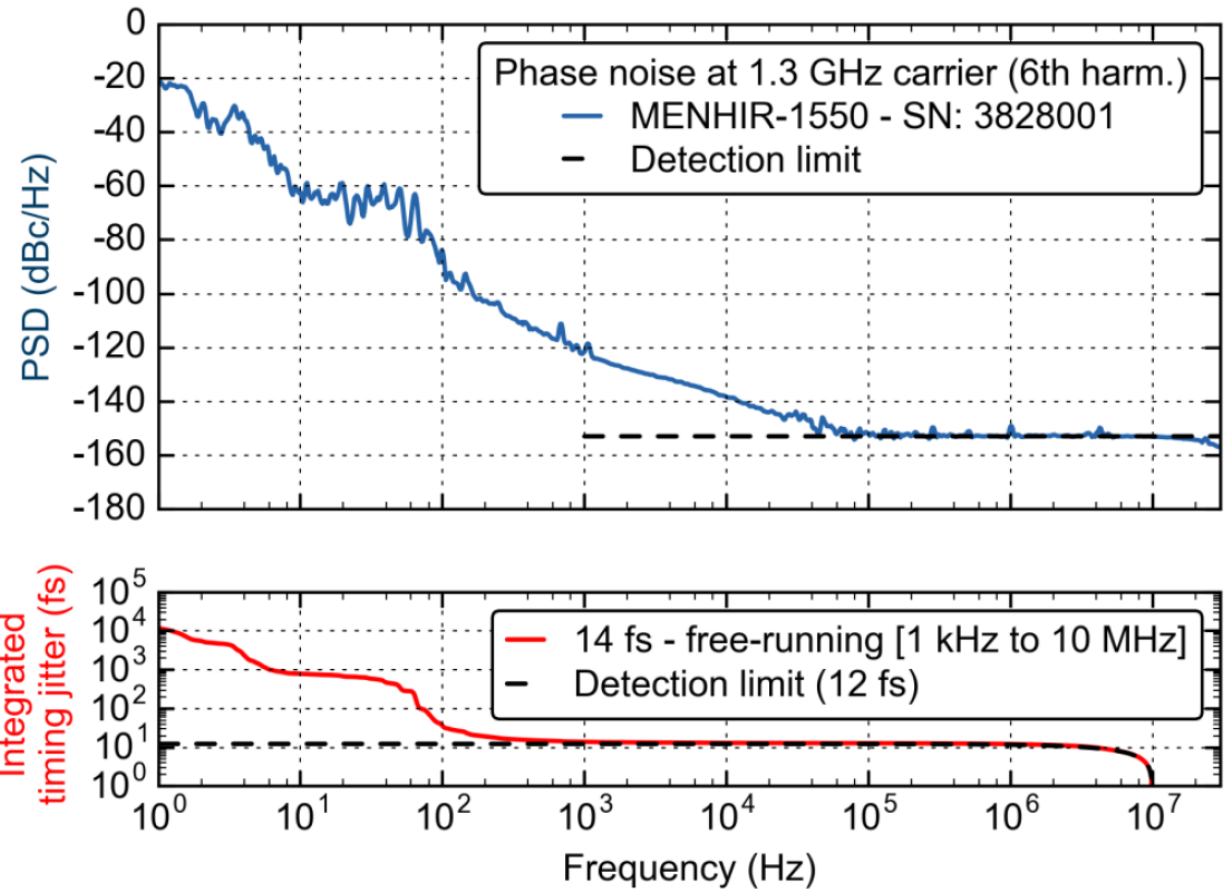
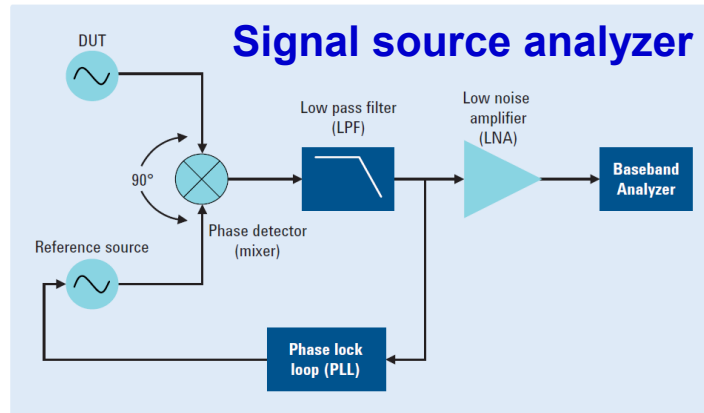
Timing drift is used for timing jitter at very lower frequency range.

Phase noise and timing jitter detection of optical signal

Direct timing jitter detection



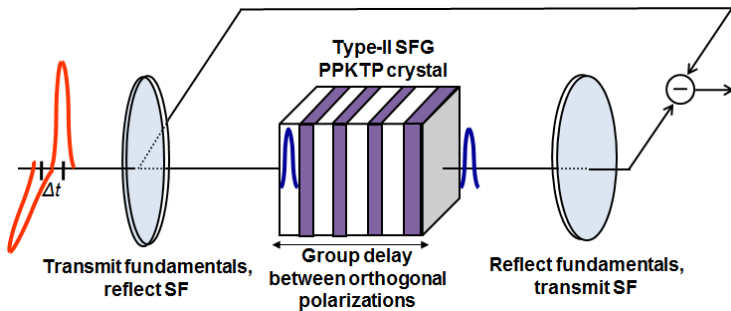
R&S FSWP



For stable signals, the measurement is limited by **AM-PM** noise.

Phase noise and timing jitter detection of optical signal

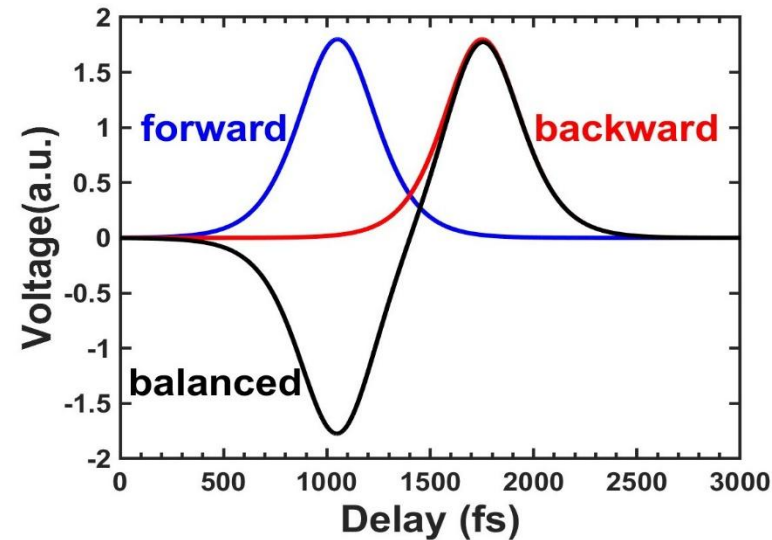
Balanced optical cross-correlator (BOC) between two signals of the same wavelength



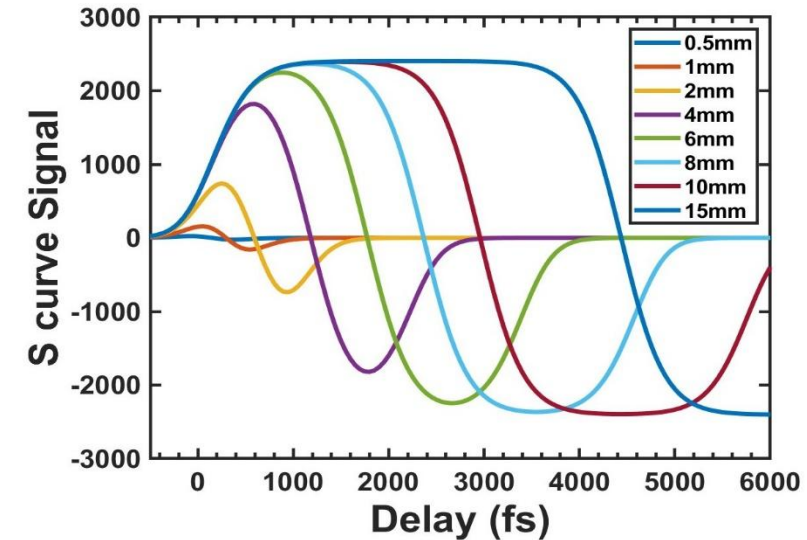
T. R. Schibli, et al., *Opt. Lett.* 28, 947–949 (2003)

- Sum frequency generation
- Type II phase matching
- Quasi-phase matching
- Balanced detection

Highly sensitive balanced detection



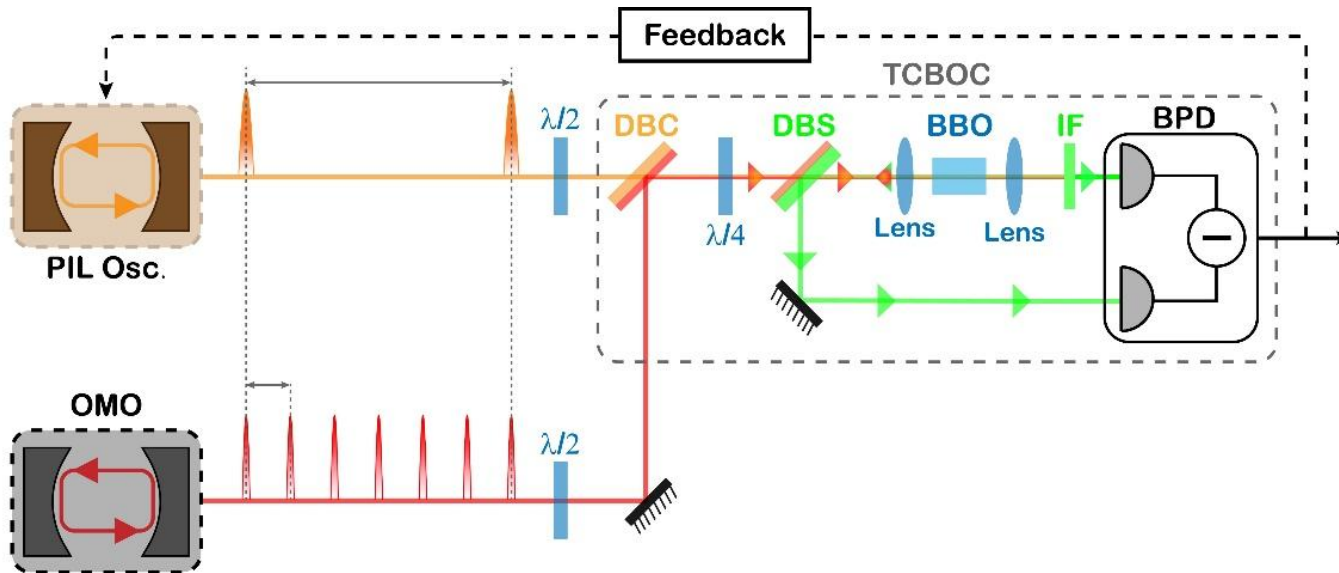
Optimized PPKTP parameter



BOC achieves **attosecond-level** sensitivity for 1550 nm, 1030 nm and 800 nm.

Phase noise and timing jitter measurement of optical signal

Two color balanced optical cross-correlator (TCBOC) for signals of different wavelengths.

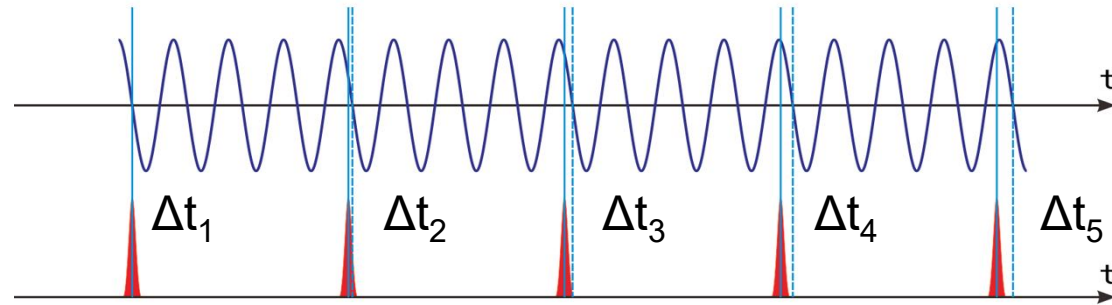


- Insensitive to laser pulse fluctuations
- Sensitivity >10 mV/fs
- Approved for solid-state and fiber lasers
- Core device of laser synchronization

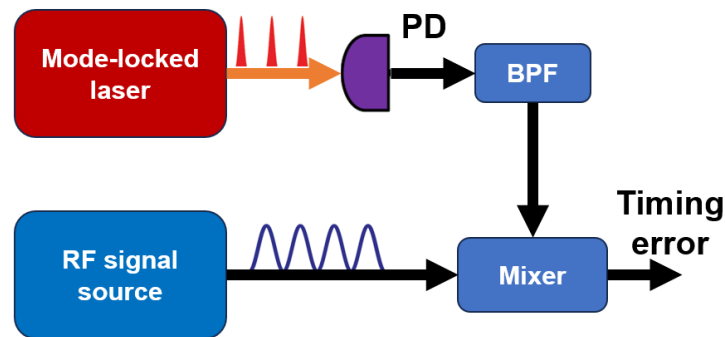
TCBOC achieves **sub-fs** synchronization for 1030/1550 nm lasers.

Phase noise and timing jitter measurement of microwave signal

Optical-Microwave



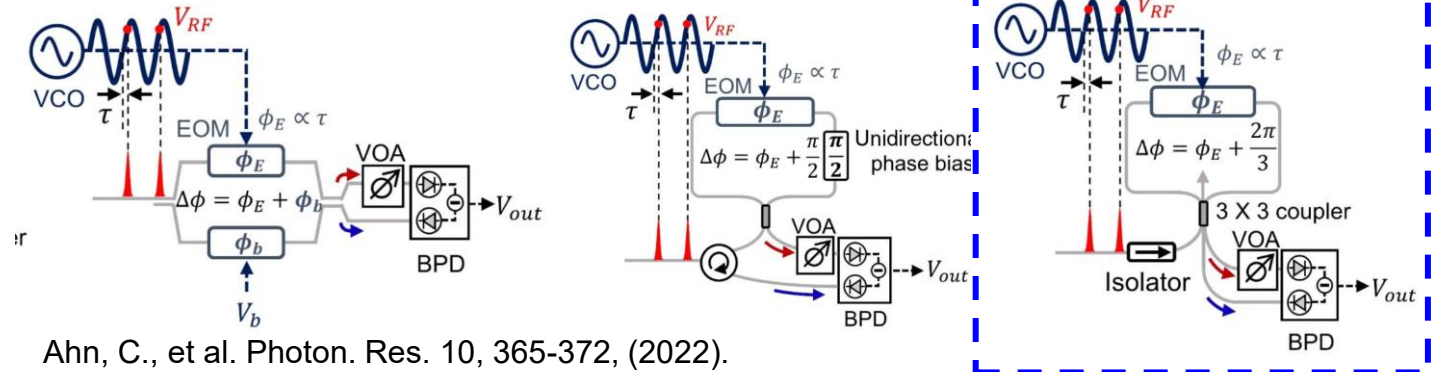
Traditional electronic method



Limited timing resolution **~50 fs**:

- amplitude-to-phase conversion of photodiode
- phase discrimination resolution of microwave mixer

Microwave photonics methods

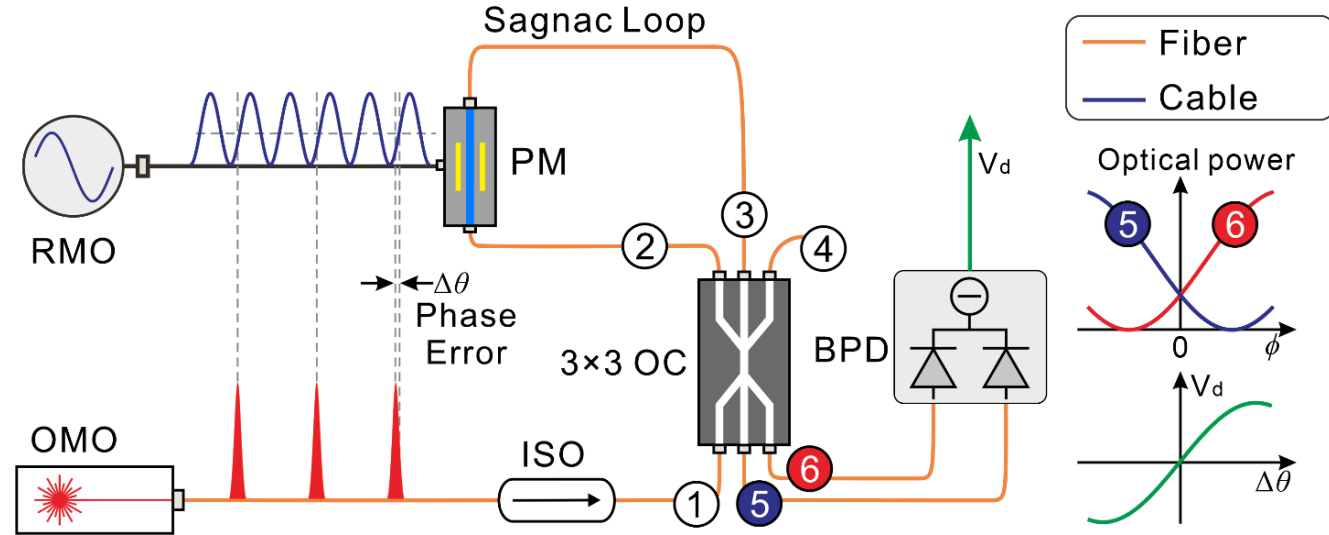


Precise timing discrimination

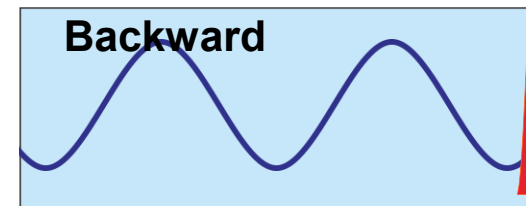
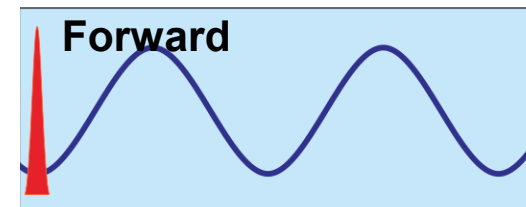
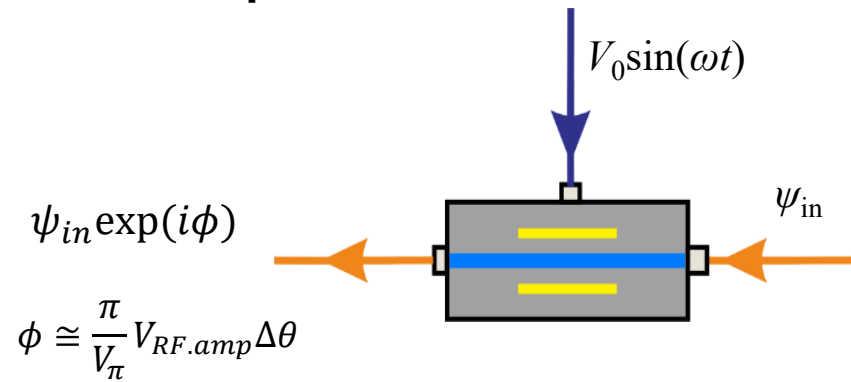
- **sub-fs** timing discrimination
- high amplitude-to-phase suppression ratio by balanced detection

Timing jitter detection of microwave signal

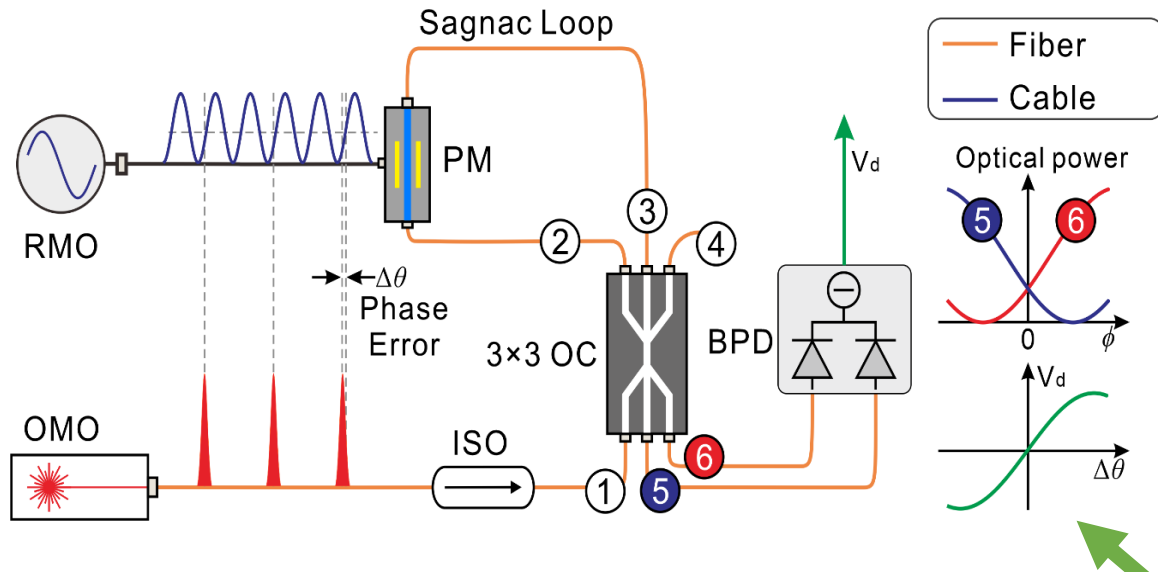
BOMPD: Balanced optical-microwave phase detector



Uni-directional phase modulator



BOMPD



- Insensitive to laser pulse fluctuations
- Sensitivity 0.5 mV/fs
- Locking BW~5 kHz

ψ_{1235} : port 1 → port 2 → port 3 → port 5

ψ_{1236} : port 1 → port 2 → port 3 → port 6

Interference

Interference

ψ_{1325} : port 1 → port 3 → port 2 → port 5

ψ_{1326} : port 1 → port 3 → port 2 → port 6

Port 5

$$\begin{aligned}
 P_5 &= |\psi_{1235} + \psi_{1325}|^2 \\
 &= \left| \frac{1}{3} \psi_1 \left[\exp(i\phi) + \exp\left(i\frac{4\pi}{3}\right) \right] \right|^2 \\
 &= \frac{4}{9} |\psi_1|^2 \cos^2\left(\frac{\phi}{2} - \frac{2\pi}{3}\right)
 \end{aligned}$$

Port 6

$$\begin{aligned}
 P_6 &= |\psi_{1236} + \psi_{1326}|^2 \\
 &= \left| \frac{1}{3} \psi_1 \exp\left(i\frac{2\pi}{3}\right) \left[\exp(i\phi) + \exp\left(i\frac{2\pi}{3}\right) \right] \right|^2 \\
 &= \frac{4}{9} |\psi_1|^2 \cos^2\left(\frac{\phi}{2} - \frac{\pi}{3}\right)
 \end{aligned}$$

$$\begin{aligned}
 V_d &= c[P_6 - P_5] = \frac{4}{9} c |\psi_1|^2 \sin\left(\frac{\pi}{3}\right) \sin(\phi) \\
 &\approx \frac{4}{9} \pi c |\psi_1|^2 \sin\left(\frac{\pi}{3}\right) \frac{V_{RF}}{V_\pi} \Delta\theta
 \end{aligned}$$

BOMPD

RF input

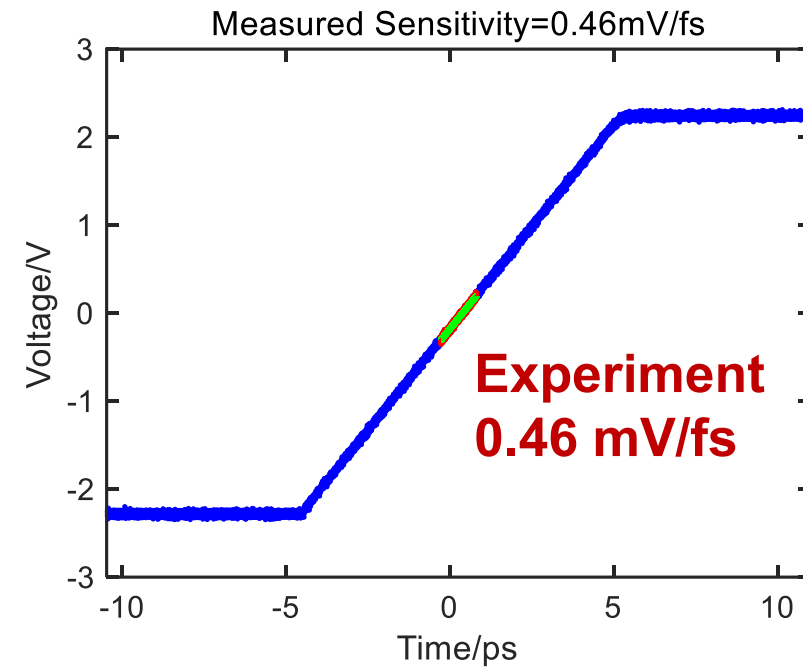
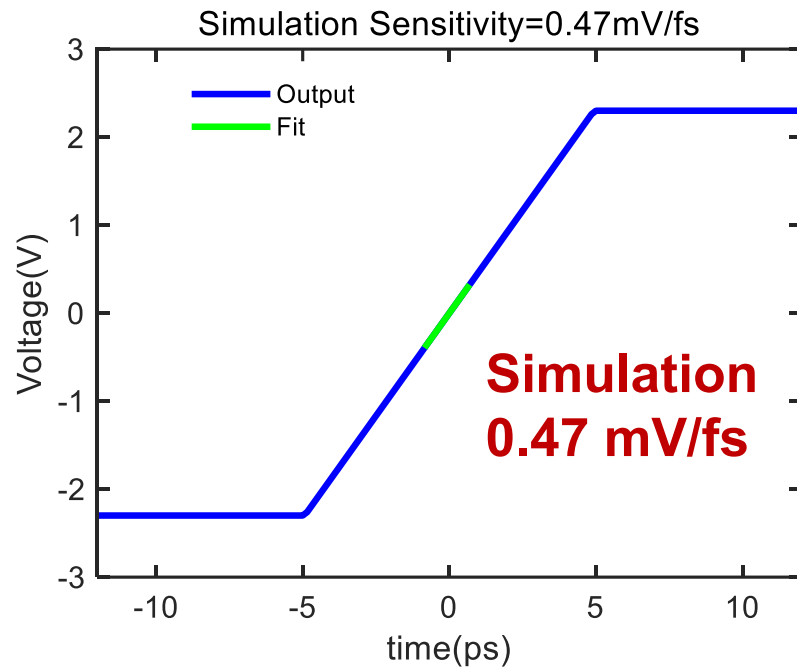
- $P = 8.5 \text{ dBm}$
- $f = 1.3 \text{ GHz}$

Laser input

- $P = 15 \text{ mW}$
- $f' = 216.667 \text{ MHz}$
- $\lambda = 1550 \text{ nm}$

BPD

- $\text{Gain} = 0.5e5 \text{ V/A}$



Simulation result meets experiment's very well !

Timing and synchronization system for FEL facilities

Task:

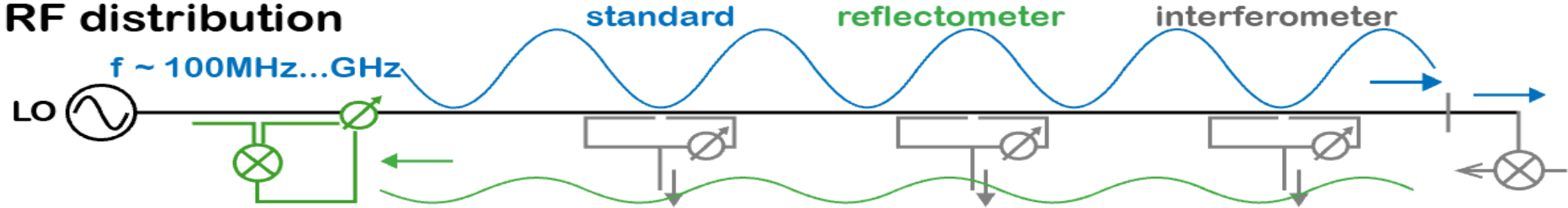
- Providing femtosecond level stable master clock;
- Distributing fs-level reference signal over the whole FEL facility for slave lasers, LLRF, and diagnostics;
- Tight locking of optical and microwave clients;

Function:

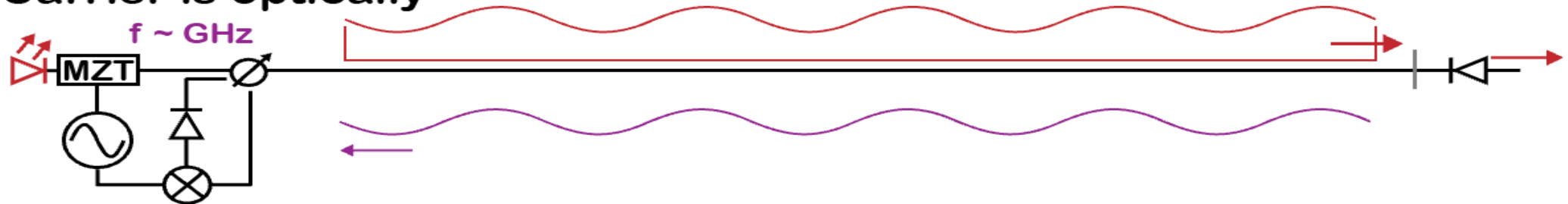
- Stable operation of FEL;
- Highest time-resolution for diagnostics and pump-probe experiments;

Synchronization methods

1) RF distribution



2) Carrier is optically



3) Pulsed optical source



Adapted from Dr. Holger Schlarb slide (DESY)

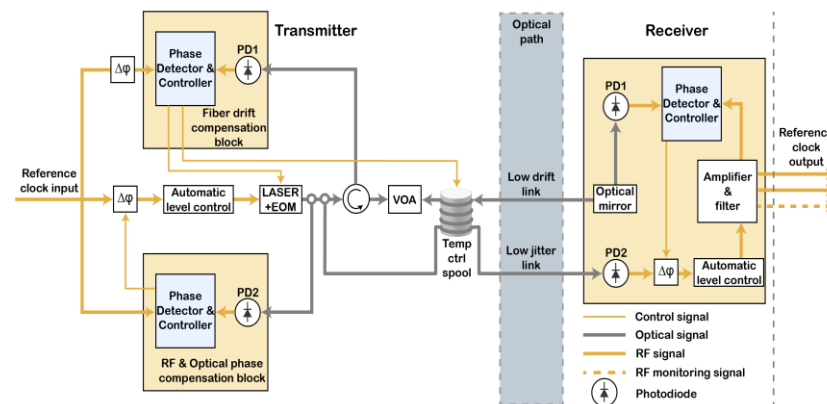
Synchronization methods

Various synchronization methods have been developed

Approach		Developed by	Jitter fs[10Hz,10MHz]	Drift fs/24h/km	Advantage	Disadvantage
Fiber Link	Pulsed laser + PMF	MIT-CFEL DESY	<0.5	~5	Lowest jitter Drift free Large distance	Expensive
	CW laser + SMF	PSI LBNL	<20	~40	Low jitter Long distance	Certain drift
RF Link	Temperature stabilized RF	PAL SACLA	<10	>250	Stable	Large drift



European XFEL



SwissFEL



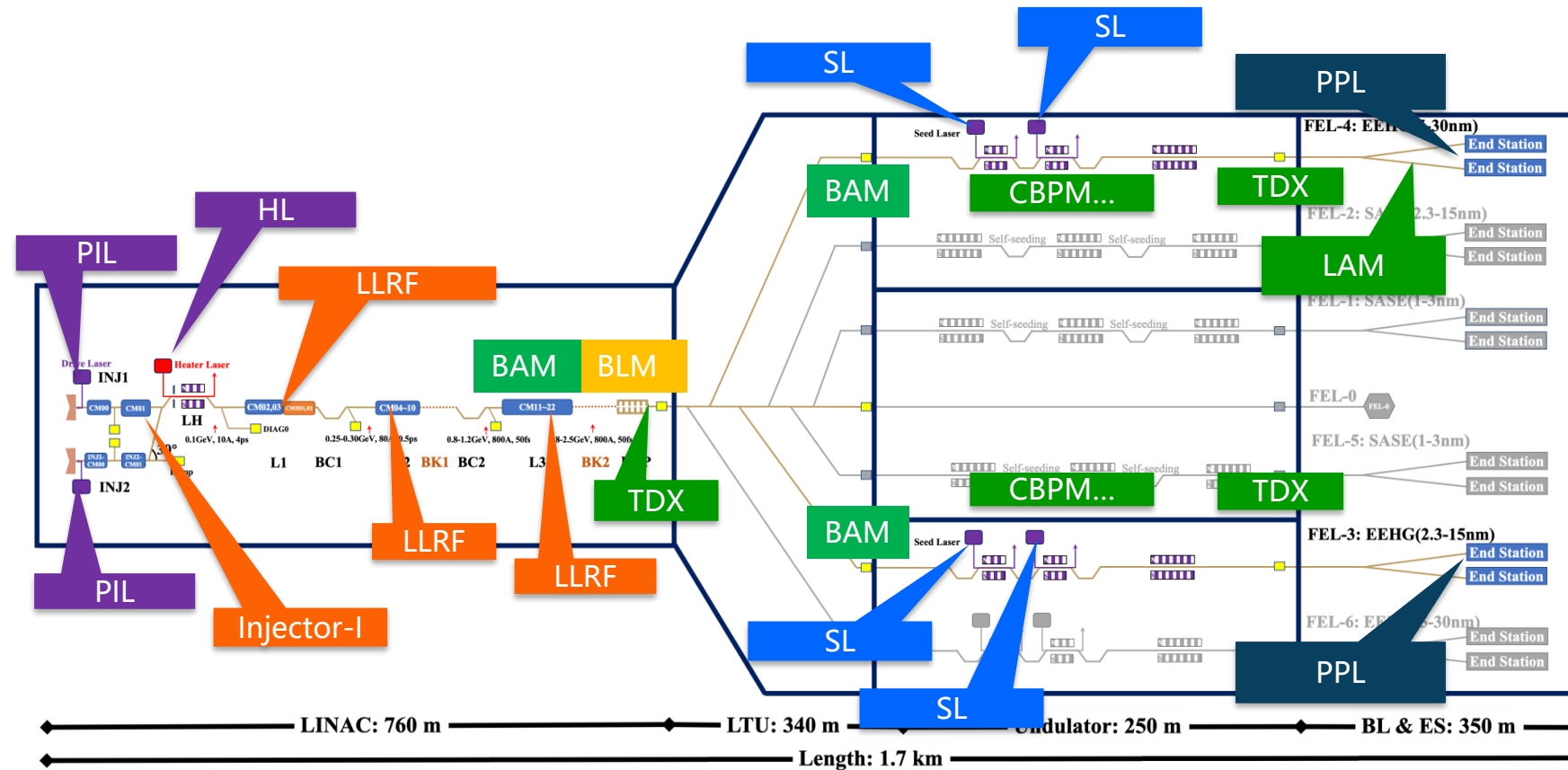
PAL XFEL

Synchronization methods

Hybrid SYN with pulsed laser link as the backbone has become the mainstream design of large-scale FELs.

Facility	RF	CW Laser Link	Pulsed Laser Link
LCLS	•		•
SwissFEL	•	•	•
FERMI	•	•	•
PAL XFEL	•	•	
SACLA	•	•	
SXFEL	•	•	•
DCLS	•		•
LCLS-II	•	•	•
FLASH	•		•
EuXFEL	•		•
SHINE	•		•
S ³ FEL	•	•	•
DALS	•	•	•

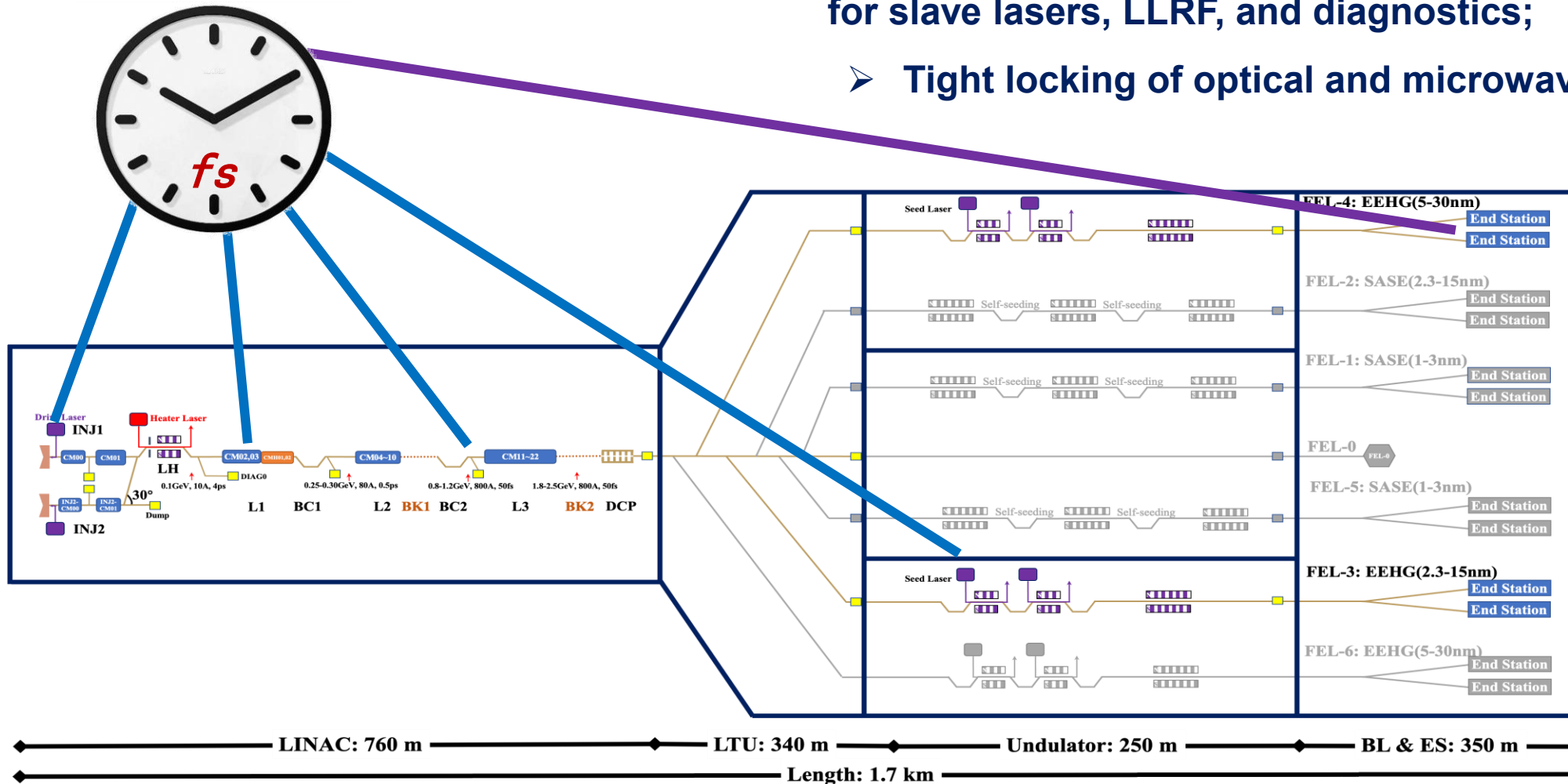
Requirements of Synchronization System for S³FEL



High-precision, kilometer scale, hundreds clients, robustness

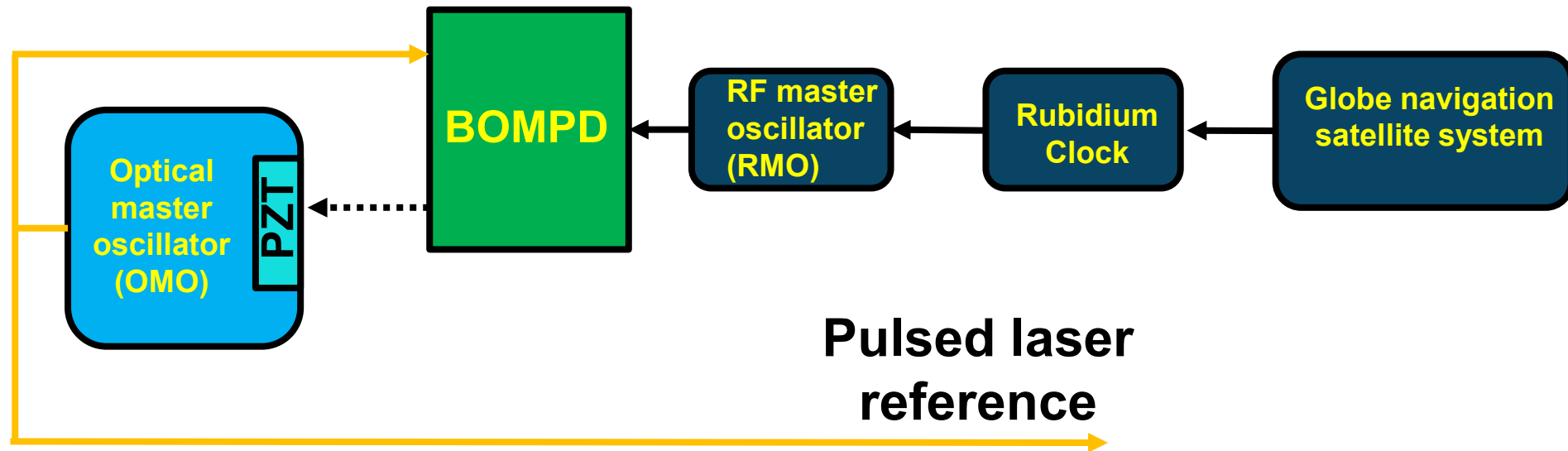
S³FEL Synchronization System

- Providing femtosecond level stable reference signal;
- Distributing reference signal over the whole FEL facility for slave lasers, LLRF, and diagnostics;
- Tight locking of optical and microwave clients;



Master Clock

Master clock provides ultra-stable timing reference



Architecture diagram

GPSDRO + RMO + OMO → ultra-stable master clock with extremely low jitter and outstanding long-term accuracy.

Master Clock

Main components of master clock

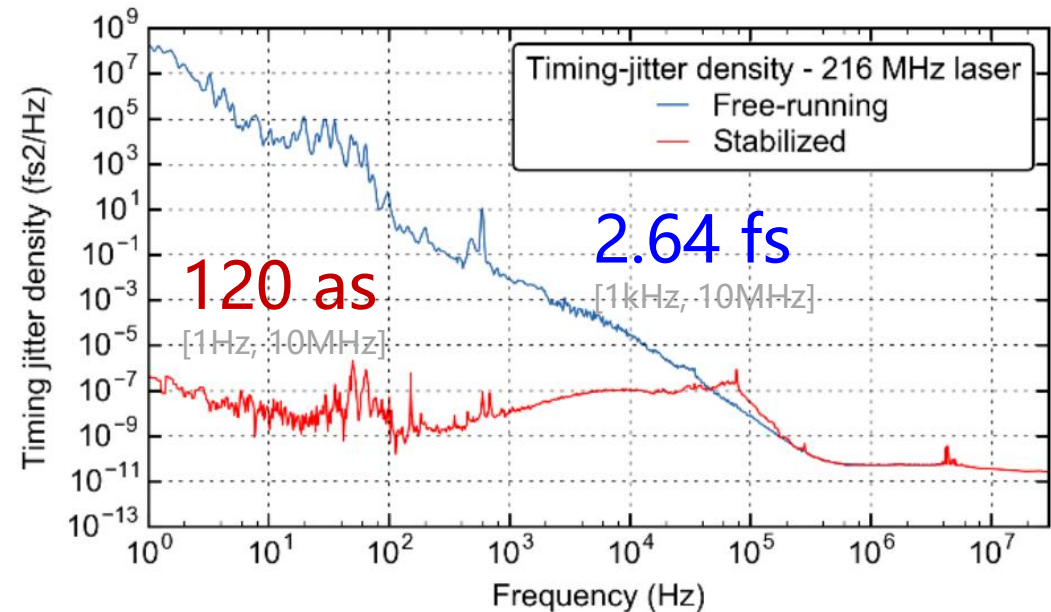
GPSDRO + RMO

- Commercial RF signal oscillator
- 1.3 GHz
- Ultra-low phase noise
- 24/7 operation

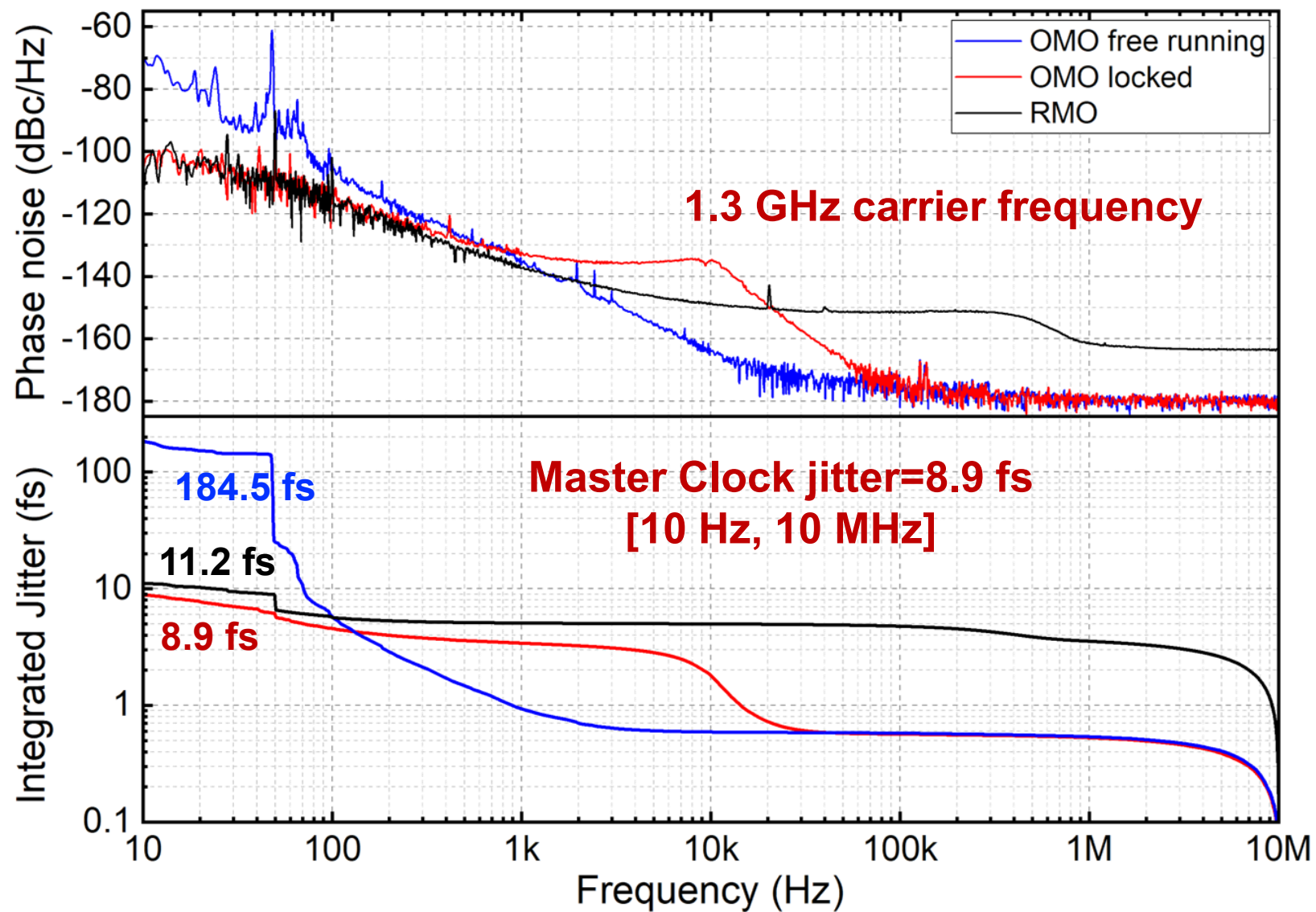


OMO

- Commercial mode-locked laser
- 216.667 MHz, 1550 nm
- Ultra-low jitter
- 24/7 operation

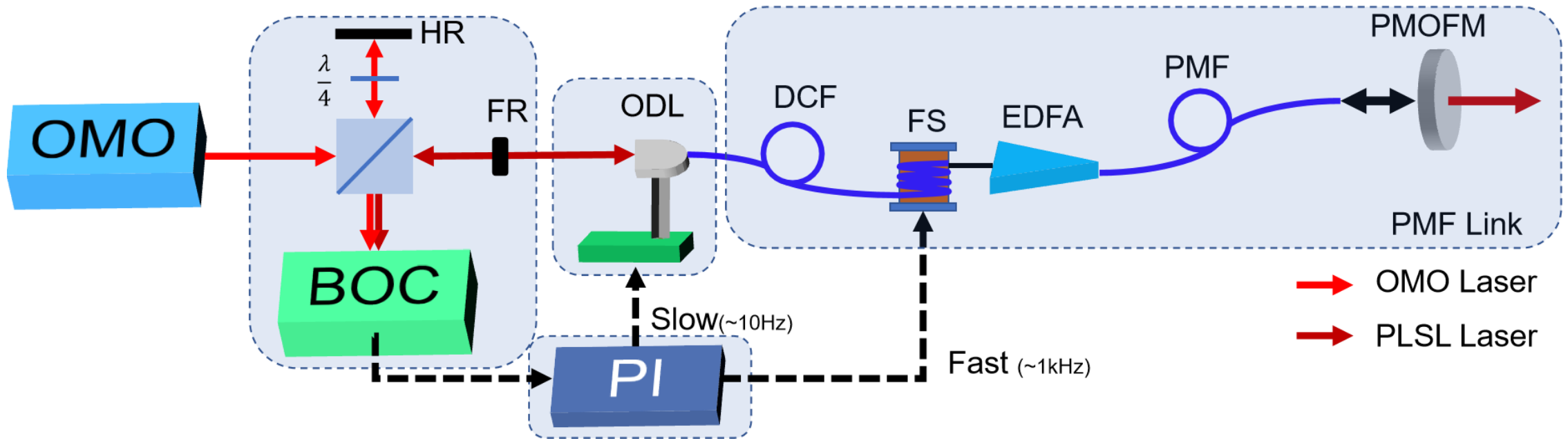


Master clock



Stabilized reference distribution

Pulse laser phase-stabilized link (PLSL) distributes reference signal to laser/ RF clients with sub-fs level additive jitter.



Schematic diagram of PLSL

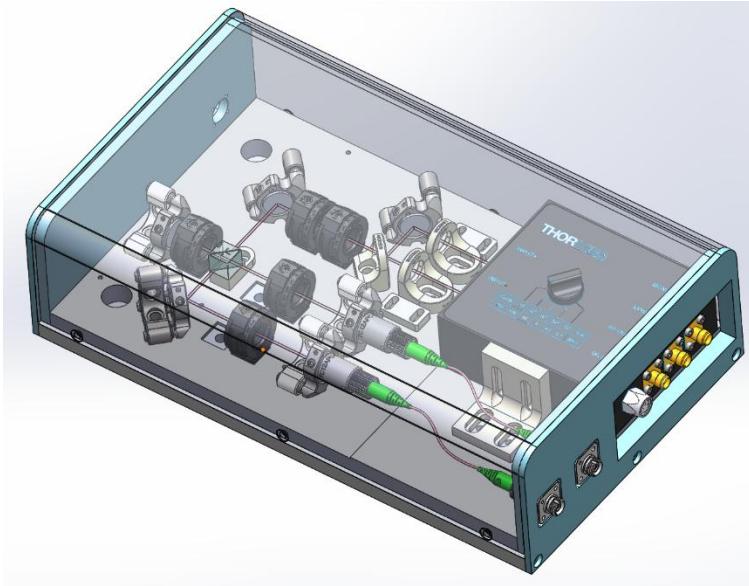
Pulse laser phase stabilized link

Timing detection

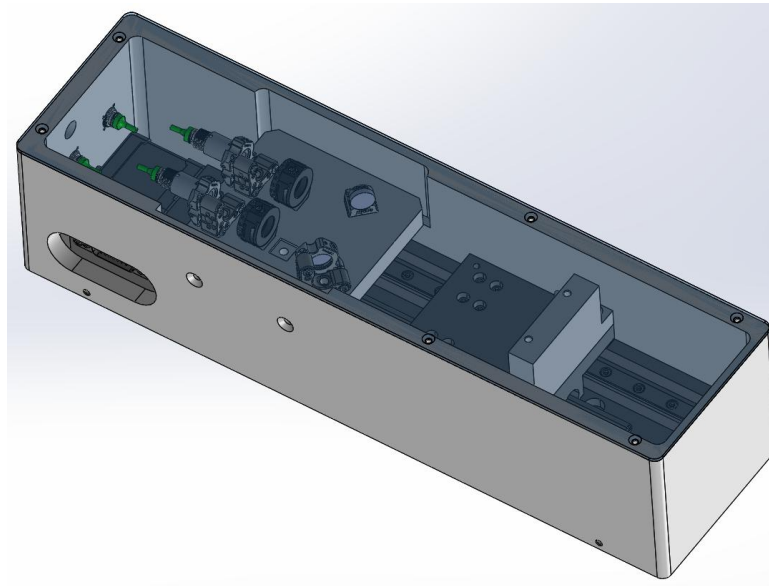
- **Balanced optical cross-correlation (BOC)**
- **Insensitive to laser pulse fluctuations**
- **Sensitivity 15 mV/fs**

Feedback

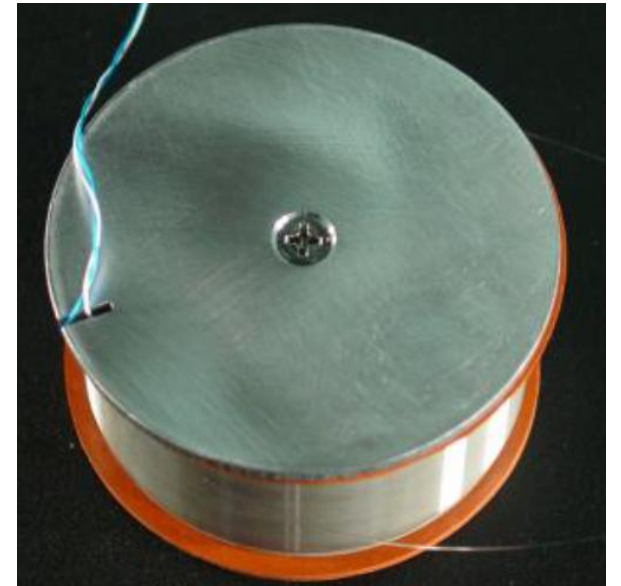
- **Fast: Piezo-based fiber stretcher BW~5 kHz**
- **Slow: Optical delay line ~1.3 ns**



BOC-1550 nm



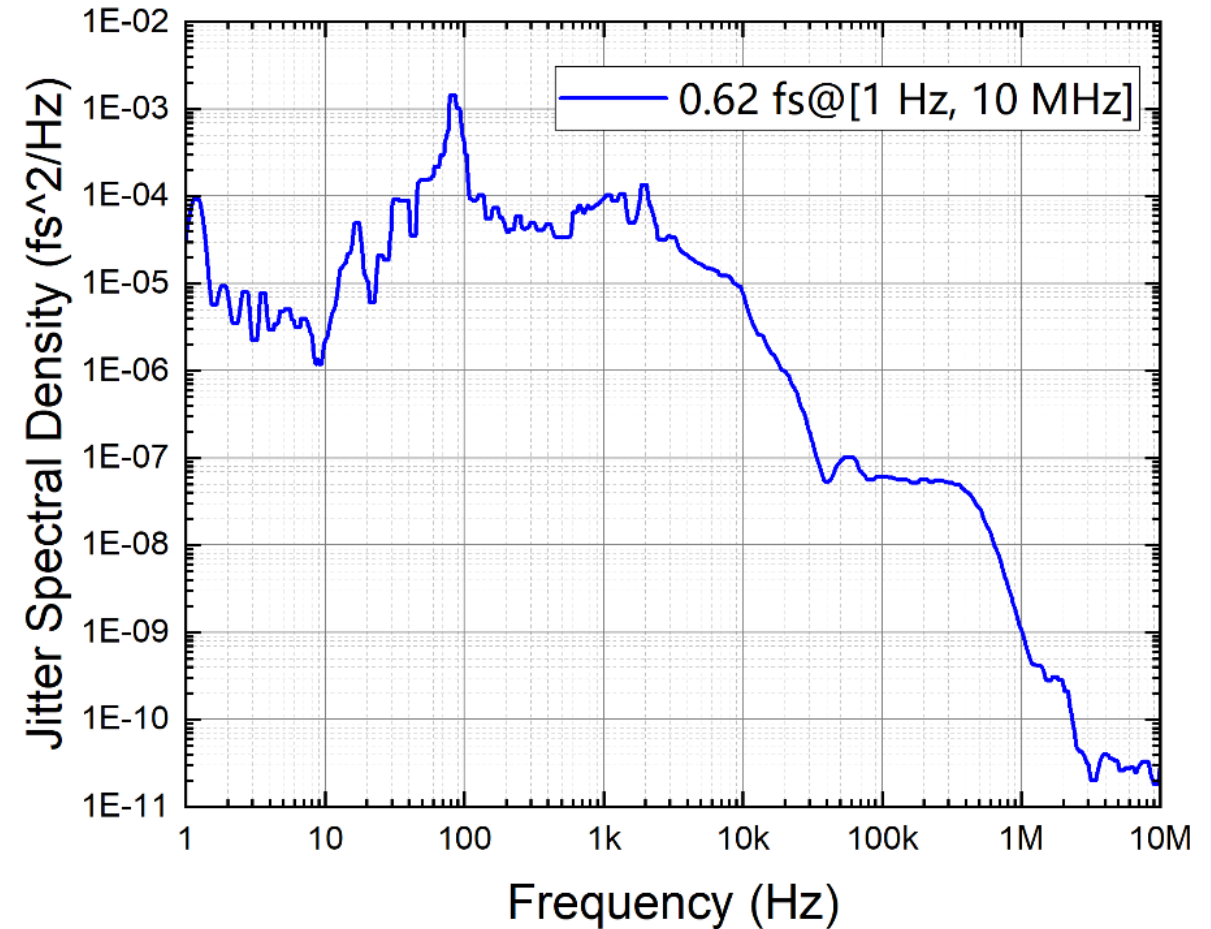
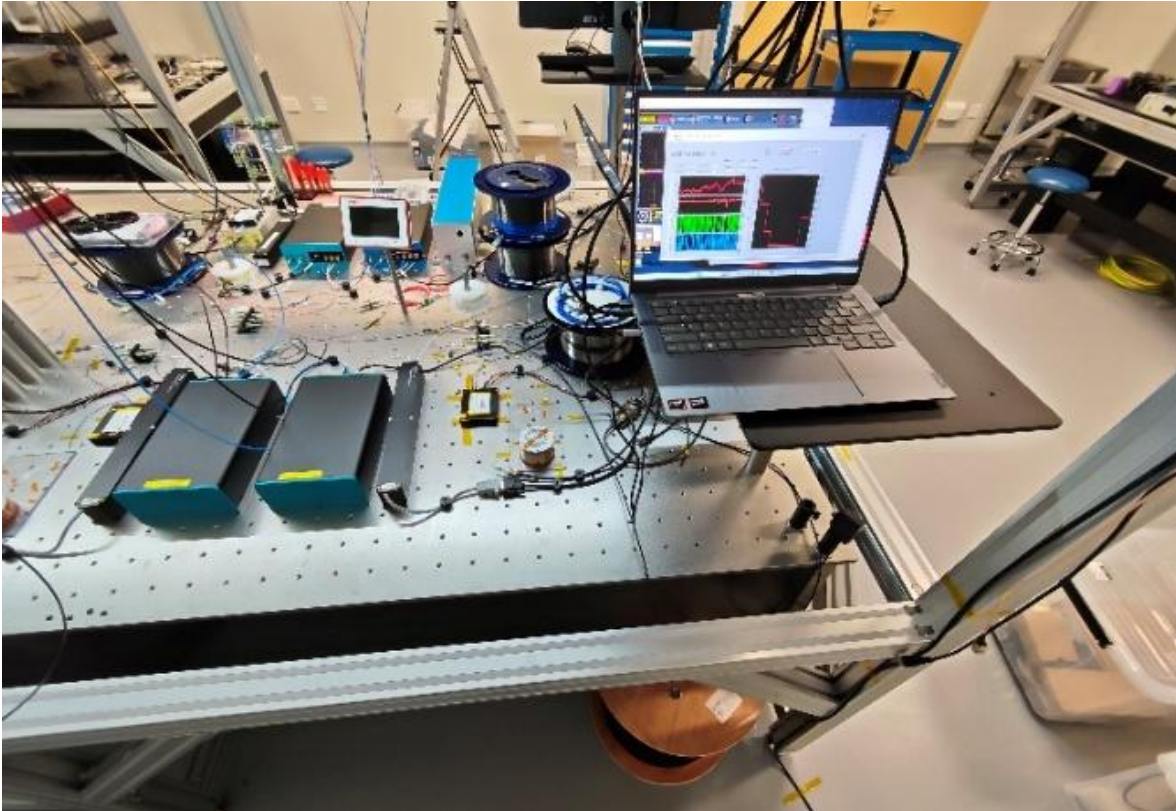
Optical delay line



**Piezo-based
fiber stretcher**

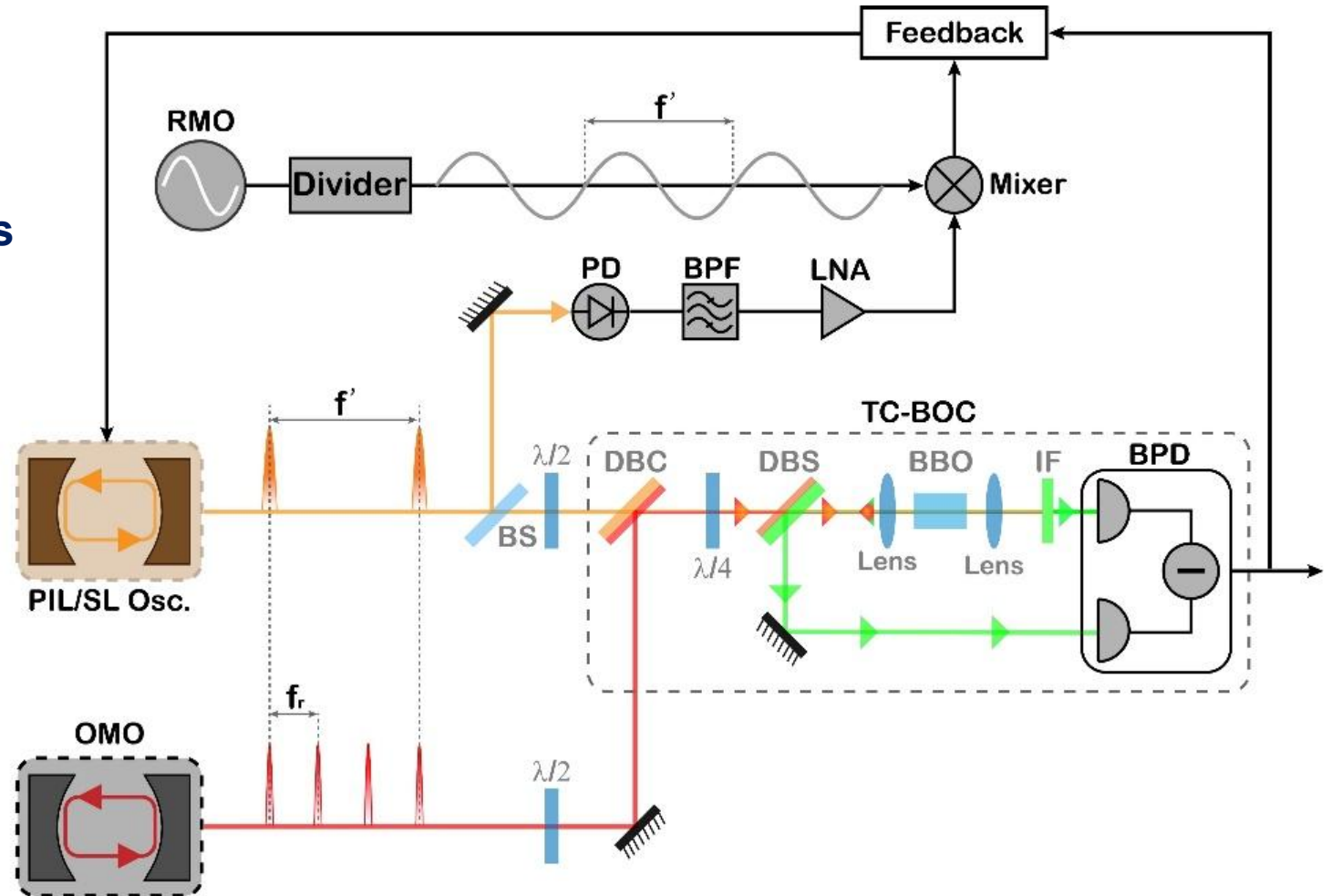
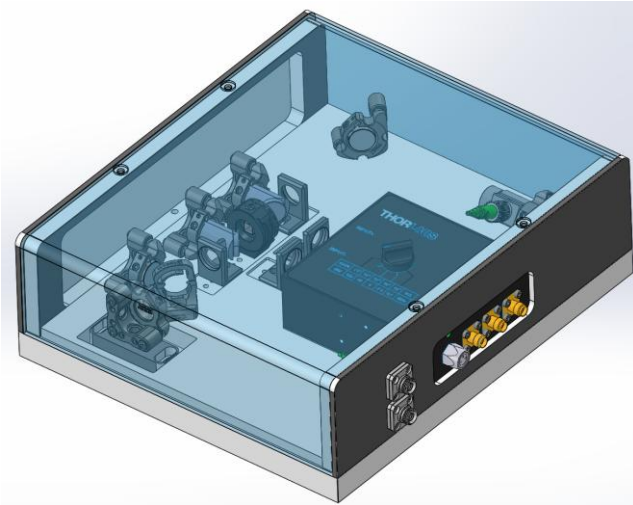
Pulse laser phase stabilized link

PLSL- 2 km has achieved sub-fs jitter.

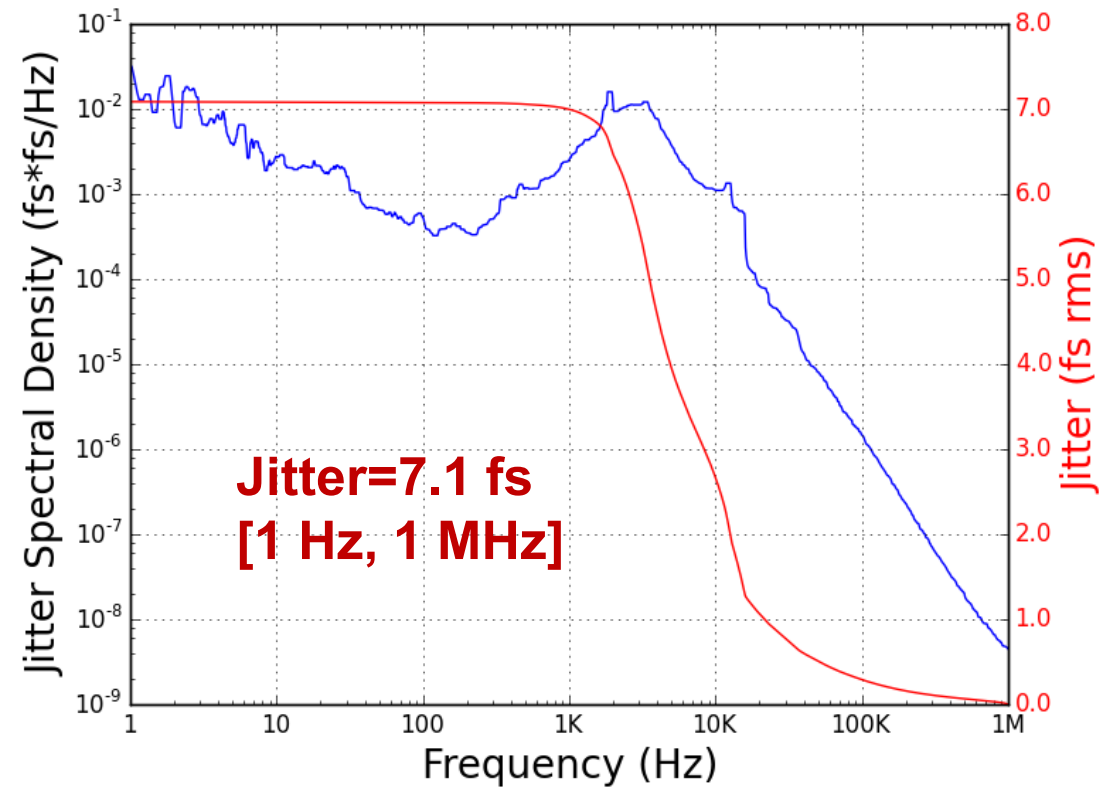
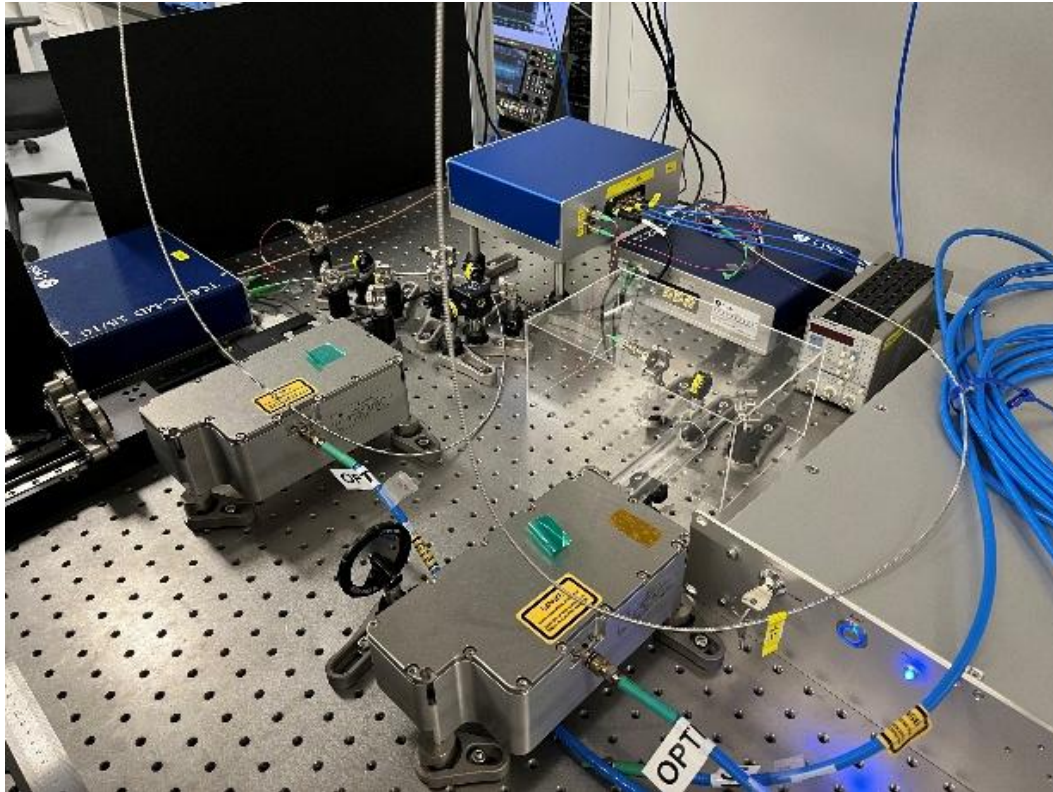


Locking of laser clients

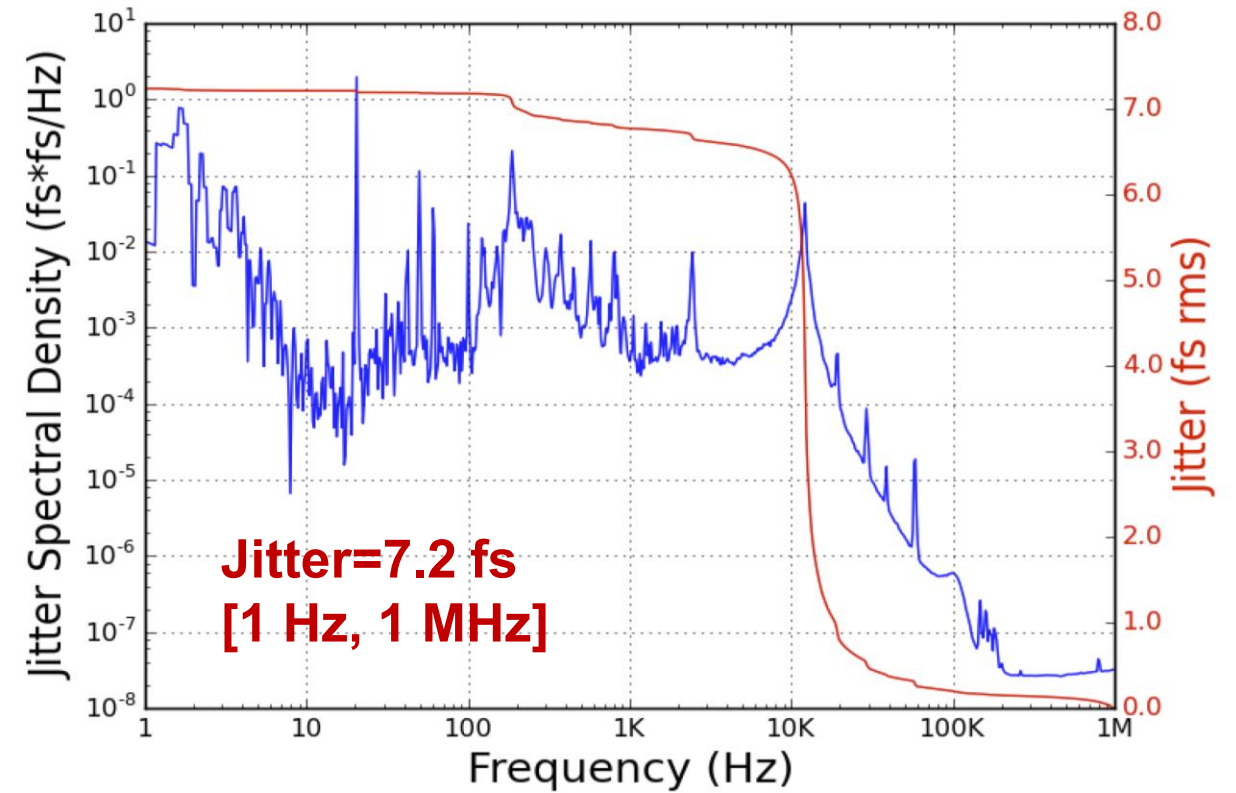
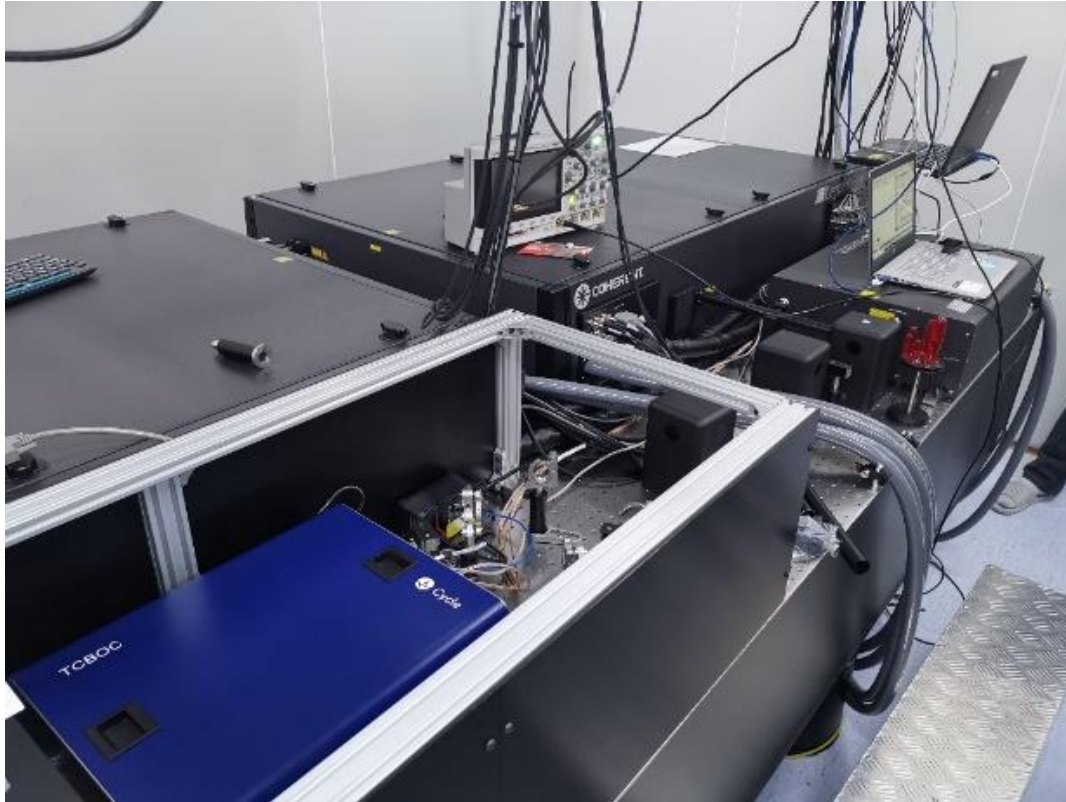
- **Two-color balanced optical cross-correlation (TCBOC)**
- **RF Pre-lock function**
- **Insensitive to laser pulse fluctuations**
- **Slope 10 mV/fs**
- **Piezo-based feedback BW~5 kHz**



Locking of Yb lasers at 1030 nm

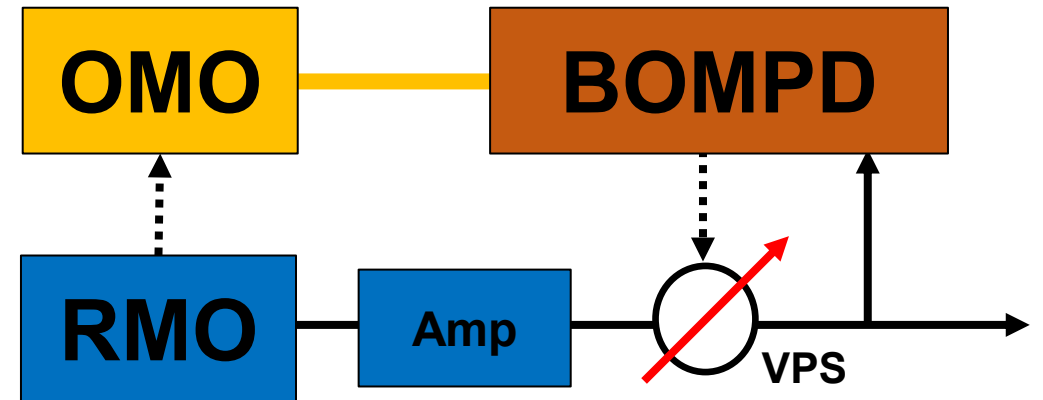
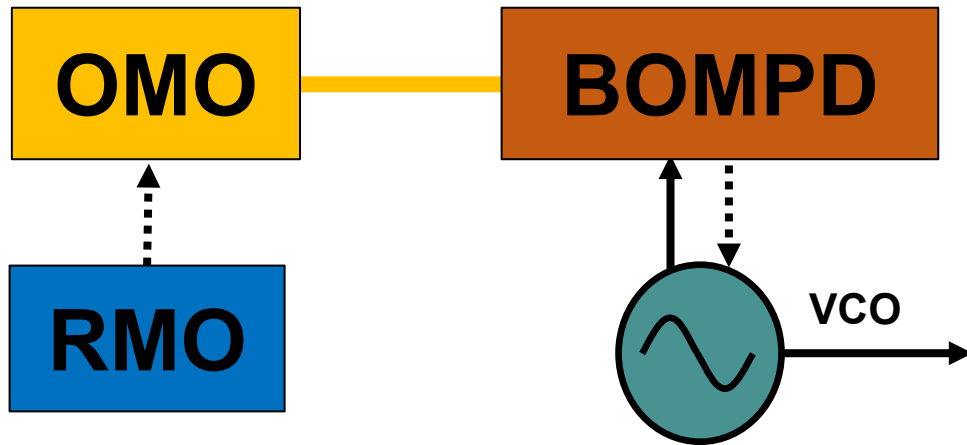


Locking of Ti:Sa lasers at 800 nm



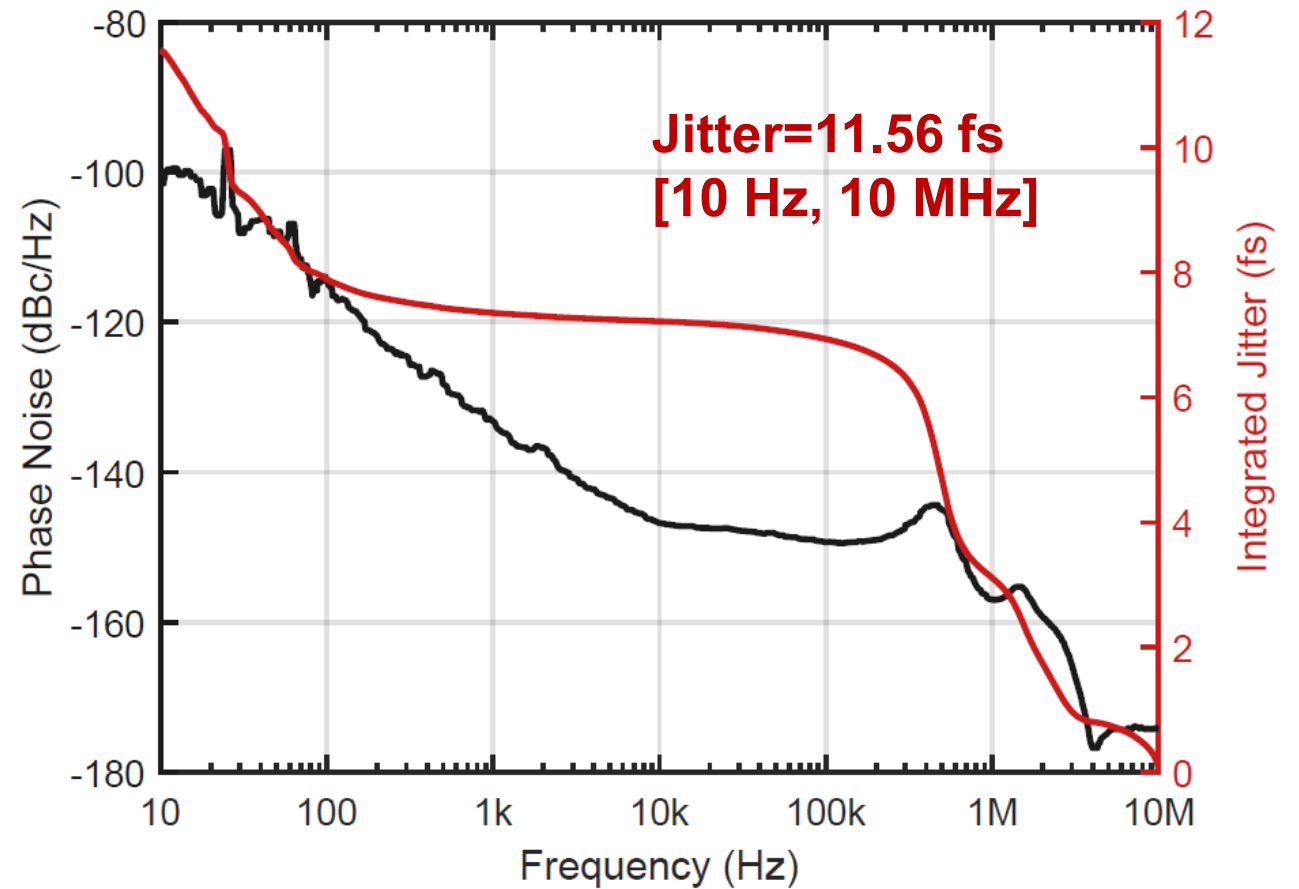
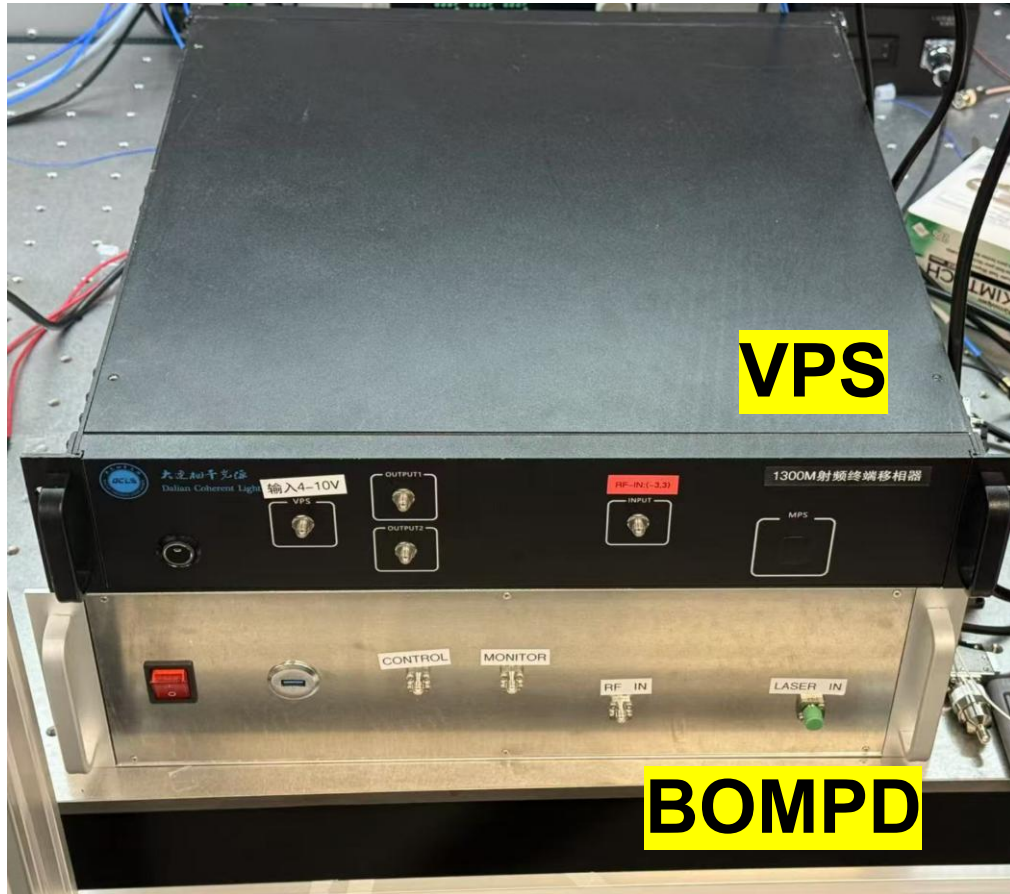
Locking of RF clients

- **Balanced optical-microwave phase detector (BOMPD)**
- **Insensitive to laser pulse fluctuations**
- **Sensitivity 1 mV/fs**
- **Locking BW~1 kHz**



High performance VCO is embargoed; VPS solution meets FEL requirements.

Locking of RF clients @ 1.3 GHz



Summary

- **Ultrafast laser systems and fs timing & synchronization system are critical sub-systems for FEL facilities.**
- **High R.R. FEL imposes stringent requirements on ultrafast laser systems, especially in achieving high pulse energy and power simultaneously, necessitating cutting-edge laser technologies.**
- **These systems are required to operate continuously 24/7, with long-term stabilities.**

Acknowledgement



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Weiyu Lai

He Zhao

Jie Chen

Junsong Ding

Baoqi Xue



Thanks!

